Spectral Method for Unbounded Domains

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May 8, 2013

1 Laguerre Polynomials/Functions

1.1 (Generalized) Laguerre Polynomials

• Laguerre polynomials, denoted by $\mathcal{L}_n(x)$, are orthogonal w.r.t. $\omega(x) = e^{-x}$, on the half line $\mathbb{R}_+ := (0, +\infty)$, i.e.

$$\int_{0}^{+\infty} \mathcal{L}_{n}(x)\mathcal{L}_{m}(x)e^{-x} \,\mathrm{d}x = \delta_{m\,n}.$$
 (1)

• The Generalized Laguerre polynomials (GLPs), denoted by $\mathcal{L}_{n}^{(\alpha)}(x), \alpha > -1$ are orthongal w.r.t. $\omega_{\alpha}(x) = x^{\alpha} e^{-x}$ on \mathbb{R}_{+} , i.e.

$$\int_{0}^{+\infty} \mathcal{L}_{n}(x)\mathcal{L}_{m}(x)\omega_{\alpha}(x) \,\mathrm{d}x = \gamma_{n}^{(\alpha)}\delta_{m\,n},\qquad(2)$$

where

$$\gamma_n^{(\alpha)} = \frac{\Gamma(n+\alpha+1)}{n!}.$$
(3)

Basic properties of GLPs

1. Three-term recurrence formula

 $(n+1)\mathcal{L}_{n+1}^{(\alpha)}(x) = (2n+\alpha+1-x)\mathcal{L}_{n}^{(\alpha)}(x) - (n+\alpha)\mathcal{L}_{n-1}^{(\alpha)}(x),$ $\mathcal{L}_{0}^{(\alpha)}(x) = 1, \quad \mathcal{L}_{1}^{(\alpha)}(x) = -x+\alpha+1.$ Leading coefficient: $k_{n}^{(\alpha)} = \frac{(-1)^{n}}{1-1},$

2. Leading coefficient: $k_n^{(\alpha)} = \frac{(-1)^n}{n!}$, value at 0: $\mathcal{L}_n^{(\alpha)}(0) = \frac{\Gamma(n+\alpha+1)}{n!\Gamma(\alpha+1)}$ 3. Sturm–Liouville equation $(\lambda_n = n !)$

$$x^{-\alpha} e^x \partial_x \left(x^{\alpha+1} e^{-x} \partial_x \mathcal{L}_n^{(\alpha)}(x) \right) + \lambda_n \, \mathcal{L}_n^{(\alpha)}(x) = 0, \qquad (4)$$

$$x \,\partial_x^2 \mathcal{L}_n^{(\alpha)}(x) + (\alpha + 1 - x) \,\partial_x \mathcal{L}_n^{(\alpha)}(x) + \lambda_n \mathcal{L}_n^{(\alpha)}(x) = 0.$$
(5)

4.
$$\left\{\partial_{x}\mathcal{L}_{n}^{(\alpha)}\right\}$$
 orthogonal w.r.t. $\omega_{\alpha+1}$
$$\int_{0}^{+\infty} \partial_{x}\mathcal{L}_{n}^{(\alpha)}(x) \,\partial_{x}\mathcal{L}_{m}^{(\alpha)}(x) \,\omega_{\alpha+1}(x) \,\mathrm{d}x = \lambda_{n}\gamma_{n}^{(\alpha)} \,\delta_{mn}.$$
(6)

5. Rodrigues's formula

$$\mathcal{L}_{n}^{(\alpha)}(x) = \frac{x^{-\alpha} e^{x}}{n!} \frac{\mathrm{d}^{n}}{\mathrm{d}x^{n}} \{x^{n+\alpha} e^{-x}\}.$$
(7)

Explicit expression

$$\mathcal{L}_{n}^{(\alpha)}(x) = \sum_{k=0}^{n} \frac{(-1)^{k}}{k!} \binom{n+a}{n-k} x^{k}$$

6. Derivatives

$$\partial_{x}\mathcal{L}_{n}^{(\alpha)}(x) = -\mathcal{L}_{n-1}^{(\alpha+1)}(x) = -\sum_{k=0}^{n-1} \mathcal{L}_{k}^{(\alpha)}(x),$$

$$\mathcal{L}_{n}^{(\alpha)}(x) = \partial_{x}\mathcal{L}_{n}^{(\alpha)}(x) - \partial_{x}\mathcal{L}_{n+1}^{(\alpha)}(x),$$

$$x \partial_{x}\mathcal{L}_{n}^{(\alpha)}(x) = n\mathcal{L}_{n}^{(\alpha)}(x) - (n+\alpha)\mathcal{L}_{n-1}^{(\alpha)}(x),$$
(8)

- 7. Asymptotic properties of the GLPs.
 - For large x, $\mathcal{L}_{n}^{(\alpha)}(x)$ is dominated by leading term $\frac{(-1)^{n}}{n!}x^{n}$
 - Theorem 8.22.5 of Szegö (1975)

$$\mathcal{L}_{n}^{(\alpha)}(x) = \sqrt{\frac{n^{\alpha - \frac{1}{2}e^{x}}}{\pi x^{\alpha + \frac{1}{2}}}} \left[\cos\left(2\sqrt{nx} - \frac{2\alpha + 1}{4}\pi\right) + \frac{O(1)}{\sqrt{nx}} \right]$$
(9)

for all $x \in [c/n, b]$.

• upper bounds

$$\left|\mathcal{L}_{n}^{(\alpha)}(x)\right| \leqslant \begin{cases} \mathcal{L}_{n}^{(\alpha)}(0) e^{x/2}, & \text{if } \alpha \ge 0, \\ \left(2 - \mathcal{L}_{n}^{(\alpha)}(0)\right) e^{x/2}, & \text{if } \alpha < 0. \end{cases}$$
(10)



1.2 Generalized Laguerre Functions (GLFs)

Definition:

$$\hat{\mathcal{L}}_n^{(\alpha)}(x) := e^{-x/2} \mathcal{L}_n^{(\alpha)}(x), \quad x \in \mathbb{R}_+, \quad \alpha > -1.$$

GLFs are orthogonal w.r.t weight function $\hat{\omega}_{\alpha} = x^{\alpha}$, i.e.

$$\int_{0}^{+\infty} \hat{\mathcal{L}}_{n}^{(\alpha)}(x) \hat{\mathcal{L}}_{n}^{(\alpha)}(x) \hat{\omega}_{\alpha}(x) \,\mathrm{d}x = \gamma_{n}^{(\alpha)} \delta_{n\,m}.$$
 (11)

Introducing

$$\hat{\partial}_x = \partial_x + \frac{1}{2} \tag{12}$$

then

$$\partial_x \mathcal{L}_n^{(\alpha)}(x) = e^{x/2} \hat{\partial}_x \hat{\mathcal{L}}_n^{(\alpha)}(x).$$
(13)

Basic properties of GLFs

• Three-term recurrence relation

$$(n+1)\hat{\mathcal{L}}_{n+1}^{(\alpha)}(x) = (2n+\alpha+1-x)\hat{\mathcal{L}}_{n}^{(\alpha)}(x) - (n+\alpha)\hat{\mathcal{L}}_{n-1}^{(\alpha)}(x),$$
$$\hat{\mathcal{L}}_{0}^{(\alpha)}(x) = e^{-x/2}, \quad \hat{\mathcal{L}}_{1}^{(\alpha)}(x) = (\alpha+1-x)e^{-x/2}.$$

• Decay property

$$\left| \hat{\mathcal{L}}_{n}^{(\alpha)}(x) \right| \to 0, \quad x \to +\infty$$

$$\left| \hat{\mathcal{L}}_{n}^{(\alpha)}(x) \right| \leqslant 1, \quad \forall x \in [0, +\infty)$$

$$(14)$$

• Sturm–Liouville equation

$$x^{-\alpha}e^{x/2}\partial_x\left(x^{\alpha+1}e^{-x/2}\hat{\partial}_x\hat{\mathcal{L}}_n^{(\alpha)}(x)\right) + n\,\hat{\mathcal{L}}_n^{(\alpha)}(x) = 0.$$
(16)

• Orthogonality

$$\int_{0}^{+\infty} \hat{\partial}_{x} \hat{\mathcal{L}}_{n}^{(\alpha)}(x) \,\hat{\partial}_{x} \hat{\mathcal{L}}_{m}^{(\alpha)}(x) \,\hat{\omega}_{\alpha+1}(x) \,\mathrm{d}x = \lambda_{n} \gamma_{n}^{(\alpha)} \,\delta_{mn}.$$
(17)

• Derivative relation

$$\hat{\partial}_{x}\hat{\mathcal{L}}_{n}^{(\alpha)}(x) = -\hat{\mathcal{L}}_{n-1}^{(\alpha+1)}(x) = -\sum_{k=0}^{n-1} \hat{\mathcal{L}}_{k}^{(\alpha)}(x),
\hat{\mathcal{L}}_{n}^{(\alpha)}(x) = \hat{\partial}_{x}\hat{\mathcal{L}}_{n}^{(\alpha)}(x) - \hat{\partial}_{x}\hat{\mathcal{L}}_{n+1}^{(\alpha)}(x),
x\,\hat{\partial}_{x}\mathcal{L}_{n}^{(\alpha)}(x) = n\hat{\mathcal{L}}_{n}^{(\alpha)}(x) - (n+\alpha)\hat{\mathcal{L}}_{n-1}^{(\alpha)}(x),$$
(18)

1.3 Laguerre–Gauss Type Quadarature

Theorem 1. Let $\{x_j^{\alpha}, \omega_j^{\alpha}\}_{j=0}^N$ be the set of LG or LGR nodes and weights.

• Laguerre-Gauss (LG): $\{x_j^{\alpha}\}_{j=0}^N$ are zeros of $\mathcal{L}_{N+1}^{(\alpha)}(x)$,

$$\omega_{j}^{\alpha} = -\frac{\Gamma(N+\alpha+1)}{(N+1)!} \frac{1}{\mathcal{L}_{N}^{(\alpha)}(x_{j}^{\alpha})\partial_{x}\mathcal{L}_{N+1}^{(\alpha)}(x_{j}^{\alpha})} = \frac{\Gamma(N+\alpha+1)}{(N+\alpha+1)(N+1)!} \frac{x_{j}^{\alpha}}{\left[\mathcal{L}_{N}^{(\alpha)}(x_{j}^{\alpha})\right]^{2}}, \ 0 \leq j \leq N.$$

$$(19)$$

• Laguerre-Gauss-Radau (LGR): $x_0^{\alpha} = 0$ and $\{x_j^{\alpha}\}_{j=1}^N$ are zeros of $\partial_x \mathcal{L}_{N+1}^{(\alpha)}(x)$;

$$\omega_0^{\alpha} = \frac{(\alpha+1)\Gamma^2(\alpha+1)N!}{\Gamma(N+\alpha+2)}, \quad \omega_j^{\alpha} = \frac{\Gamma(N+\alpha+1)}{N!(N+\alpha+1)} \frac{1}{\left[\partial_x \mathcal{L}_N^{(\alpha)}(x_j^{\alpha})\right]^2} = \frac{\Gamma(N+\alpha+1)}{N!(N+\alpha+1)} \frac{1}{\left[\mathcal{L}_N^{(\alpha)}(x_j^{\alpha})\right]^2}, \quad 1 \le j < N.$$

$$(20)$$

with $\delta = 1, 0$ for LG and LGR quadrature, respectively, we have

$$\int_0^{+\infty} p(x) x^{\alpha} e^{-x} \, \mathrm{d}x = \sum_{j=0}^N p(x_j^{\alpha}) \omega_j^{\alpha}, \quad \forall \, p \in P_{2N+\delta},$$
(21)

Remark 2. Denote $\{x_j, \omega_j\}_{j=0}^N$ the usual $(\alpha = 0)$ Laguerre–Gauss type nodes and weights

• LG case:
$$\{x_j\}_{j=0}^N$$
 zeros of $\mathcal{L}_{N+1}(x)$,

$$\omega_j = \frac{x_j}{(N+1)\mathcal{L}_N^2(x_j)}, \quad 0 \le j \le N.$$
(22)

• LGR case: $x_0 = 0$, $\{x_j\}_{j=1}^N$ zeros of $\partial_x \mathcal{L}_{N+1}(x)$,

$$\omega_j = \frac{1}{(N+1)\mathcal{L}_N^2(x_j)}, \quad 0 \le j \le N.$$
(23)

We find from (9) that weights are exponentially small for large x_j .

$$\mathcal{L}_N(x) \sim e^{x/2} (Nx)^{-1/4} \Longrightarrow \omega_j \sim e^{-x_j} \sqrt{\frac{x_j}{N}}$$

Theorem 3. Let $\{x_j^{\alpha}, \omega_j^{\alpha}\}_{j=0}^N$ be the LG or LGR quadrature nodes and weights. Define

$$\hat{\omega}_{j}^{\alpha} = e^{x_{j}^{\alpha}} \omega_{j}^{\alpha}, \quad 0 \leqslant j \leqslant N,$$
(24)

and

$$\hat{P}_N := \left\{ \phi : \phi = e^{-x/2} \psi, \ \forall \psi \in P_N \right\}.$$
(25)

Then we have the modified quadrature formula

$$\int_0^{+\infty} p(x) x^{\alpha} \, \mathrm{d}x = \sum_{j=0}^N p(x_j^{\alpha}) \, \hat{\omega}_j^{\alpha}, \quad \forall \, p \in \hat{P}_{2N+\delta}, \tag{26}$$

where $\delta = 1, 0$ for the modified Laguerre-Gauss rule and the modified Laguerre-Gauss-Radau rule, respectively.

Remark 4. For $\alpha = 0$, $\hat{\omega}_j = \frac{1}{(N+1)[\hat{\mathcal{L}}_N(x_j)]^2}$, $0 \leq j \leq N \Longrightarrow \hat{\omega}_j = O(\sqrt{x_j/N})$.

Computation of Nodes and Weights

• Eigenvalue method to calculate the nodes $\{x_j^{\alpha}\}_{j=0}^N$

$$A_{N+1} = \begin{pmatrix} a_0 & -\sqrt{b_1} & & \\ -\sqrt{b_1} & a_1 & -\sqrt{b_2} & & \\ & \ddots & \ddots & \ddots & \\ & & -\sqrt{b_{N-1}} & a_{N-1} & -\sqrt{b_N} \\ & & & -\sqrt{b_N} & a_N \end{pmatrix},$$

 $a_j = 2j + \alpha + 1, \quad 0 \leq j \leq N; \quad b_j = j(j + \alpha), \ 1 \leq j \leq N.$

• First calculate weights $\{\hat{\omega}_{j}^{\alpha}\}_{j=0}^{N}$, then compute $\{\omega_{j}^{\alpha}\}_{j=0}^{N}$ by

$$\omega_j^{\alpha} = e^{-x_j^{\alpha}} \hat{\omega}_j^{\alpha}, \quad 0 \leqslant j \leqslant N.$$

• Nodes distribution: $\min_{j} \{x_{j}^{\alpha}\} \sim N^{-1}, \max_{j} \{x_{j}^{\alpha}\} \sim 4N, \min_{j} |x_{j+1}^{\alpha} - x_{j}^{\alpha}| \sim N^{-1}$







1.4 Interpolation and Discrete Laguerre Transform $\{x_j^{\alpha}, \hat{\omega}_j^{\alpha}\}_{j=0}^N$ modified LG or LGR nodes and weights, define $\langle u, v \rangle_{N, \hat{\omega}_{\alpha}} = \sum_{j=0}^N u(x_j^{\alpha})v(x_j^{\alpha})\hat{\omega}_j^{\alpha}, \quad \|u\|_{N, \hat{\omega}_{\alpha}} = \sqrt{\langle u, u \rangle_{N, \hat{\omega}_{\alpha}}}.$

The exactness of the Gauss quadrature implies

$$\langle p,q \rangle_{N,\hat{\omega}_{\alpha}} = (p,q)_{\hat{\omega}_{\alpha}}, \quad \forall p \cdot q \in \hat{P}_{2N+\delta},$$
(27)

where $\delta = 1, 0$ for the modified LG and LGR quadrature, respectively. Define interpolation $\hat{I}_N^{(\alpha)}: C[0, +\infty) \to \hat{P}_N$ as

$$\left(\hat{I}_N^{(\alpha)}u\right)(x) = \sum_{n=0}^N \tilde{u}_n^\alpha \hat{\mathcal{L}}_n^{(\alpha)}(x).$$

sucht that

$$\left(\hat{I}_{N}^{(\alpha)}u\right)(x_{j}^{\alpha}) = u(x_{j}^{\alpha}), \quad 0 \leq j \leq N.$$

Given the physical values $\{u(x_j^{\alpha})\}_{j=0}^N$, the coefficients $\{\tilde{u}_n^{\alpha}\}_{n=0}^N$ can be determined by the forward discrete transform:

$$\tilde{u}_n^{\alpha} = \frac{1}{\gamma_n^{(\alpha)}} \sum_{j=0}^N u(x_j^{\alpha}) \hat{\mathcal{L}}_n^{(\alpha)}(x_j^{\alpha}) \hat{\omega}_j^{\alpha}, \quad 0 \le n \le N.$$

Backward discrete transform calculates $\{u(x_j^{\alpha})\}_{j=0}^N$ using $\{\tilde{u}_n^{\alpha}\}_{n=0}$:

$$u(x_j^{\alpha}) = \sum_{n=0}^{N} \tilde{u}_n^{\alpha} \hat{\mathcal{L}}_n^{(\alpha)}(x_j^{\alpha}), \quad 0 \leq j \leq N.$$

The above definitions extends to usual Laguerre-Gauss transform directly (by removing "^").

1.5 Differentiation

Differentiation in Physical Space

• Differentiation in Laguerre polynomials space P_N : $\{h_j\}$ Lagrange polynomials assoc. with LGR points $\{x_j^{\alpha}\}_{j=0}^N$. $u \in P_N, \ \boldsymbol{u}^{(m)} = (u^{(m)}(x_0^{\alpha}), ..., u^{(m)}(x_N^{\alpha}))^T, \ \boldsymbol{u} := \boldsymbol{u}^{(0)}, \text{ then}$ $\boldsymbol{u}^{(m)} = D^m \boldsymbol{u}, \quad m \ge 1,$ (28)

where $D = (d_{kj})_{k,j=0,...,N}, d_{kj} = h'_j(x_k).$

• Differentiation in Laguerre function space \hat{P}_N :

$$\hat{h}_j(x) = e^{-x/2} / e^{-x_j^{\alpha}/2} h_j(x).$$
(29)

then

$$\hat{h}_{j}(x_{k}^{\alpha}) = \delta_{kj}, \quad 0 \leqslant k, j \leqslant N; \quad \hat{P}_{N} = \operatorname{span} \left\{ \hat{h}_{j} : 0 \leqslant j \leqslant N \right\}.$$

Moreover $\hat{h}_{j}' \in \hat{P}_{N}$ and

$$\hat{d}_{kj} := \hat{h}'_j(x_k^{\alpha}) = e^{-x_k^{\alpha}/2} / e^{-x_j^{\alpha}/2} d_{kj} - \frac{1}{2} \delta_{kj}, \quad 0 \le j, k \le N.$$
(30)

Differentiation in spectral space

$$u(x) = \sum_{n=0}^{N} \hat{u}_n \mathcal{L}_n^{(\alpha)}(x), \quad \hat{u}_n = \frac{1}{\gamma_n^{(\alpha)}} \int_0^{+\infty} u(x) \mathcal{L}_n^{(a)}(x) \omega_\alpha(x) \, \mathrm{d}x,$$
$$u'(x) = \sum_{n=1}^{N} \hat{u}_n \partial_x \mathcal{L}_n^{(\alpha)}(x) = \sum_{n=0}^{N} \hat{u}_n^{(1)} \mathcal{L}_n^{(\alpha)}(x) \in P_{N-1}, \quad \hat{u}_N^{(1)} = 0.$$

The coefficients $\{\hat{u}_n^{(1)}\}_{n=0}^{N-1}$ can be calculated from $\{\hat{u}_n\}_{n=1}^N$ by

$$\begin{cases} \hat{u}_{n-1}^{(1)} = \hat{u}_n^{(1)} - \hat{u}_n, & n = N, N - 1, \dots, 1, \\ \hat{u}_N^{(1)} = 0. \end{cases}$$
(31)

Similarly, for the Laguerre function case (in space \hat{P}_N)

$$\begin{cases} \hat{v}_{n}^{(1)} = \hat{v}_{n+1}^{(1)} - \frac{1}{2}(\hat{v}_{n} + \hat{v}_{n+1}), & n = N - 1, \dots, 0, \\ \hat{v}_{N}^{(1)} = -\frac{1}{2}\hat{v}_{N}. \end{cases}$$
(32)

2 Hermite Polynomials/Functions

2.1 Hermite polynomials

1. $H_m(x)$ defined on $\mathbb{R} := (-\infty, +\infty)$, weight $\omega(x) = e^{-x^2}$

$$\int_{-\infty}^{\infty} H_m(x)H_n(x)\omega(x)\,\mathrm{d}x = \gamma_n\,\delta_{m\,n}, \quad \gamma_n = \sqrt{\pi}2^n n!. \tag{33}$$

2. Three-term recurrence relation

$$H_{n+1}(x) = 2 x H_n(x) - 2n H_{n-1}(x), \quad n \ge 1, \qquad (34)$$
$$H_0(x) = 1, \quad H_1(x) = 2x.$$

Leading coefficient $k_n = 2^n$.

3. Connection with Laguerre polynomials

$$H_{2n}(x) = (-1)^n 2^{2n} n! \mathcal{L}_n^{(-1/2)}(x^2),$$

$$H_{2n+1}(x) = (-1)^n 2^{2n+1} n! x \mathcal{L}_n^{(1/2)}(x^2).$$
(35)

4. Values at 0:

$$H_{2n}(0) = (-2)^n (2n-1)!!, \quad H_{2n+1}(0) = 0.$$
(36)

5. Sturm–Liouville equation:

$$e^{x^{2}} \left(e^{-x^{2}} H_{n}'(x) \right)' + \lambda_{n} H_{n}(x) = 0, \quad \lambda_{n} = 2n, \quad (37)$$

$$H_n''(x) - 2x H_n'(x) + \lambda_n H_n(x) = 0.$$
(38)

6. Rodrigues' formula

$$H_n(x) = (-1)^n e^{x^2} \frac{\mathrm{d}^n}{\mathrm{d}x^n} \{e^{-x^2}\},\tag{39}$$

explicit expression

$$H_n(x) = \sum_{k=0}^{[n/2]} \frac{(-1)^k n!}{k! (n-2k)!} (2x)^{n-2k}.$$
 (40)

7. Upper bound

$$|H_n(x)| < c \, 2^{n/2} \sqrt{n!} e^{x^2/2}, \quad c \approx 1.086435. \tag{41}$$

8. Derivatives

$$H'_n(x) = \lambda_n H_{n-1}(x), \quad n \ge 1.$$
(42)

$$H'_{n}(x) = 2x H_{n}(x) - H_{n+1}(x), \quad n \ge 0.$$
(43)

9. Asymptotics

$$\frac{\Gamma(n/2+1)}{n!}e^{-x^2/2}H_n(x) = \cos\left(\sqrt{2n+1}x - \frac{n\pi}{2}\right) + \frac{x^3}{6\sqrt{2n+1}}\sin\left(\sqrt{2n+1}x - \frac{n\pi}{2}\right) + O(n^{-1}).$$
(44)

2.2 Hermite functions



$$\hat{H}_n(x) = \frac{1}{\pi^{1/4}\sqrt{2^n n!}} e^{-x^2/2} H_n(x), \quad n \ge 0, \ x \in \mathbb{R}$$
(45)

• Orthogonality

$$\int_{\mathbb{R}} \hat{H}_n(x) \hat{H}_m(x) \,\mathrm{d}x = \delta_{m\,n}.$$
(46)

• Three-term recurrence

$$\hat{H}_{n+1}(x) = x \sqrt{\frac{2}{n+1}} \hat{H}_n(x) - \sqrt{\frac{n}{n+1}} \hat{H}_{n-1}(x), n \ge 1$$

$$\hat{H}_0 = \pi^{-1/4} e^{-x^2/2}, \quad \hat{H}_1(x) = \sqrt{2}\pi^{-1/4} x \, e^{-x^2/2}.$$
(47)

• Differentiation equation

$$\hat{H}_{n}^{\prime\prime}(x) + (2n+1-x^{2})\hat{H}_{n}(x) = 0.$$
(48)

• Derivatives

$$\hat{H}'_{n}(x) = \sqrt{2n} \hat{H}_{n-1}(x) - x \hat{H}_{n}(x) = \sqrt{\frac{n}{2}} \hat{H}_{n-1}(x) - \sqrt{\frac{n+1}{2}} \hat{H}_{n+1}(x)$$
(49)

leads to

$$\int_{\mathbb{R}} \hat{H}'_{m}(x) \hat{H}'_{n}(x) \, \mathrm{d}x = \begin{cases} -\frac{\sqrt{n(n-1)}}{2}, & m=n-2, \\ n+\frac{1}{2}, & m=n, \\ -\frac{\sqrt{(n+1)(n+2)}}{2}, & m=n+2, \\ 0, & \text{otherwise.} \end{cases}$$
(50)

• Asymptotic $(\hat{H}_n(x) \sim n^{-1/4}, n \to \infty)$

$$\hat{H}_n(x) = \pi^{-1/4} \sqrt{2^n n!} \sum_{k=0}^{\lfloor n/2 \rfloor} \frac{(-1)^k}{4^k k! (n-2k)!} \left(x^{n-2k} e^{-x^2/2} \right).$$
(51)

2.3 Hermite-Gauss Quadrature

Theorem 5. (Hermite-Gauss Quadrature) Let $\{x_j\}_{j=0}^N$ be zeros of $H_{N+1}(x)$, and let $\{\omega_j\}_{j=0}^N$ be given by

$$\omega_j = \frac{\sqrt{\pi} 2^N N!}{(N+1) H_N^2(x_j)}, \quad 0 \le j \le N.$$
(52)

Then,

$$\int_{-\infty}^{+\infty} p(x)e^{-x^2} dx = \sum_{j=0}^{N} p(x_j)\omega_j, \quad \forall p \in P_{2N+1}.$$
 (53)

Theorem 6. (Modified Hermite-Gauss Quadrature) Let $\{x_j, \omega_j\}_{j=0}^N$ be the Hermite-Gauss quadrature nodes and weights, Define

$$\hat{\omega}_{j} = e^{x_{j}^{2}} \omega_{j} = \frac{1}{(N+1)\hat{H}_{N}^{2}(x_{j})}, \quad 0 \leq j \leq N.$$
(54)

Then

$$\int_{-\infty}^{+\infty} p(x)q(x) \, \mathrm{d}x = \sum_{j=0}^{N} p(x_j)q(x_j)\hat{\omega}_j, \ \forall p \cdot q \in \hat{P}_{2N+1}, \tag{55}$$

where

$$\hat{P}_M := \left\{ \phi : \phi = e^{-x^2/2} \psi, \ \forall \psi \in P_M \right\}.$$
(56)

Compution of Nodes and Weights

• Zeros $\{x_j\}_{j=0}^N$ of $H_{n+1}(x)$ are eigenvalues of

$$A_{N+1} = \begin{pmatrix} a_0 & \sqrt{b_1} & & \\ \sqrt{b_1} & a_1 & \sqrt{b_2} & & \\ & \ddots & \ddots & \ddots & \\ & & \sqrt{b_{N-1}} & a_{N-1} & \sqrt{b_N} \\ & & & \sqrt{b_N} & a_N \end{pmatrix},$$
(57)

$$a_j = 0, \ 0 \leqslant j \leqslant N; \quad b_j = j/2, \ 1 \leqslant j \leqslant N.$$

• Computing $\{\omega_j\}_{j=0}^N$ using (52) is not stable. Instead using

$$\omega_j = e^{-x_j^2} \hat{\omega}_j,$$

 $\hat{\omega}_j$ are given by (54).

• Node distribution: $\max_j |x_j| \sim \sqrt{2N},$

 $\max_{j} |x_{j}| \sim \sqrt{2N}, \qquad \min_{j} |x_{j} - x_{j-1}| \sim 1/\sqrt{N}$



Fig. 7.4 (a) Distribution of the Hermite-Gauss nodes $\{x_j\}_{j=0}^N$ with N = 8, 16, 24, 32; (b) Growth of the largest node against the asymptotic estimate: $\sqrt{2(N+1) - (2(N+1))^{1/3}}$ (*dashed line*) with various N

2.4 Interpolation and Discrete Hermite Transforms

• Using Hermite polynomials in P_N : $I_N^h v \in P_N$ intepolates $v \in C(\mathbb{R})$ at Hermite-Gauss pts $\{x_j\}_{j=0}^N$:

$$I_{N}^{h}v = \sum_{n=0}^{N} \tilde{v}_{n}H_{n}(x); \quad (I_{N}^{h}v)(x_{j}) = v(x_{j}), \quad 0 \leq j \leq N.$$

Forward discrete transform:

$$\tilde{v}_n = \frac{1}{\gamma_n} \sum_{j=0}^N v(x_j) H_n(x_j) \omega_j, \quad 0 \leqslant n \leqslant N$$
(58)

Backward discrete transform

$$v(x_j) = \sum_{n=0}^{N} \tilde{v}_n H_n(x_j), \quad 0 \leqslant j \leqslant N.$$
(59)

• Using Hermite functions in \hat{P}_N $\hat{I}_N^h u \in \hat{P}_N$ intepolates $u \in C(\mathbb{R})$ at Hermite-Gauss pts $\{x_j\}_{j=0}^N$:

$$\hat{I}_N^h u = \sum_{n=0}^N \tilde{u}_n \hat{H}_n(x); \quad \left(\hat{I}_N^h u\right)(x_j) = u(x_j), \quad 0 \leq j \leq N.$$

Forward discrete transform:

$$\tilde{u}_n = \frac{1}{\gamma_n} \sum_{j=0}^N u(x_j) \hat{H}_n(x_j) \hat{\omega}_j, \quad 0 \leqslant n \leqslant N$$
(60)

Backward discrete transform

$$u(x_j) = \sum_{n=0}^{N} \tilde{u}_n \hat{H}_n(x_j), \quad 0 \leqslant j \leqslant N.$$
(61)

2.5 Differentiation

Differentiation in Physical Space

• Hermite polynomials case (space P_N):

$$\boldsymbol{u}^{(m)} = D^m \boldsymbol{u}, \quad m \ge 1, \tag{62}$$

where the entries of D are given by

$$d_{kj} = \begin{cases} \frac{H_N(x_k)}{H_N(x_j)} \frac{1}{x_k - x_j}, & k \neq j, \\ x_k, & k = j. \end{cases}$$
(63)

• Hermite functions case (space \hat{P}_N):

$$\hat{d}_{kj} = -x_k \delta_{kj} + \frac{e^{-x_k^2/2}}{e^{-x_j^2/2}} d_{kj} = \begin{cases} \frac{\hat{H}_N(x_k)}{\hat{H}_N(x_j)} \frac{1}{x_k - x_j}, & k \neq j, \\ 0, & k = j. \end{cases}$$
(64)

Differentiation in Frequency Space

For $u \in P_N$, $u' \in P_{N-1}$:

$$u(x) = \sum_{n=0}^{N} \hat{u}_n H_n(x), \quad u'(x) = \sum_{n=1}^{N} \hat{u}_n H'_n(x) = \sum_{n=0}^{N} \hat{u}_n^{(1)} H_n(x),$$
$$\hat{u}_N^{(1)} = 0; \quad \hat{u}_n^{(1)} = 2(n+1)\hat{u}_{n+1}, \ n = N-1, N-2, \dots, 0.$$
(65)
For $v \in \hat{P}_N, \ v' \in \hat{P}_{N+1}$:

$$v(x) = \sum_{n=0}^{N} \hat{v}_n \hat{H}_n(x), \quad v'(x) = \sum_{n=0}^{N} \hat{v}_n \hat{H}'_n(x) = \sum_{n=0}^{N+1} \hat{v}_n^{(1)} \hat{H}_n(x).$$

$$\hat{v}_{n}^{(1)} = \sqrt{\frac{n+1}{2}} \hat{v}_{n+1} - \sqrt{\frac{n}{2}} \hat{v}_{n-1}, \quad n = N+1, N, \dots, 0.$$
(66)

with $\hat{v}_{-1} = \hat{v}_{N+1} = \hat{v}_{N+2} = 0.$

3 Approximation Estimates

3.1 Inverse Inequalities

For Laguerre approximation, $\omega_{\alpha} = x^{\alpha} e^{-x}$ and $\hat{\omega}_{\alpha} = x^{\alpha}$.

Theorem 7. For $\alpha > -1$ and any $\phi \in P_N$, $\|\partial_x^m \phi\|_{\omega_{\alpha+m}} \lesssim N^{m/2} \|\phi\|_{\omega_{\alpha}}, \quad m \ge 0.$ (67) **Corollary 8.** For $\alpha > -1$ and any $\psi \in \hat{P}_N$,

$$\left\|\hat{\partial}_{x}^{m}\psi\right\|_{\hat{\omega}_{\alpha+m}} \leqslant N^{m/2} \|\psi\|_{\hat{\omega}_{\alpha}}, \quad m \ge 0.$$
(68)

Theorem 9. For $\alpha \ge 0$ and any $\phi \in P_N$,

$$\|\partial_x^m \phi\|_{\omega_\alpha} \leqslant N^m \|\phi\|_{\omega_\alpha}, \quad m \ge 0.$$
(69)

Corollary 10. For $\alpha \ge 0$ and any $\psi \in \hat{P}_N$, $\|\hat{\partial}_x^m \psi\|_{\hat{\omega}_{\alpha}} \lesssim N^m \|\psi\|_{\hat{\omega}_{\alpha}}, \quad m \ge 0.$ (70)

For the Hermite case, only one weight $\omega(x) = e^{-x^2}$. We have

Theorem 11. For any $\phi \in P_N$, $\|\partial_x \phi\|_{\omega} \lesssim N \|\phi\|_{\omega}$, moreover, let $\hat{\partial} = \partial_x + x$. Then for any $\psi \in \hat{P}_N$, $\|\hat{\partial}_x \psi\| \lesssim N \|\psi\|$.
3.2 Orthogonal Projections

Laguerre polynomial $L^2_{\omega_{\alpha}}$ projection: $\Pi_{N,\alpha}: L^2_{\omega_{\alpha}}(\mathbb{R}_+) \to P_N$

$$(\Pi_{N,\alpha}u - u, v_N)_{\omega_{\alpha}} = 0, \quad \forall v_N \in P_N$$
(71)

we have

$$\Pi_{N,\alpha} u = \sum_{n=0}^{N} \hat{u}_n^{(\alpha)} \mathcal{L}_n^{(\alpha)}(x), \quad \hat{u}_n^{(\alpha)} = \frac{1}{\gamma_n^{(\alpha)}} \left(u, \mathcal{L}_n^{(\alpha)} \right)_{\omega_\alpha}.$$

Define

$$B^m_{\alpha}(\mathbb{R}_+) := \{ u : \partial^k_x u \in L^2_{\omega_{\alpha+k}}(\mathbb{R}_+), \ 0 \leqslant k \leqslant m \},$$
(72)

 $\|u\|_{B^m_{\alpha}} = \left(\sum_{k=0}^m \|\partial_x^k u\|_{\omega_{\alpha+k}}^2\right)^{1/2}, \quad \|u\|_{B^m_{\alpha}} = \|\partial_x^m u\|_{\omega_{\alpha+m}}.$

Theorem 12. Let $\alpha > -1$. If $u \in B^m_{\alpha}(\mathbb{R}_+)$ and $0 \leq m \leq N+1$, then

$$\|\partial_{x}^{l}(\Pi_{N,\alpha}u - u)\|_{\omega_{\alpha+l}} \leqslant \sqrt{\frac{(N - m + 1)!}{(N - l + 1)!}} \|\partial_{x}^{m}u\|_{\omega_{\alpha+m}}.$$
 (73)

Laguerre function case: $\hat{\omega}_{\alpha} = x^{\alpha}$ For $u \in L^2_{\hat{\omega}_a}(\mathbb{R}_+)$, we have $u e^{x/2} \in L^2_{\omega_{\alpha}}(\mathbb{R}_+)$. Define $\hat{\Pi}_{N,\alpha} u = e^{-x/2} \Pi_{N,\alpha} (u e^{x/2}). \in \hat{P}_N$ (74)

Clearly

$$\left(\hat{\Pi}_{N,\alpha}u-u,v_N\right)_{\hat{\omega}_{\alpha}}=0,\quad\forall v_N\in\hat{P}_N.$$

Define

$$\hat{B}^{m}_{\alpha}(\mathbb{R}_{+}) := \left\{ u : \hat{\partial}^{k}_{x} u \in L^{2}_{\hat{\omega}_{\alpha+k}}(\mathbb{R}_{+}), \ 0 \leqslant k \leqslant m \right\},$$
(75)

$$\|u\|_{\hat{B}^{m}_{\alpha}} = \left(\sum_{k=0}^{m} \|\hat{\partial}^{k}_{x}u\|_{\hat{\omega}_{\alpha+k}}^{2}\right)^{1/2}, \quad |u|_{\hat{B}^{m}_{\alpha}} = \|\hat{\partial}^{m}_{x}u\|_{\hat{\omega}_{\alpha+m}}.$$

Theorem 13. Let $\alpha > -1$. If $u \in \hat{B}^m_{\alpha}(\mathbb{R}_+)$ and $0 \leq m \leq N+1$, then

$$\left\|\hat{\partial}_{x}^{l}(\Pi_{N,\alpha}u-u)\right\|_{\hat{\omega}_{\alpha+l}} \leqslant \sqrt{\frac{(N-m+1)!}{(N-l+1)!}} \left\|\hat{\partial}_{x}^{m}u\right\|_{\hat{\omega}_{\alpha+m}}.$$
(76)

Remark 14. Remarks on Laguerre approximation:

- 1. The convergence rate is about $N^{(l-m)/2}$, is only half of the classical Jacobi approximation. This is a direct consequence of the linear growth of the eigenvalues in the Sturm-Liouville problem.
- 2. $B^m_{\alpha}(\mathbb{R}_+)$ includes functions that do not decay at infinity. A fast convergence rate in $\|\cdot\|_{\omega_{\alpha+l}}$ norm does not mean that the error would decay rapidly for large x.
- 3. $u \in \hat{B}^{m}_{\alpha}(\mathbb{R}_{+})$ requires u decays at infinity. So theorem (13) doesn't apply to functions like $\sin(x)$. On the other hand, for $u(x) = (1+x)^{-h}$ and $u(x) = \sin(kx)/(1+x)^{h}$, we have for both $\|\hat{\partial}^{m}_{x}u\|_{\hat{\omega}_{\alpha+m}} < \infty$ if $m < 2h \alpha 1$, so we have

$$\left\| u - \hat{\Pi}_{N,\alpha} u \right\|_{\hat{\omega}a} \lesssim N^{-(2h-\alpha-1)/2}.$$
(77)

*H*¹-type projections. (Here we only consider the case $\alpha = 0$) Let $\omega(x) = e^{-x}$. Denote

 $H^{1}_{0,\omega}(\mathbb{R}_{+}) = \{ u \in H^{1}_{\omega}(\mathbb{R}_{+}) : u(0) = 0 \}, \quad P^{0}_{N} = \{ \phi \in P_{N} : \phi(0) = 0 \}.$ Define $\Pi^{1,0}_{N} : H^{1}_{0,\omega}(\mathbb{R}_{+}) \to P^{0}_{N}$ as

$$\left(\left(u - \Pi_N^{1,0} u\right)', v_N'\right)_{\omega} = 0, \quad \forall v_N \in P_N^0.$$
 (78)

Theorem 15. If $u \in H^1_{0,\omega}(\mathbb{R}_+)$ and $\partial_x u \in B^{m-1}_0(\mathbb{R}_+)$, then for $1 \leq m \leq N+1$,

$$\left\| \Pi_{N,\alpha}^{1,0} u - u \right\|_{1,\omega} \lesssim \sqrt{\frac{(N-m+1)!}{N!}} \|\partial_x^m u\|_{\omega_{m-1}},\tag{79}$$

Proof: Let $\phi(x) = \int_0^x \Pi_{N-1,0} u'(y) \, \mathrm{d}y$. Then $u - \phi \in H^1_{0,\omega}(\mathbb{R}_+)$. (B.35b)

$$\|\Pi_{N}^{1,0}u - u\|_{1,\omega} \leq \|\phi - u\|_{1,\omega} \leq c \|\partial_{x}(\phi - u)\|_{\omega}$$

Laguerre function case For $u \in H_0^1(\mathbb{R}_+)$, we have $u e^{x/2} \in H_{0,\omega}^1(\mathbb{R}_+)$. Define

$$\hat{\Pi}_N^{1,0} u = e^{-x/2} \Pi_N^{1,0} (u \, e^{x/2}) \quad \in \hat{P}_N^0.$$

Theorem 16. For any $u \in H_0^1(\mathbb{R}_+)$, we have

$$\left(\left(u - \hat{\Pi}_{N}^{1,0} u \right)', v_{N}' \right) + \frac{1}{4} \left(u - \hat{\Pi}_{N}^{1,0} u, v_{N} \right) = 0, \quad \forall v_{N} \in \hat{P}_{N}^{0}.$$

$$Let \ \hat{\partial_{x}} = \partial_{x} + \frac{1}{2}. \ If \ u \in H_{0}^{1}(\mathbb{R}_{+}) \ and \ \hat{\partial_{x}} u \in \hat{B}_{n}^{m-1}(\mathbb{R}_{+}), \ then$$

$$\left\| \hat{\Pi}_{N}^{1,0} u - u \right\|_{1} \leqslant c \sqrt{\frac{(N-m+1)!}{N!}} \left\| \hat{\partial_{x}}^{m} u \right\|_{\hat{\omega}_{m-1}},$$

$$(81)$$

where c is a positive constant independent of m, N and u.

H^2 -type orthogonal projections

Define

$$H_{0,\omega}^{2}(\mathbb{R}_{+}) = \{ v \in H_{\omega}^{2}(\mathbb{R}_{+}) : v(0) = v'(0) = 0 \}, \quad X_{N} = H_{0,\omega}^{2}(\mathbb{R}_{+}) \cap P_{N},$$

$$H_0^2(\mathbb{R}_+) = \{ v \in H^2(\mathbb{R}_+) : v(0) = v'(0) = 0 \}, \quad \hat{X}_N = H_0^2(\mathbb{R}_+) \cap \hat{P}_N.$$

Define $\Pi^{2,0}_N: H^2_{0,\omega}(\mathbb{R}_+) \to X_N$, as

$$\left(\left(v - \Pi_N^{2,0} u\right)'', v_N''\right)_{\omega} = 0, \quad \forall v_N \in X_N.$$
 (82)

Define $\hat{\Pi}_N^{2,0}: H_0^2(\mathbb{R}_+) \to \hat{X}_N$ as

$$\hat{\Pi}_{N}^{2,0} u = e^{-x/2} \Pi_{N}^{2,0} (u \, e^{x/2}). \tag{83}$$

Theorem 17. If $v \in H^2_{0,\omega}(\mathbb{R}_+)$ and $\partial_x^2 v \in B^{m-2}_0(\mathbb{R}_+)$ with $2 \leq m \leq N+1$, then we have

$$\left\| \Pi_{N}^{2,0} v - v \right\|_{2,\omega} \leq c_{\sqrt{\frac{(N-m+1)!}{(N-1)!}}} \|\partial_{x}^{m} v\|_{\omega_{m-2}}.$$
(84)

For $u \in H_0^2(\mathbb{R}_+)$ and all $u_N \in \hat{X}_N$, we have $((u - \hat{\Pi}_N^{2,0}u)'', u_N'') + \frac{1}{2}((u - \hat{\Pi}_N^{2,0}u)', u_N') + \frac{1}{16}(u - \hat{\Pi}_N^{2,0}u, u_N) = 0.$ (85) Moreover, if $u \in H_0^2(\mathbb{R}_+)$ and $\hat{\partial}_x^2 u \in \hat{B}_0^{m-2}(\mathbb{R}_+)$ with $2 \leq m \leq N+1$,

Moreover, if $u \in H_0^-(\mathbb{R}_+)$ and $\mathcal{O}_x u \in \mathcal{B}_0$ (\mathbb{R}_+) with $2 \leq m \leq N+1$, then

$$\left\|\hat{\Pi}_{N}^{2,0}u - u\right\|_{2,\hat{\omega}} \leqslant c_{\sqrt{\frac{(N-m+1)!}{(N-1)!}}} \left\|\hat{\partial}_{x}^{m}u\right\|_{\hat{\omega}_{m-2}}.$$
(86)

Hermite projections $(\omega(x) = e^{-x^2})$ $\Pi_N: L^2_{\omega}(\mathbb{R}) \to P_N:$

$$(u - \Pi_N u, v_N)_\omega = 0, \quad \forall v_N \in P_{N}.$$
(87)

Clearly

$$\Pi_N u(x) = \sum_{n=0}^N \hat{u}_n H_n(x), \quad \hat{u}_n = \frac{1}{\gamma_n} (u, H_n)_{\omega}.$$

Theorem 18. For any $u \in H^m_{\omega}(\mathbb{R})$ with $0 \leq m \leq N+1$,

$$\|\partial_x^l (\Pi_N u - u)\|_{\omega} \leq 2^{(l-m)/2} \sqrt{\frac{(N-m+1)!}{(N-l+1)!}} \|\partial_x^m\|_{\omega}, \quad 0 \leq l \leq m.$$
(88)

<u>Reamrks</u>: The $L^2_{\omega}(\mathbb{R})$ -orthogonal projection is simultaneously optimal in the $H^l_{\omega}(\mathbb{R})$ -norm with $l \ge 1$.

Hermite function approximation

For $u \in L^2(\mathbb{R})$, we have $u e^{x^2/2} \in L^2_{\omega}(\mathbb{R})$, define

$$\hat{\Pi}_N u := e^{-x^2/2} \Pi_N \left(u \, e^{x^2/2} \right) \quad \in \hat{P}_N, \tag{89}$$

which satisfies

$$(u - \hat{\Pi}_N u, v_N) = (u e^{x^2/2} - \Pi_N (u e^{x^2/2}), v_N e^{x^2/2})_\omega = 0, \forall v_N \in \hat{P}_N.$$

Theorem 19. Let $\hat{\partial}_x = \partial_x + x$. For $\hat{\partial}_x^m u \in L^2(\mathbb{R})$ with $0 \leq m \leq N+1$, $\left\| \hat{\partial}_x^l (\hat{\Pi}_N u - u) \right\| \lesssim 2^{(l-m)/2} \sqrt{\frac{(N-m+1)!}{(N-l+1)!}} \left\| \hat{\partial}_x^m u \right\|, \quad 0 \leq l \leq m.$ (90)

Theorem 20. For any $\hat{\partial}_x^m u \in L^2(\mathbb{R})$ with $2 \leq m \leq N+1$,

$$\left\|\partial_{x}^{l}(\hat{\Pi}_{N}u-u)\right\| \lesssim \sqrt{\frac{(N-m+1)!}{2^{m}(N-l+1)!}} \left\|\hat{\partial}_{x}^{m}u\right\|, \quad l=0,1,2.$$
(91)

3.3 Interpolations

Theorem 21. (Laguerre-Gauss interpolation[Guo et al 2006b]) Let $\alpha > -1$. If $u \in C(\mathbb{R}_+) \cap B^m_{\alpha}(\mathbb{R}_+)$ and $\partial_x u \in B^{m-1}_{\alpha}(\mathbb{R}_+)$ with $1 \leq m \leq N+1$, then

$$\left\|I_N^{(\alpha)}u - u\right\|_{\omega_{\alpha}} \lesssim \sqrt{\frac{(N-m+1)!}{N!}} \left(\|\partial_x^m u\|_{\omega_{\alpha+m-1}} + \sqrt{\ln N} \|\partial_x^m u\|_{\omega_{\alpha+m}}\right).$$

Theorem 22. Let $\alpha > -1$. If $u \in C(\mathbb{R}_+) \cap \hat{B}^m_{\alpha}(\mathbb{R}_+)$ and $\hat{\partial}_x u \in \hat{B}^{m-1}_{\alpha}(\mathbb{R}_+)$ with $1 \leq m \leq N+1$ then

$$\left\| \hat{I}_{N}^{(\alpha)}u - u \right\|_{\hat{\omega}\alpha} \lesssim \sqrt{\frac{(N-m+1)!}{N!}} \Big(\left\| \hat{\partial}_{x}^{m}u \right\|_{\hat{\omega}_{\alpha+m-1}} + \sqrt{\ln N} \left\| \hat{\partial}_{x}^{m}u \right\|_{\hat{\omega}_{\alpha+m}} \Big)$$

Hermite Interpolation

Theorem 23. ([Guo and Xu 2000])For $u \in C(\mathbb{R}) \cap H^m_{\omega}(\mathbb{R})$ with $m \ge 1$, we have

$$\|\partial_x^l (I_N^h u - u)\|_{\omega} \lesssim N^{\frac{1}{6} + \frac{l - m}{2}} \|\partial_x^m u\|_{\omega}, \quad 0 \leqslant l \leqslant m.$$

$$(92)$$

Theorem 24. Let $\hat{\partial}_x = \partial_x + x$. For $u \in C(\mathbb{R})$ and $\hat{\partial}_x^m u \in L^2(\mathbb{R})$ with fixed $m \ge 1$, we have

$$\left\|\hat{\partial}_x^l (\hat{I}_N^h u - u)\right\| \lesssim N^{\frac{1}{6} + \frac{l - m}{2}} \left\|\hat{\partial}_x^m u\right\|, \quad 0 \leqslant l \leqslant m.$$

$$(93)$$

4 Spectral Methods Using Laguerre and Hermite Functions 4.1 Laguerre-Galerkin Method

$$-u_{xx} + \gamma u = f, \quad x \in \mathbb{R}_+, \, \gamma > 0; \\ u(0) = 0, \quad \lim_{x \to +\infty} u(x) = 0.$$
(94)

Weak formulation (for $f \in (H_0^1(\mathbb{R}_+))'$)

Find
$$u \in H_0^1(\mathbb{R}_+)$$
 such that
 $a(u,v) := (u',v') + \gamma(u,v) = (f,v), \quad \forall v \in H_0^1(\mathbb{R}_+)$

$$(95)$$

Laguerre spectral-Galerkin approximation to (94)

$$\begin{cases} \operatorname{Find} u_N \in \hat{P}_N^0 \operatorname{such} \operatorname{that} \\ a(u_N, v_N) = (\hat{I}_N f, v_N), \quad \forall v_N \in \hat{P}_N^0 \end{cases}$$
(96)

Define

$$\hat{\phi}_k(x) = \left(\mathcal{L}_k(x) - \mathcal{L}_{k+1}(x)\right) e^{-x/2} = \hat{\mathcal{L}}_k(x) - \hat{\mathcal{L}}_{k+1}(x) \tag{97}$$

then $\hat{P}_N^0 = \operatorname{span} \{ \hat{\phi}_0, \hat{\phi}_1, \dots, \hat{\phi}_{N-1} \}$. By setting

$$u_{N} = \sum_{k=0}^{N-1} \hat{u}_{k} \hat{\phi}_{k}, \quad \boldsymbol{u} = (\hat{u}_{0}, \hat{u}_{1}, ..., \hat{u}_{N-1})^{T};$$

$$f_{j} = (\hat{I}_{N}f, \hat{\phi}_{j}), \qquad \boldsymbol{f} = (f_{0}, f_{1}, ..., f_{N-1})^{T};$$

$$s_{jk} = (\hat{\phi}'_{k}, \hat{\phi}'_{j}), \qquad \boldsymbol{S} = (s_{jk})_{0 \leq j,k \leq N-1};$$

$$m_{jk} = (\hat{\phi}_{k}, \hat{\phi}_{j}), \qquad \boldsymbol{M} = (m_{jk})_{0 \leq j,k \leq N-1};$$

we find M is a symmetric tridiagonal matrix and $S = I - \frac{1}{4}M$, so (96) reduces to

$$(I+(\gamma-\frac{1}{4})\boldsymbol{M})\boldsymbol{u}=\boldsymbol{f}.$$
 (98)

Theorem 25. If $u \in H_0^1(\mathbb{R}_+)$, $\hat{\partial}_x u \in \hat{B}_0^{m-1}(\mathbb{R}_+)$, $f \in C(\overline{\mathbb{R}}_+) \cap \hat{B}_0^k(\mathbb{R}_+)$ and $\hat{\partial}_x f \in \hat{B}_0^{k-1}(\mathbb{R}_+)$ with $1 \leq k, m \leq N+1$, then we have

$$\|u - u_N\|_1 \lesssim \sqrt{\frac{(N - m + 1)!}{N!}} \|\hat{\partial}_x^m u\|_{\hat{\omega}_{m-1}} + \sqrt{\frac{(N - k + 1)!}{N!}} \Big(\|\hat{\partial}_x^k f\|_{\hat{\omega}_{k-1}} + \sqrt{\ln N} \|\hat{\partial}_x^k f\|_{\hat{\omega}_k} \Big).$$

$$(99)$$

Proof. Let $e_N = u_N - \hat{\Pi}_N^{1,0} u$ and $\tilde{e}_N = u - \hat{\Pi}_N^{1,0} u$, then by

$$a(u_N - u, v_N) = (\hat{I}_N f - f, v_N), \quad \forall v_N \in \hat{P}_N^0$$
$$a(e_N, v_N) = a(\tilde{e}_N, v_N) + (\hat{I}_N f - f, v_N), \quad \forall v_N \in \hat{P}_N^0.$$

Taking $v_N = e_N$, we find

$$||e_N||_1 \leq c(||\tilde{e}_N||_1 + ||\hat{I}_N f - f||).$$

4.2 Hermite–Galerkin Method

$$-u_{xx} + \gamma u = f, \quad x \in \mathbb{R}, \ \gamma > 0; \quad \lim_{|x| \to \infty} u(x) = 0.$$
 (100)

Weak formulation

$$\begin{cases} \operatorname{Find} u \in H^1(\mathbb{R}) \text{ such that} \\ (\partial_x u, \partial_x v) + \gamma(u, v) = (f, v), \quad \forall v \in H^1(\mathbb{R}). \end{cases}$$
(101)

Hermite–Galerkin method

 $\begin{cases} \operatorname{Find} u_N \in \hat{P}_N \operatorname{such} \operatorname{that} \\ (\partial_x u, \partial_x v) + \gamma(u, v) = \left(\hat{I}_N^h f, v_N\right), \quad \forall v_N \in \hat{P}_N. \end{cases}$ (102)

Theorem 26. Let $\gamma > 0$ and $\hat{\partial}_x = \partial_x + x$. If $u \in H^1(\mathbb{R})$ with $\hat{\partial}_x^m u \in L^2(\mathbb{R})$, and $f \in C(\mathbb{R})$ with $\hat{\partial}_x^k f \in L^2(\mathbb{R})$ and fixed $k, m \ge 1$, then

$$\|u_N - u\|_1 \lesssim N^{\frac{1-m}{2}} \|\hat{\partial}_x^m u\| + N^{\frac{1}{6} - \frac{k}{2}} \|\hat{\partial}_x^k f\|.$$
(103)

4.3 Numerical Results

Exact solution $u(x)$	$x \in (0,\infty)$	$x \in (-\infty, \infty)$
$u_1(x)$	$e^{-x}\sin(kx)$	$e^{-x^2}\sin\left(kx\right)$
$u_2(x)$	$\frac{1}{(1+x)^h}$	$\frac{1}{(1+x^2)^h}$
$u_3(x)$	$\frac{\sin\left(kx\right)}{(1+x)^h}$	$\frac{\sin\left(kx\right)}{(1+x^2)^h}$



Figure 1. (Laguerre) Exponetial decay solution $u_1(x)$ with $k = 4, \gamma = 1$.



Figure 2. (Laguerre) Algebraic decay solution $u_2(x)$: different *h* and *N*.



Figure 3. (Hermite) Exponetial decay $u_1(x)$, Algebraic deacy $u_2(x)$.

4.4 Scaling Factor

Assume that

 $|u(x)| \leq \varepsilon \quad \text{for } x > M.$

Scaling factor $(x_N \text{ is the largest L-G-R point})$

 $\beta_N = x_N / M,$

Scaled equation

 $-\beta_N^2 v_{yy} + \gamma v = g(y); \quad v(0) = 0, \quad \lim_{y \to +\infty} v(y) = 0.$ (104)



Figure 4. Scaling effects. $\beta_N = 15$.

5 Mapped Spectral Methods

5.1 Mappings

 $x = g(y; s), \quad y \in (-1, 1), \quad x \in \Lambda := (0, +\infty) \text{ or } (-\infty, +\infty)$ (105) with s > 0 and

$$\frac{\mathrm{d}x}{\mathrm{d}y} = g'(y;s) > 0.$$

$$g(-1;s) = 0, g(1;s) = +\infty, \quad \text{if } \Lambda = (0, +\infty)$$

$$g(\pm 1;s) = \pm\infty, \quad \text{if } \Lambda = (-\infty, +\infty)$$
(106)

- 1. Mapping between $x \in (-\infty, +\infty)$ and $y \in (-1, 1)$ with s > 0.
 - Algebraic mapping:

$$x = \frac{s y}{\sqrt{1 - y^2}}, \quad y = \frac{x}{\sqrt{x^2 + s^2}}.$$
 (107)

• Logarithmic mapping:

$$x = \frac{s}{2} \ln \frac{1+y}{1-y}, \quad y = \tanh \frac{x}{s} \tag{108}$$

• Exponential mapping:

$$x = \sinh(s y), \quad y = \frac{1}{s} \ln(x + \sqrt{x^2 + 1}), \quad (109)$$

where $y \in (-1, 1), \ x \in (-L_s, L_s), \text{ and } L_s = \sinh(s).$

- 2. Mapping between $x \in (0, +\infty)$ and $y \in (-1, 1)$ with s > 0:
 - Algebraic mapping:

$$x = s \frac{1+y}{1-y}, \quad y = \frac{x-s}{x+s}.$$
 (110)

• Logarithmic mapping

$$x = \frac{s}{2} \ln \frac{3+y}{1-y}, \quad y = 1 - 2 \tanh\left(\frac{x}{s}\right).$$
 (111)

• Exponential mapping

$$x = \sinh\left(s\frac{1+y}{2}\right), \quad y = \frac{2}{s}\ln\left(x+\sqrt{x^2+1}\right) - 1, \quad (112)$$

where $x \in (0, L_s), L_s = \sinh(s)$.

5.2 Approximation by Mapped Jacobi Polynomials

For the algebraic mapping, we define rational Jacobi polynomials:

$$j_{n,s}^{\alpha,\beta}(x) := J_n^{\alpha,\beta}(y) \tag{113}$$

then

$$\int_{\Lambda} j_{n,s}^{\alpha,\beta}(x) j_{n,s}^{\alpha,\beta}(x) \,\omega_s^{\alpha,\beta}(x) \,\mathrm{d}x = \gamma_n^{\alpha,\beta} \delta_{mn}, \qquad (114)$$

where $\gamma_n^{\alpha,\beta} = (J_n^{\alpha,\beta}, J_n^{\alpha,\beta})_{\omega^{\alpha,\beta}}$ and

$$\omega_s^{\alpha,\beta}(x) = \omega^{\alpha,\beta}(y) \frac{\mathrm{d}y}{\mathrm{d}x} = \omega^{\alpha,\beta}(y) \left(\frac{\mathrm{d}x}{\mathrm{d}y}\right)^{-1} > 0 \tag{115}$$

Note that for mapping (107) and (110), we have respectively

$$\frac{\mathrm{d}x}{\mathrm{d}y} = \frac{s}{\sqrt{1-y^2}^3} = s\,\omega^{-\frac{3}{2},-\frac{3}{2}}, \quad \frac{\mathrm{d}x}{\mathrm{d}y} = \frac{2\,s}{(1-y)^2} = 2\,s\,\omega^{-2,0}.$$



Figure 5. (Left) Graph of $j_{n,1}^{0,0}(x)$ with n = 8, 12 using mapping (107). (Right) Graphs of $j_{8,s}^{0,0}(x)$ with s = 0.2, 0.6, 1 using mapping (110).

Approximation

$$V_{N,s}^{\alpha,\beta} = \operatorname{span}\{j_{n,s}^{\alpha,\beta}(x) : n = 0, 1, ..., N\}, \quad s > 0.$$
(116)

Consider the orthogonal projection $\pi_{N,s}^{\alpha,\beta}: L^2_{\omega_s^{\alpha,\beta}}(\Lambda) \to V^{\alpha,\beta}_{N,s}:$

$$\left(\pi_{N,s}^{\alpha,\beta}u-u,v_{N}\right)_{\omega_{s}^{\alpha,\beta}}=0,\quad\forall v_{N}\in V_{N,s}^{\alpha,\beta}.$$
(117)

For a given mapping x = g(y, s), define

$$a_s(x) := \frac{\mathrm{d}x}{\mathrm{d}y}(>0), \quad U_s(y) := u(x) = u(g(y,s)).$$
 (118)

Then by defining $D_x u := a_s \frac{\mathrm{du}}{\mathrm{dx}}$, we have

$$\frac{\mathrm{d}U_s}{\mathrm{d}y} = a_s \frac{\mathrm{d}u}{\mathrm{d}x} = D_x u, \quad \frac{\mathrm{d}^k U_s}{\mathrm{d}y^k} = D_x^k u. \tag{119}$$

Define space

 $\tilde{B}^{m}_{\alpha,\beta}(\Lambda) = \left\{ u: u \text{ is measurable in } \Lambda \text{ and } \|u\|_{\tilde{B}^{m}_{\alpha,\beta}} < \infty \right\}$ (120) equipped with the norm and semi-norm

$$\begin{split} \|u\|_{\tilde{B}^{m}_{\alpha,\beta}} &= \left(\sum_{k=0}^{m} \|D^{k}_{x}u\|_{\omega_{s}^{\alpha+k,\beta+k}}^{2}\right)^{1/2}, \quad |u|_{\tilde{B}^{m}_{\alpha,\beta}} = \|D^{m}_{x}u\|_{\omega_{s}^{\alpha+m,\beta+m}}.\\ \\ \mathbf{Theorem 27.} \ \ Let \ \alpha, \beta > -1. \ \ If \ u \in \tilde{B}^{m}_{\alpha,\beta}(\Lambda), \ then \ for \ 0 \leqslant m \leqslant N+1, \\ \|\pi^{\alpha,\beta}_{N,s}u - u\|_{\omega_{s}^{\alpha,\beta}} \lesssim \sqrt{\frac{(N-m+1)!}{(N+1)!}} (N+m)^{-\frac{m}{2}} \|D^{m}_{x}u\|_{\omega_{s}^{\alpha+m,\beta+m}}, \quad (121) \\ and \ for \ 1 \leqslant m \leqslant N+1, \end{split}$$

$$\left\|\partial_x \left(\pi_{N,s}^{\alpha,\beta} u - u\right)\right\|_{\tilde{w}_s^{\alpha,\beta}} \lesssim \sqrt{\frac{(N-m+1)!}{N!}} (N+m)^{(1-m)/2} \|D_x^m u\|_{\omega_s^{\alpha+m,\beta+m}}, \quad (122)$$

where $\tilde{\omega}_s^{\alpha,\beta}(x) = \omega^{\alpha+1,\beta+1}g'(y;s), \quad y = h(x;s).$

For algebraic mapping on $(0, +\infty)$

1. For
$$u(x) = \frac{1}{(1+x)^{h}}$$
, $\|D_{x}^{m}u\|_{\omega_{s}^{\alpha+m,\beta+m}} < \infty$ if $m < 2h + \alpha + 1$:
 $\|u - \pi_{N,s}^{\alpha,\beta}u\|_{\omega_{s}^{\alpha,\beta}} \lesssim N^{-(2h+\alpha+1)}$, $(u(x) = (1+x)^{-h})$ (123)
2. For $u(x) = \frac{\sin(kx)}{(1+x)^{h}}$, $\|D_{x}^{m}u\|_{\omega_{s}^{\alpha+m,\beta+m}} < \infty$ if $m < \frac{2h+\alpha+1}{3}$:
 $\|u - \pi_{N,s}^{\alpha,\beta}u\|_{\omega_{s}^{\alpha,\beta}} \lesssim N^{-(2h+\alpha+1)/3}$, $\left(u(x) = \frac{\sin(kx)}{(1+x)^{h}}\right)$ (124)

For algebraic mapping on $(-\infty, +\infty)$

1. Algebraic decay

$$\|u - \pi_{N,s}^{\alpha,\beta} u\|_{\omega_s^{\alpha,\beta}} \lesssim N^{-(2h+\alpha+1)}, \quad (u(x) = (1+x^2)^{-h}) \quad (125)$$

2. Algebraic decay with oscillation

$$\left\| u - \pi_{N,s}^{\alpha,\beta} u \right\|_{\omega_s^{\alpha,\beta}} \lesssim N^{-(2h+\alpha+1)/2}, \quad (u(x) = \frac{\sin(kx)}{(1+x^2)^h}) \quad (126)$$

Remark 28.

- 1. For $u = (1+x)^{-h}$ and $u = (1+x^2)^{-h}$
 - a. if h positive integer, expressed exactly by finite sum of mapped rational functions
 - b. for other case, algebraic convergence rate but faster than approx. by Laguerre functions or Hermite functions.
- 2. For solutions with oscillation $u(x) = \frac{\sin(kx)}{(1+x)^h}$, $\frac{\sin(kx)}{(1+x^2)^h}$. The convergence rates are much slower than solution without oscillations. And also slower than Laguerre functions and Hermite functions methods.
- 3. For solutions with exponetial decay at infinity, the convergence rate is faster than any algebraic rate. $\sim e^{-c\sqrt{N}}$.

Mapped Interpolation

Given Jacobi-Gauss points $\{\xi_{N,j}^{\alpha,\beta}\}_{j=0}^N$, define $\zeta_{N,j,s}^{\alpha,\beta} := g(\xi_{N,j}^{\alpha,\beta}; s)$, and mapped Jacobi–Gauss interpolation $I_{N,s}^{\alpha,\beta}: C(\Lambda) \to V_{N,s}^{\alpha,\beta}$ as:

$$I_{N,s}^{\alpha,\beta} u \in V_{N,s}^{\alpha,\beta} \ s.t. \ \left(I_{N,s}^{\alpha,\beta} u\right) \left(\zeta_{N,j,s}^{\alpha,\beta}\right) = u\left(\zeta_{N,j,s}^{\alpha,\beta}\right), j = 0, 1, \dots, N.$$
(127)

Theorem 29. Let $\alpha, \beta > -1$. If $u \in \tilde{B}^m_{\alpha,\beta}(\Lambda)$ with $1 \leq m \leq N+1$, then

$$\|\partial_{x} (I_{N,s}^{\alpha,\beta}u - u)\|_{\tilde{\omega}_{s}^{\alpha,\beta}} + N \|I_{N,s}^{\alpha,\beta}u - u\|_{\omega_{s}^{\alpha,\beta}}$$

$$\lesssim \sqrt{\frac{(N - m + 1)!}{N!}} (N + m)^{(1 - m)/2} \|D_{x}^{m}u\|_{\omega_{s}^{\alpha + m,\beta + m}}.$$
(128)



Figure 6. (Left): Hermite–Gauss points ("o") vs. mapped Legendre–Gauss points using the algebraic map (107) with s = 1 ("•"). $\min_j |x_{j+1} - x_j| \sim O(N^{-1})$ for mapped Legendre–Gauss points.

(Right) Laguerre–Gauss–Radau points ("o") vs. mapped Legendre–Gauss-Radau points using the algebraic map (110) with s = 1 (" \star "). min_j $|x_{j+1} - x_j| \sim O(N^{-2})$ for mapped Legendre–Gauss–Radau points. 5.3 Spectral Methods Using Mapped Jacobi Polynomials

$$\gamma u - \partial_x (a(x)\partial_x u) = f, \quad x \in \Lambda = (-\infty, +\infty), \quad \gamma > 0$$
(129)

For a given map x = g(y; s), the weighted weak form

 $\begin{cases} \operatorname{Find} u \in \tilde{B}^{1}_{\alpha,\beta}(\Lambda) \text{ such that} \\ \gamma(u,v)_{\omega_{s}^{\alpha,\beta}} + \left(a(x)\partial_{x}u, \partial_{x}\left(v\,\omega_{s}^{\alpha,\beta}\right)\right) = (f,v)_{\omega_{s}^{\alpha,\beta}}, \forall v \in \tilde{B}^{1}_{\alpha,\beta}(\Lambda). \end{cases}$ (130)

The corresponding mapped Jacobi–Galerkin method is

Find
$$u_N \in V_{N,s}^{\alpha,\beta}$$
 such that $\forall v_N \in V_{N,s}^{\alpha,\beta}$
 $\gamma(u_N, v_N)_{\omega_s^{\alpha,\beta}} + (a(x)\partial_x u_N, \partial_x (v_N \omega_s^{\alpha,\beta})) = (I_{N,s}^{\alpha,\beta} f, v_N)_{\omega_s^{\alpha,\beta}}.$
(131)

A second approach: Jacobi approximation for transformed problem:

$$\gamma U_s(y) - \frac{1}{g'(y;s)} \partial_y \left(\frac{a(g(y;s))}{g'(y;s)} \partial_y U_s(y) \right) = F_s(y), \tag{132}$$

where $U_s(y) = u(g(y;s))$ and $F_s(y) = f(g(y;s))$. Let $\hat{\omega}_s^{\alpha,\beta}(y) = \omega^{\alpha,\beta}(y)g'(y;s)$. Jacobi-Gelerkin method for (132) is

$$\begin{cases} \operatorname{Find} \tilde{u}_{N} \in P_{N} s.t. \ \forall \ \tilde{v}_{N} \in P_{N} \\ \gamma(\tilde{u}_{N}, \tilde{v}_{N})_{\omega^{\alpha,\beta}} + \left(\frac{a(g(y;s))}{g'(y;s)} \partial_{y} \tilde{u}_{N}, \partial_{y} \left(\tilde{v}_{N} \hat{\omega}_{s}^{\alpha,\beta}\right)\right) = \left(I_{N}^{\alpha,\beta} F_{s}, \tilde{v}_{N}\right)_{\omega^{\alpha,\beta}} \end{cases}$$
(133)

This approach is in general more difficult to analyze, but can be easily implemented using the standard Jacobi–collocation method.

Error Estimates for a Model Problem

 $\gamma u(x) - \partial_x^2 u(x) = f(x), \quad x \in \Lambda = (0, \infty), \quad \gamma > 0; \quad u(0) = 0.$ (134)

Weak formulation

$$\begin{cases} \text{Find } u \in H^1_{0,\omega}(\Lambda) \text{ such that} \\ a_{\omega}(u,v) = (f,v)_{\omega}, \quad \forall v \in H^1_{0,\omega}(\Lambda), \end{cases}$$
(135)

where $\omega := \omega_s^{\alpha,\beta}$, $H_{0,\omega}^1 = \{ u \in H_\omega^1(\Lambda) : u(0) = 0 \}$, and

$$a_{\omega}(u,v) = \gamma(u,v)_{\omega} + (\partial_x u, \partial_x(v\,\omega)), \quad \forall u, v \in H^1_{0,\omega}(\Lambda).$$
(136)

Denote $X_N = \{ u \in V_{N,s}^{\alpha,\beta} : u(0) = 0 \}$. Jacobi–Galerkin approximation is

$$\begin{cases} \operatorname{Find} u_N \in X_N \operatorname{such} \operatorname{that} \\ a_{\omega}(u_N, v_N) = \left(I_{N,s}^{\alpha,\beta} f, v_N \right)_{\omega}, \quad \forall v_N \in X_n, \end{cases}$$
(137)

Lemma 30. Assume that

$$d_1 = \max_{x \in \bar{\Lambda}} |\omega^{-1}(x)\partial_x \omega(x)|, \quad d_2 = \max_{x \in \bar{\Lambda}} |\omega^{-1}(x)\partial_x^2 \omega(x)|$$

are finite. Then, for any $u, v \in H^1_{\omega}(\Lambda)$,

 $a_{\omega}(u,v) \leq (d_1+1) |u|_{1,\omega} ||v||_{1,\omega} + \gamma ||u||_{\omega} ||v||_{\omega}.$

If, in addition $v^2(x)\omega'(x)|_{x=0}=0$ and $\lim_{x\to\infty}v^2\omega'(x) \ge 0$, then

$$a_{\omega}(v,v) \ge |v|_{1,\omega}^2 + (\gamma - d_2/2) \|v\|_{\omega}^2, \quad \forall v \in H^1_{\omega}(\Lambda).$$
(138)

Theorem 31. Assume that the condition of Lemma 30 are satisfied and $\gamma - \frac{d_2}{2} > 0$. Then the problem (137) admits a unique solution. We have

$$\|u - u_N\|_{1,\omega} \lesssim \inf_{v_N \in X_N} \|u - v_N\|_{1,\omega} + \|f - I_{N,s}^{\alpha,\beta}f\|_{\omega}.$$
 (139)
For mapped Legendre method $(d_1 \leq 2, d_2 \leq 6)$

Corollary 32. Let u and u_N be respectively the solution of (135) and (137) with $(\alpha, \beta) = (0, 0)$ and the mapping (110) with s = 1. Assume that $u \in \tilde{B}_{0,0}^m(\Lambda)$ and $f \in \tilde{B}_{0,0}^k(\Lambda)$, $k, m \ge 1$ and $\gamma > 3$, we have

$$\|u - u_N\|_{1,\omega_1^{0,0}} \lesssim N^{1-m} \|D_x^m u\|_{\omega_1^{m,m}} + N^{-k} \|D_x^k f\|_{\omega_1^{k,k}}, \qquad (140)$$

For mapped Chebyshev method. We have $d_1, d_2 = \infty$. but $a_{\omega}(\cdot, \cdot)$ is still continuous and coercive [cf. Guo et al. 2002].

Corollary 33. Let u and u_N be the solution of (135) and (137) with $(\alpha, \beta) = (-1/2, -1/2)$ and the mapping (110) with s = 1. Assume that $u \in \tilde{B}^m_{-1/2, -1/2}(\Lambda)$ and $f \in \tilde{B}^m_{-1/2, -1/2}(\Lambda)$ and that $\gamma > \frac{14}{17}$, we have

$$\|u - u_N\|_{1,\omega_1^{-1/2,-1/2}} \lesssim N^{1-m} \|D_x^m u\|_{\omega_1^{m-1/2,m-1/2}} + N^{-k} \|D_x^k f\|_{\omega_1^{k-1/2,k-1/2}},$$
(141)

Implementation and Numerical Results

Take $\phi_k(x) = j_{s,k}^{0,0}(x) + j_{s,k+1}^{0,0}$ with s = 1 as basis of X_N , then $\omega(x) = \frac{2}{(1+x)^2}$.

(137) leads to a non-symmetric sparse system (M tridiagonal, S seven diagonal)

$$(\gamma M + S)\boldsymbol{u} = \boldsymbol{f} \tag{142}$$



Figure 7. (Left) $u(x) = \sin(2x)e^{-x}$; (Right) $u(x) = 1/(1+x)^{5/2}$.



Figure 8. (Left) $u(x) = \sin(2x)/(1+x)^{7/2}$; (Right) $u(x) = \sin(2x)e^{-x^2}$.



Figure 9. (Left) $u(x) = 1/(1+x^2)^{5/2}$; (Right) $u(x) = \frac{\sin(2x)}{(1+x^2)^{7/2}}$.

Conclusion

- 1. Two approaches to solve problem in unbounded domain
 - a. Laguerre/Hermite method
 - b. Mapped Jacobi method
- 2. The convergence rates of two approaches are compatible
- 3. Scaling parameter affects the convergence rates.