Multilevel Iterative Methods

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Abstract

Over the last few decades, intensive research has been done on developing efficient parallel iterative solvers for PDEs. One useful mathematical technique that has drawn a lot of attention is multilevel iterative solvers and preconditioners. The lecture note is originally prepared for a semester-long course at LSEC. In this note, we will discuss algorithm, analysis, and applications of multilevel iterative solvers for partial differential equations. The note is mainly based on Prof. Jinchao Xu's short courses at the Peking University in 2013 and at the Academy of Mathematics and Systems Science in 2016 and Prof. Ludmil Zikatanov's summer school lectures at the Academy of Mathematics and Systems Science in 2015.

Contents

Part I

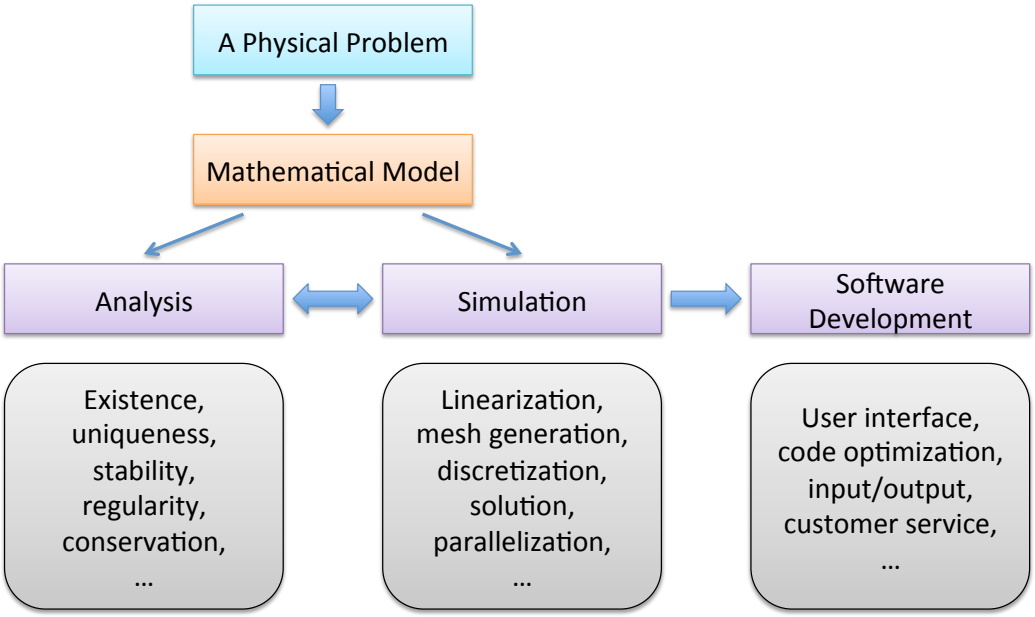
Theory of Multilevel Iterative Methods

Chapter 1

Introduction

ch:intro

Computer simulation has become an important tool in engineering and sciences. Many physical problems in scientific and engineering computing can be reduced to the numerical solution of certain partial differential equations (PDEs). Finding a viable solution to underlying discretized systems is often expensive, generally consuming a significant portion of the overall cost in a numerical solution procedure of PDEs. Various fast solution techniques, such as adaptive mesh refinement (AMR), domain decomposition (DD) methods, and multigrid (MG) methods, have been developed to address this issue.



ig:simulation

Figure 1.1: Numerical simulation of a physical problem.

The above diagram gives a simple illustration of how a physical problem is "solved" via numerical simulation in general. It is basically an interplay of modeling, mathematical analy-

sis, numerical analysis, scientific computing, and software engineering. A successful computer simulation of complicated physical phenomena requires expertise in many scientific subjects. Hence, nowadays it is difficult for one person to manage all these areas and close collaborations of experts from different areas become crucial.

Effective linear solvers play a key role in many application areas in scientific computing. There are many different types of algorithms for solving linear systems. In this lecture, we focus on studying algorithmic and theoretical aspects of multilevel iterative methods, including GMG and AMG. The basic problem setting for our discussion is: Given an invertible matrix $A : \mathbb{R}^{N\times N}$ and a vector $\vec{f} \in \mathbb{R}^N$, find $\vec{u} \in \mathbb{R}^N$ such that $A\vec{u} = \vec{f}$. There are many features of linear solver that we desire in practice, including:

- Convergence — The method should converge to somewhere at least.
- Robustness — The method should behave similarly for different conditions.
- Optimality — The method can give a solution with $O(N)$ computational cost.
- Efficiency — The method can give a solution with "reasonably short" wall time.
- Scalability — The method can scale well on modern parallel architectures.
- Reliability — The method should converge to a solution with limited amount of time.
- Usability — The method can be implemented and used relatively easily.

These above features sometimes contradict with each other and we have to find a good balance in practice. There are many different solution methods available. In this lecture, including direct solvers and iterative solvers. In this lecture, we will discuss several popular multilevel iterative methods, including the overlapping domain decomposition methods with coarse space corrections, two-grid methods, geometric multigrid methods, algebraic multigrid methods. And we will mainly study the convergence theory of these methods using the subspace correction framework.

1.1 The model equation

Let $\Omega \subset \mathbb{R}^d$ be an open and bounded domain with Lipschitz boundary and $f \in L^2(\Omega)$. We consider solution of the Poisson's equation with Dirichlet boundary condition

$$\begin{cases} -\Delta u = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{cases} \tag{1.1}$$

eqn:Poisson

This equation will be our main model equation in most part of this lecture.

Remark 1.1 (Diffusion equation in various applications)**.** The Poisson's equation, or more generally the diffusion equation, appears in many areas of physics, for example, Fick's law for chemical concentration, Fourier's law for temperature, Ohm's law for electrostatic potential, Darcy's law for porous media flow.

Derivation and classical solution ⋆

The concept of diffusion is widely used in physics, chemistry, biology, sociology, economics, and finance. It is the net movement of the object (e.g. molecules or atoms) from a region of high concentration (or high chemical potential) to a region of low concentration (or low chemical potential). This is also referred to as the movement of a substance down a concentration gradient.

Let $u(x)$ be some diffusive quantity, like pressure, temperature, or concentration of a biological species. We define the operator $\nabla := (\partial_1, \dots, \partial_d)^T$. So the gradient of scalar function $u : \Omega \mapsto \mathbb{R}$ can be denoted by ∇u. The Laplace operator can be written as $\Delta u = \nabla \cdot \nabla u$. A diffusive flux $\vec{F}$ is usually proportional to the gradient of u, i.e.,

$$\vec{F} = -\kappa \nabla u. \tag{1.2}$$ eqn:flux

where κ is the diffusivity (e.g., heat conductivity or permeability). Note that $-\nabla u$ is the so-called steepest descent direction. If a flow is controlled solely by diffusion, then the mass conservation in any volume ω can be written, in the integral form, as

$$\frac{\partial}{\partial t} \int_\omega u \, dx = -\int_{\partial\omega} \vec{F} \cdot \vec{\nu} \, dS$$

or, in the strong form, as

$$\frac{\partial}{\partial t} u = -\nabla \cdot \vec{F}. \tag{1.3}$$ eqn:masscon

This can be seen by applying the Divergence Theorem

$$\int_\omega \nabla \cdot \vec{F} \, dx = \int_{\partial\omega} \vec{F} \cdot \vec{\nu} dS. \tag{1.4}$$ eqn:divthm

Now, by plugging (1.2) into (7.10), we obtain an equation

$$\frac{\partial}{\partial t} u = \nabla \cdot (\kappa \nabla u). \tag{1.5}$$ eqn:heat1

If we assume $\kappa \equiv 1$ or just a constant and there is a source/sink term f on Ω, then we arrive at the heat equation

$$\frac{\partial}{\partial t} u - \Delta u = f. \tag{1.6}$$ eqn:heat

The steady-state solution of equation (1.6) satisfies the well-known Poisson's equation

$$-\Delta u = f. \tag{1.7}$$ eqn:poisson

Remark 1.2 (Laplace equation)**.** In case of the body force or source/sink term is zero, the equation is usually referred to as the Laplace equation

$$-\Delta u = 0. \tag{1.8}$$ eqn:laplace

If $u \in C^2(\Omega)$ and $-\Delta u = 0$, the u is called a *harmonic function*.

We have the fundamental solution of the Poisson's equation

$$\Phi(x) := \begin{cases} -\frac{1}{2\pi}\log|x|, & d = 2 \\ \frac{1}{d(d-2)\alpha(d)}|x|^{2-d}, & d \geqslant 3 \end{cases} \tag{1.9}$$ eqn:fsolu

where $\alpha(d)$ is the volume of the unit ball in $\mathbb{R}^d$. It is well-known that

$$u(x) = \Phi * f := \int_{\mathbb{R}^d} \Phi(x-y) f(y)\, dy$$

satisfies $-\Delta u = f$ in $\mathbb{R}^d$ and $u \in C^2(\mathbb{R}^d)$; see Evans [30].

Theorem 1.1 (Strong Maximum Principle)**.** *If $u \in C^2(\Omega)\bigcap C(\overline{\Omega})$ is harmonic in Ω, then the maximal value*

$$\max_{x\in\overline{\Omega}} u(x) = \max_{x\in\partial\Omega} u(x).$$

If the domain Ω is connected, then $u \equiv C$ if there exist $x_0 \in \Omega$ such that

$$u(x_0) = \max_{x\in\overline{\Omega}} u(x).$$

Theorem 1.2 (Uniqueness of solution)**.** *If $f \in C(\Omega)$, then there exists at most one solution $u \in C^2(\Omega)\bigcap C(\overline{\Omega})$.*

Sobolev spaces ⋆

The standard L^∞-norm and L^2-norm will be denoted by $\|\cdot\|_\infty$ and $\|\cdot\|_0$, respectively. The symbol $L^2_0(\Omega)$ denotes a subspace of $L^2(\Omega)$ consisting of functions that have a zero average. The bilinear forms $(\cdot,\cdot)$ and $\langle\cdot,\cdot\rangle$ denote the classical L^2*-inner product* and the *duality pair*, respectively.

Given a natural number $k \in \mathbb{N}$ and $1 \leqslant p \leqslant \infty$, we define the Sobolev spaces

$$W_p^k(\Omega) := \big\{v : \Omega \mapsto \mathbb{R} \,:\, \nabla^\alpha v \in L^p(\Omega),\ \text{for all } |\alpha| \leqslant k\big\}, \tag{1.10}$$ eqn:W_kp

where $\alpha = [\alpha_1, \ldots, \alpha_d]$ is a multi-index and $\nabla^\alpha v := \partial_{x_1}^{\alpha_1}\cdots\partial_{x_d}^{\alpha_d} v$ is the weak derivative of order α. The corresponding norm and semi-norm are then defined as follows: for $1 \leqslant p < \infty$,

$$\|v\|_{W_p^k(\Omega)} := \Big(\sum_{|\alpha|\leqslant k} \|\nabla^\alpha v\|^p_{L^p(\Omega)}\Big)^{\frac{1}{p}}, \quad |v|_{W_p^k(\Omega)} := \Big(\sum_{|\alpha|= k} \|\nabla^\alpha v\|^p_{L^p(\Omega)}\Big)^{\frac{1}{p}}, \tag{1.11}$$ eqn:wkp_norm

and, for $p = \infty$,

$$\|v\|_{W_\infty^k(\Omega)} := \sup_{|\alpha|\leqslant k} \|\nabla^\alpha v\|_{L^\infty(\Omega)}, \quad |v|_{W_\infty^k(\Omega)} := \sup_{|\alpha|= k} \|\nabla^\alpha v\|_{L^\infty(\Omega)}. \tag{1.12}$$ eqn:wkinf_nor

Definition 1.1 (Sobolev number)**.** Consider a change of variables $\hat{x} = x/h$ for $h > 0$ and all $x \in \Omega$, which transforms the domain Ω to $\hat{\Omega}$. With this transformation, a function $v(x)$ on Ω becomes $\hat{v}(\hat{x})$ on $\hat{\Omega}$. We notice that

$$|\hat{v}|_{W_p^k(\hat{\Omega})} = h^{\mathrm{sob}(W_p^k)} |v|_{W_p^k(\Omega)}, \tag{1.13}$$ eqn:scaling

where the Sobolev number is defined by

$$\mathrm{sob}(W_p^k) := k - \frac{d}{p}. \tag{1.14}$$ eqn:sobnum

Remark 1.3 (Natural scaling)**.** There is a natural scaling for the semi-norm $|\cdot|_{W_p^k(\Omega)}$. For $h > 0$, we apply the change of variable $\hat{x} = x/h : \Omega \mapsto \hat{\Omega}$. Then the following scaling result holds

$$\big|\hat{v}\big|_{W_p^k(\hat{\Omega})} = h^{k-\frac{d}{p}} \big|v\big|_{W_p^k(\Omega)} = h^{\mathrm{sob}(W_p^k)} \big|v\big|_{W_p^k(\Omega)}.$$

This property is useful in scaling argument (or homogeneity argument) for finite element error estimates.

If $p = 2$, the spaces $W_2^k(\Omega)$ are Hilbert spaces and we denote them by $H^k(\Omega)$ for short. The inner product is given by

$$(u, v)_{k,\Omega} := (u, v)_{H^k(\Omega)} := \sum_{|\alpha| \leqslant k} \int_\Omega \nabla^\alpha u \, \nabla^\alpha v \, dx.$$

The induced norm of this scalar product is the $W_2^k(\Omega)$-norm. We denote $H_0^k(\Omega)$ the completion of $C_0^\infty(\Omega)$ in $H^k(\Omega)$. We will also use the fractional Sobolev space $H_0^{k+\sigma}(\Omega)$ where $0 < \sigma < 1$. It is defined as the completion of $C_0^\infty(\Omega)$ in the fraction norm:

$$\|v\|_{H^{k+\sigma}(\Omega)} := \Big(\|v\|^2_{H^k(\Omega)} + |v|^2_{H^{k+\sigma}(\Omega)} \Big)^{\frac{1}{2}},$$

where

$$|v|_{H^{k+\sigma}(\Omega)} := \Bigg(\sum_{|\alpha|=k} \int_\Omega \int_\Omega \frac{|D^\alpha v(x) - D^\alpha v(y)|^2}{|x-y|^{d+2\sigma}} \, dx \, dy \Bigg)^{\frac{1}{2}}.$$

Before we discuss the Poisson's equation in weak formulation, we introduce a few important properties of the Sobolev spaces, which will become important in our later analysis for multigrid methods.

rop:embedding **Proposition 1.1** (Sobolev embedding)**.** *Let $0 \leqslant k < m$. If $sob(W_p^m) > sob(W_q^k)$, then the embedding $W_p^m(\Omega) \hookrightarrow W_q^k(\Omega)$ is compact.*

prop:P-W **Proposition 1.2** (Poincaré-Wirtinger inequality)**.** *For any $v \in H^1(\Omega)$, we have*

$$\Big\| v - |\Omega|^{-1} \int_\Omega v \, dx \Big\|_{0,\Omega} \leqslant C(\Omega) \big|v\big|_{1,\Omega}.$$

prop:Poincare **Proposition 1.3** (Poincaré inequality). *For any $v \in H_0^1(\Omega)$, we have*

$$\|v\|_{0,\Omega} \leqslant C_d |\Omega|^{1/d} |v|_{1,\Omega}.$$

It is a special case of the more general Friedrichs' inequality on $W_p^k(\Omega)$ with zero trace and it is sometimes referred to as the Friedrichs–Poincaré inequality.

prop:trace **Proposition 1.4** (Trace theorem). *There exists a unique linear operator* trace $: H^1(\Omega) \mapsto L^2(\partial\Omega)$, *such that* $\mathrm{trace}(v) = v$, *if $v \in C^0(\overline{\Omega}) \bigcap H^1(\Omega)$, and*

$$\|\mathrm{trace}(v)\|_{0,\partial\Omega} \leqslant C(\Omega)\|v\|_{1,\Omega}, \quad \forall v \in H^1(\Omega).$$

Moreover, if $g \in H^{\frac{1}{2}}(\partial\Omega)$, there exists $\phi \in H^1(\Omega)$ such that $\phi|_{\partial\Omega} = g$ and

$$\|\phi\|_{1,\Omega} \leqslant C\|g\|_{\frac{1}{2},\partial\Omega}.$$

Weak formulation

Now we consider the Poisson's equation in a weaker sense. A simple motivation is to convert from a point-wise view to an average view:

$$u(x) = 0, \ \text{a.e.} \quad \Longleftrightarrow \quad \int_\Omega uv\,dx = 0, \quad \forall v \in C_0^\infty(\Omega).$$

Similarly, we can write the Poisson's equation in the weak form (i.e., the integral form). In the one-dimensional case, it is easy to see that

$$-u'' = f, \ \text{a.e.} \quad \Longleftrightarrow \quad -\int_\Omega (u'' + f)v\,dx = 0, \quad \forall v \in C_0^\infty(\Omega).$$

Let $\mathscr{U}$ be a Hilbert space with an inner product $(\cdot,\cdot)_{\mathscr{U}}$ and its induced norm $\|\cdot\|_{\mathscr{U}}$. Let $\mathscr{V}$ be a Hilbert space with another inner product $(\cdot,\cdot)_{\mathscr{V}}$ and its induced norm $\|\cdot\|_{\mathscr{V}}$. Denote by $\mathscr{V}'$ the dual space of $\mathscr{V}$ equipped with the norm

$$\|f\|_{\mathscr{V}'} := \sup_{v\in\mathscr{V}} \frac{\langle f, v\rangle}{\|v\|_{\mathscr{V}}}, \quad \forall f \in \mathscr{V}'.$$

Definition 1.2 (Continuity). *A bilinear form $a[\cdot,\cdot] : \mathscr{U} \times \mathscr{V} \mapsto \mathbb{R}$ is called continuous if and only if there exists a constant C_a such that*

$$a[u,v] \leqslant C_a\|u\|_{\mathscr{U}}\|v\|_{\mathscr{V}}, \quad \forall\, u \in \mathscr{U}, v \in \mathscr{V}. \tag{1.15}$$ eqn:cont_bili

Consider a continuous bilinear form $a[\cdot,\cdot] : \mathscr{U} \times \mathscr{V} \to \mathbb{R}$ and $f \in \mathscr{V}'$. We formulate a model problem: Find $u \in \mathscr{U}$ such that $\mathcal{A}u = f$ in $\mathscr{V}'$. Or in the weak form, find $u \in \mathscr{U}$ such that

$$a[u,v] = \langle f, v\rangle, \quad \forall v \in \mathscr{V}. \tag{1.16}$$ prob:model

g:WeakPoisson **Example 1.1** (The Poisson equation). The Poisson problem with homogenous Dirichlet boundary was given in (1.1). In this case, we have $\mathcal{A}u := -\Delta u$ and $a[u,v] := (\nabla u, \nabla v)$. Apparently, the bilinear form $a[\cdot,\cdot]$ is continuous due to the Cauchy–Schwarz inequality and $\mathscr{U} = \mathscr{V} = H_0^1(\Omega)$. □

Well-posedness of the weak problem ⋆

ssc:bt

We denote the space of all linear and continuous operators from $\mathscr{U}$ to $\mathscr{V}$ as $\mathscr{L}(\mathscr{U};\mathscr{V})$. Here we review a few results on the inf-sup condition due to Nečas [46].

:banach_necas **Theorem 1.3** (Banach–Nečas Theorem)**.** *Let $a[\cdot,\cdot]:\mathscr{U}\times\mathscr{V}\mapsto\mathbb{R}$ be a continuous bilinear form with a norm defined as*

$$\|a[\cdot,\cdot]\| := \sup_{u\in\mathscr{U}}\sup_{v\in\mathscr{V}}\frac{a[u,v]}{\|u\|_{\mathscr{U}}\|v\|_{\mathscr{V}}}.$$

(i) Then there exists a unique linear operator $\mathcal{A}\in\mathscr{L}(\mathscr{U};\mathscr{V})$ such that

$$(\mathcal{A}u,v)_{\mathscr{V}} = a[u,v],\quad \forall u\in\mathscr{U}, v\in\mathscr{V},$$

with the operator norm

$$\|\mathcal{A}\|_{\mathscr{L}(\mathscr{U};\mathscr{V})} = \|a[\cdot,\cdot]\|.$$

(ii) Moreover, the bilinear form $a[\cdot,\cdot]$ satisfies the inf-sup condition:

$$\exists\,\alpha>0,\ \text{ such that }\ \alpha\|u\|_{\mathscr{U}}\leqslant\sup_{v\in\mathscr{V}}\frac{a[u,v]}{\|v\|_{\mathscr{V}}},\quad\forall\,u\in\mathscr{U}, \tag{1.17}$$ cond:infsup

$$\text{for any }\ 0\neq v\in\mathscr{V},\ \text{ there exists }\ u\in\mathscr{U},\ \text{ such that }\ a[u,v]\neq 0, \tag{1.18}$$ cond:nonsin

if and only if $\mathcal{A}:\mathscr{U}\mapsto\mathscr{V}$ is an isomorphism and

$$\|\mathcal{A}^{-1}\|_{\mathscr{L}(\mathscr{V};\mathscr{U})}\leqslant\alpha^{-1}. \tag{1.19}$$ eqn:AinvBound

Proof. (i) For any fixed $u\in\mathscr{U}$, the mapping $a[u,\cdot]$ belongs to the dual space $\mathscr{V}'$. By the Riesz representation theorem, there exists $\mathcal{A}u\in\mathscr{V}$ such that

$$(\mathcal{A}u,v)_{\mathscr{V}} = a[u,v],\quad\forall\,v\in\mathscr{V}.$$

Since $a[\cdot,\cdot]$ is continuous, we obtain a bounded operator $\mathcal{A}\in\mathscr{L}(\mathscr{U};\mathscr{V})$. Furthermore,

$$\|\mathcal{A}\|_{\mathscr{L}(\mathscr{U};\mathscr{V})} = \sup_{u\in\mathscr{U}}\frac{\|\mathcal{A}u\|_{\mathscr{V}}}{\|u\|_{\mathscr{U}}} = \sup_{u\in\mathscr{U}}\sup_{v\in\mathscr{V}}\frac{(\mathcal{A}u,v)_{\mathscr{V}}}{\|u\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} = \sup_{u\in\mathscr{U}}\sup_{v\in\mathscr{V}}\frac{a[u,v]}{\|u\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} = \|a[\cdot,\cdot]\|.$$

(ii) $\Longrightarrow$ The inf-sup condition (1.17) guarantees that there exists $\alpha>0$ such that

$$\alpha\|u\|_{\mathscr{U}}\leqslant\sup_{v\in\mathscr{V}}\frac{a[u,v]}{\|v\|_{\mathscr{V}}} = \sup_{v\in\mathscr{V}}\frac{(\mathcal{A}u,v)_{\mathscr{V}}}{\|v\|_{\mathscr{V}}} = \|\mathcal{A}u\|_{\mathscr{V}},\quad\forall u\in\mathscr{U}. \tag{1.20}$$ eqn:ALowerBou

This implies that $\mathcal{A}$ is injective. Let $\{u_k\}_{k=0}^{\infty}\subset\mathscr{U}$ and $v_k := \mathcal{A}u_k$ be a sequence such that $v_k\to v\in\mathscr{V}$. In order to show the range of $\mathcal{A}$ is closed, we need to show $v\in\mathcal{A}(\mathscr{U})$. From the inequality (1.20), we have

$$\alpha\|u_k-u_j\|_{\mathscr{U}}\leqslant\|\mathcal{A}(u_k-u_j)\|_{\mathscr{V}} = \|v_k-v_j\|_{\mathscr{V}}\to 0.$$

Hence, $\{u_k\}_{k=0}^{\infty}$ is a Cauchy sequence and $u_k \to u \in \mathscr{U}$. Moreover,

$$v = \lim_{k\to\infty} v_k = \lim_{k\to\infty} \mathcal{A}u_k = \mathcal{A}u \in \mathcal{A}(\mathscr{U}).$$

Now we assume that $\mathcal{A}(\mathscr{U}) \neq \mathscr{V}$. Since $\mathcal{A}(\mathscr{U})$ is closed, we can decompose $\mathscr{V}$ as

$$\mathscr{V} = \mathcal{A}(\mathscr{U}) \oplus \mathcal{A}(\mathscr{U})^{\perp}$$

and $\mathcal{A}(\mathscr{U})^{\perp}$ is non-trivial. That is to say, there exists $0 \neq v_{\perp} \in \mathcal{A}(\mathscr{U})^{\perp}$, which contradicts the condition (1.18). Hence the assumption $\mathcal{A}(\mathscr{U}) \neq \mathscr{V}$ cannot hold, i.e., $\mathcal{A}$ is surjective. This, in turn, shows that $\mathcal{A}$ is an isomorphism from $\mathscr{U}$ onto $\mathscr{V}$. Moreover, (1.20) shows

$$\alpha \|\mathcal{A}^{-1}v\|_{\mathscr{U}} \leqslant \|v\|_{\mathscr{V}}, \quad \forall\, v \in \mathscr{V}.$$

This proves the inequality (1.19).

(ii) $\Longleftarrow$ We have

$$\begin{aligned}\inf_{u\in\mathscr{U}} \sup_{v\in\mathscr{V}} \frac{a[u,v]}{\|u\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} &= \inf_{u\in\mathscr{U}} \sup_{v\in\mathscr{V}} \frac{(\mathcal{A}u,v)}{\|u\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} = \inf_{u\in\mathscr{U}} \frac{\|\mathcal{A}u\|_{\mathscr{V}}}{\|u\|_{\mathscr{U}}} \\ &= \inf_{v\in\mathscr{V}} \frac{\|v\|_{\mathscr{V}}}{\|\mathcal{A}^{-1}v\|_{\mathscr{U}}} = \Big(\sup_{v\in\mathscr{V}} \frac{\|\mathcal{A}^{-1}v\|_{\mathscr{U}}}{\|v\|_{\mathscr{V}}}\Big)^{-1} = \|\mathcal{A}^{-1}\|^{-1}_{\mathscr{L}(\mathscr{V};\mathscr{U})} \geqslant \alpha.\end{aligned}$$

This is exactly (1.17). Since $\mathcal{A}$ is an isomorphism, for any $0 \neq v \in \mathscr{V}$, there exists $0 \neq u \in \mathscr{U}$, such that $\mathcal{A}u = v$ and

$$a[u,v] = (\mathcal{A}u, v) = \|v\|^2_{\mathscr{V}} \neq 0,$$

which is (1.18). □

thm:necas **Theorem 1.4** (Nečas Theorem)**.** *Let $a[\cdot,\cdot] : \mathscr{U} \times \mathscr{V} \mapsto \mathbb{R}$ be a continuous bilinear form. Then the equation (1.16) admits a unique solution $u \in \mathscr{U}$ for all $f \in \mathscr{V}'$, if and only if the bilinear form $a[\cdot,\cdot]$ satisfies one of the equivalent inf-sup conditions:*

(1) There exists $\alpha > 0$ such that

$$\sup_{v\in\mathscr{V}} \frac{a[w,v]}{\|v\|_{\mathscr{V}}} \geqslant \alpha\|w\|_{\mathscr{U}}, \quad \forall\, w \in \mathscr{U}; \tag{1.21}$$ eqn:inf_sup_1

and for every $0 \neq v \in \mathscr{V}$, there exists $w \in \mathscr{U}$ such that $a[w,v] \neq 0$.

(2) There holds

$$\inf_{w\in\mathscr{U}} \sup_{v\in\mathscr{V}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} > 0 \qquad \textit{and} \qquad \inf_{v\in\mathscr{V}} \sup_{w\in\mathscr{U}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} > 0. \tag{1.22}$$ eqn:inf_sup_2

(3) There exists a positive constant $\alpha > 0$ such that

$$\inf_{w\in\mathscr{U}} \sup_{v\in\mathscr{V}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} = \inf_{v\in\mathscr{V}} \sup_{w\in\mathscr{U}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} = \alpha. \tag{1.23}$$ eqn:inf_sup_3

Furthermore, the solution u satisfies the stability condition

$$\|u\|_{\mathscr{U}} \leqslant \alpha^{-1}\|f\|_{\mathscr{V}'}.$$

Proof. Let $\mathcal{J} : \mathscr{V} \mapsto \mathscr{V}'$ be the isometric Reisz isomorphism. According to Theorem 1.3, we have $\mathcal{A} \in \mathscr{L}(\mathscr{U};\mathscr{V})$, which is the linear operator corresponding to $a[\cdot,\cdot]$. In this sense, (1.16) is equivalent to

$$u \in \mathscr{U} : \quad \mathcal{A}u = \mathcal{J}^{-1}f \quad \text{in } \mathscr{V}.$$

Assume the condition (1) holds. Then, $\mathcal{A}$ is invertible by Theorem 1.3. The other direction is also easy to see.

Now the interesting part is to show the equivalence of the three conditions, (1), (2), and (3). From the proof of Theorem 1.3, we have seen that

$$\inf_{w\in\mathscr{U}} \sup_{v\in\mathscr{V}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} = \|\mathcal{A}^{-1}\|^{-1}_{\mathscr{L}(\mathscr{V};\mathscr{U})}.$$

Similarly,

$$\begin{aligned}\inf_{v\in\mathscr{V}} \sup_{w\in\mathscr{U}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} &= \inf_{v\in\mathscr{V}} \sup_{w\in\mathscr{U}} \frac{(\mathcal{A}w,v)_{\mathscr{V}}}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} = \inf_{v\in\mathscr{V}} \sup_{w\in\mathscr{U}} \frac{(w,\mathcal{A}^{\dagger}v)_{\mathscr{U}}}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} \\ &= \|\mathcal{A}^{-\dagger}\|^{-1}_{\mathscr{L}(\mathscr{U};\mathscr{V})} = \|\mathcal{A}^{-1}\|^{-1}_{\mathscr{L}(\mathscr{V};\mathscr{U})},\end{aligned}$$

where $\mathcal{A}^{\dagger}$ denotes the adjoint operator. Furthermore, if the condition

$$\inf_{v\in\mathscr{V}} \sup_{w\in\mathscr{U}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} > 0$$

holds, then for any $v \in \mathscr{V}$, we have

$$\sup_{w\in\mathscr{U}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} > 0.$$

Hence there exists $w \in \mathscr{U}$, such that $a[w,v] \neq 0$. This completes the equivalence proof. □

From the proof of the last two theorems, we have the following observations:

Remark 1.4 (Existence and uniqueness)**.** Solution of the equation (1.16) exists (i.e., $\mathcal{A}$ is *surjective* or *onto*) if and only if

$$\inf_{v\in\mathscr{V}} \sup_{w\in\mathscr{U}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} > 0. \qquad \text{existence or surjective}$$

Solution of (1.16) is unique (i.e., $\mathcal{A}$ is *injective* or *one-to-one*) if and only if

$$\inf_{w\in\mathscr{U}} \sup_{v\in\mathscr{V}} \frac{a[w,v]}{\|w\|_{\mathscr{U}}\|v\|_{\mathscr{V}}} > 0. \qquad \text{uniqueness or injective}$$

That is to say, $\mathcal{A}$ is bijective if and only if the inf-sup conditions (1.22) or its equivalent conditions hold. In finite dimensional spaces, any linear surjective or injective map is also bijective. So we only need one of the above inf-sup conditions to show well-posedness.

Remark 1.5 (Optimal constant)**.** The constant α in (1.23) is the largest possible constant in (1.21). In general, the first condition in Theorem 1.4 is easier to verify than the third condition.

Corollary 1.1 (Well-posedness and inf-sup condition)**.** *If the weak formulation* (1.16) *has a unique solution* $u \in \mathscr{U}$ *for any* $f \in \mathscr{V}'$ *so that*

$$\|u\|_{\mathscr{U}} \leqslant C\|f\|_{\mathscr{V}'},$$

then the bilinear form $a[\cdot,\cdot]$ *satisfies the inf-sup condition* (1.23) *with* $\alpha \geqslant C^{-1}$.

Proof. Since (1.16) has a unique solution for all $f \in \mathscr{V}'$, the operator $\mathcal{A} : \mathscr{L}(\mathscr{U};\mathscr{V})$ is invertible and $\mathcal{A}^{-1} : \mathscr{L}(\mathscr{V};\mathscr{U})$ is bounded. Due to the fact $\|u\|_{\mathscr{U}} \leqslant C\|f\|_{\mathscr{V}'}$, we have $\|\mathcal{A}^{-1}\|_{\mathscr{L}(\mathscr{V};\mathscr{U})} \leqslant C$. From the proof of the Nečas theorem, we can immediately see the optimal inf-sup constant $\alpha = \|\mathcal{A}^{-1}\|^{-1}_{\mathscr{L}(\mathscr{V};\mathscr{U})} \geqslant C^{-1}$. □

A simple model problem

From now on, we consider the simplest case where $\mathscr{V} = \mathscr{U}$ and $\mathcal{A}$ is coercive.

Definition 1.3 (Coercivity)**.** *A continuous bilinear form* $a[\cdot,\cdot] : \mathscr{V} \times \mathscr{V} \mapsto \mathbb{R}$ *is called coercive if there exists* $\alpha > 0$ *such that*

$$a[v,v] \geqslant \alpha\|v\|^2_{\mathscr{V}}, \quad \forall v \in \mathscr{V}. \tag{1.24}$$ eqn:coercive

We notice that $\sup_{w\in\mathscr{V}} \frac{a[v,w]}{\|w\|_{\mathscr{V}}} \geqslant \frac{a[v,v]}{\|v\|_{\mathscr{V}}} \geqslant \alpha\|v\|_{\mathscr{V}}$, which implies the first inf-sup condition in Theorem 1.4. Hence, for any $f \in \mathscr{V}'$, the coercive variational problem (1.16) has a unique solution and the solution u is continuously depends on f, i.e., $\|u\|_{\mathscr{V}} \leqslant \alpha^{-1}\|f\|_{\mathscr{V}'}$. In this case, Theorem 1.4 is reduced to the well-known *Lax-Milgram theorem*.

Corollary 1.2 (Lax-Milgram theorem)**.** *Let* $a[\cdot,\cdot] : \mathscr{V} \times \mathscr{V} \mapsto \mathbb{R}$ *be a continuous bilinear form which satisfies the coercivity condition* (1.24)*. Then* (1.16) *has a unique solution* $u \in \mathscr{V}$ *for any* $f \in \mathscr{V}'$ *and* $\|u\|_{\mathscr{V}} \leqslant \alpha^{-1}\|f\|_{\mathscr{V}'}$.

rem:energynorm **Remark 1.6** (Energy norm)**.** If the bilinear form $a[\cdot,\cdot] : \mathscr{V} \times \mathscr{V} \mapsto \mathbb{R}$ is symmetric, then, apparently, it defines an inner product on $\mathscr{V}$. Its induced norm is also called the energy norm

$$|\!|\!|v|\!|\!| := a[v,v]^{1/2}.$$

Coercivity and continuity of the bilinear form $a[\cdot,\cdot]$ imply that

$$\alpha\|v\|^2_{\mathscr{V}} \leqslant |\!|\!|v|\!|\!|^2 \leqslant \|a[\cdot,\cdot]\|\,\|v\|^2_{\mathscr{V}} = \|\mathcal{A}\|_{\mathscr{L}(\mathscr{V};\mathscr{V})}\|v\|^2_{\mathscr{V}},$$

namely, the energy norm $|\!|\!|\cdot|\!|\!|$ is equivalent to the $\|\cdot\|_{\mathscr{V}}$-norm. We will denote the dual energy norm by $|\!|\!|\cdot|\!|\!|_*$

rem:well-posed **Remark 1.7** (Poisson is "well-conditioned")**.** We notice that the Poisson's equation is well-posed in the sense that $-\Delta : \mathscr{V} \mapsto \mathscr{V}'$ is an isomorphism with $\mathscr{V} = H_0^1(\Omega)$ and $\mathscr{V}' = H^{-1}(\Omega)$. Furthermore, it is well-conditioned, i.e., there exist constants α (coercivity constant) and C_a (continuity constant), such that

$$\alpha \|v\|_{\mathscr{V}}^2 \leqslant a[v,v] = \langle -\Delta v, v\rangle \leqslant C_a \|v\|_{\mathscr{V}}^2, \quad \forall v \in \mathscr{V}.$$

Hence we have the "condition number" of the Laplace operator is bounded

$$\kappa(-\Delta) = \|-\Delta\|_{\mathscr{L}(\mathscr{V};\mathscr{V}')} \cdot \|(-\Delta)^{-1}\|_{\mathscr{L}(\mathscr{V}';\mathscr{V})} \leqslant \frac{C_a}{\alpha}.$$

The problem here lies in that we are working on two different spaces $\mathscr{V}$ and $\mathscr{V}'$. If we consider $-\Delta : L^2(\Omega) \mapsto L^2(\Omega)$ instead, then we lost boundedness. More general theory has been developed in the seminar work [4].

High-frequency and locality

Consider the eigenvalue problem for one-dimensional Laplace operator with the homogenous Dirichlet boundary condition, i.e., $-u''(x) = \lambda u(x)$ for $x \in (0,1)$ and $u(0) = u(1) = 0$. It is easy to see that the eigenvalues and the corresponding eigenfunctions are

$$\lambda_k = (k\pi)^2 \qquad \text{and} \qquad u_k(x) = \sin(k\pi x), \quad k = 1, 2, \cdots$$

For other types of boundary conditions, the eigenvalues and eigenfunctions can be obtained as well. We notice that larger eigenvalues (larger k) correspond to eigenfunctions of higher frequency. Similar results can be expected for discrete problems which will be discussed later on.

An important observation comes from the analysis to the local problem

$$-u_\delta''(x) = f(x), \;\; x \in B_\delta := (x_0 - \delta, x_0 + \delta) \qquad \text{and} \qquad u_\delta(x_0 - \delta) = u_\delta(x_0 + \delta) = 0.$$

We can obtain the eigenfunctions of this local problem: $u_{\delta,k}(x) = \sin\big(\frac{k\pi}{2\delta}(x - x_0 + \delta)\big)$, $k = 1, 2, \cdots$. Define the error $e := u - u_\delta$ in B_δ. Hence e is harmonic in B_δ. It is easy to construct a cut-off function $\theta \in C_0^\infty(B_\delta)$, such that it satisfies the following conditions:

$$\text{(i) } \theta(x) > 0; \quad \text{(ii) } \theta(x) = 1, \; \forall x \in B_{\delta/2}; \quad \text{(iii) } |\theta'(x)| \leqslant \tfrac{C}{\delta}.$$

Thus we have

$$\begin{aligned} \int_{B_{\delta/2}} |e'(x)|^2 \, dx \;\leqslant\; \int_{B_\delta} \theta^2(x) \, |e'(x)|^2 \, dx \;&=\; -\int_{B_\delta} \Big((\theta^2)' e' + \theta^2 e'' \Big) e \, dx \\ &\leqslant\; \frac{2C}{\delta} \int_{B_\delta} |\theta e' e| \, dx \;\leqslant\; \frac{2C}{\delta} \Big(\int_{B_\delta} |\theta e'|^2 \, dx \Big)^{\frac{1}{2}} \Big(\int_{B_\delta} |e|^2 \, dx \Big)^{\frac{1}{2}}. \end{aligned}$$

The first and last inequalities immediately imply that

$$\Big(\int_{B_{\delta/2}} |e'(x)|^2\,dx\Big)^{\frac{1}{2}} \leqslant \Big(\int_{B_\delta} \theta^2(x)|e'(x)|^2\,dx\Big)^{\frac{1}{2}} \leqslant \frac{2C}{\delta}\Big(\int_{B_\delta} |e|^2\,dx\Big)^{\frac{1}{2}}. \tag{1.25}$$

eqn:Harnack

If we plug in the eigenfunctions $u_{\delta,k}$ to this inequality, we can see that

$$\frac{k\pi}{2\delta} \leqslant \frac{2C}{\delta},$$

which suggests only low-frequency components are left in the error function e and oscillating components in the distance δ are accurately captured.

m:GeoHighFreq

Remark 1.8 (Geometric high-frequencies)**.** This simple result implies that the high-frequency part of u can be estimated very well by the local solution u_δ for elliptic problems. Singularity is a particular form of high-frequency. In the numerical treatments, many forms of singularity can be resolved through local mesh refinement and the reason why this type of methods is able to work is because of such local behavior of high frequencies. Motived by (1.25), we can define geometric high-frequency functions u_k as those with large $\|\nabla u_k\|_{0,\Omega}/\|u_k\|_{0,\Omega}$ ratio.

1.2 Discretization methods

scretizations

Discretization concerns the process of transferring continuous functions, models, or equations into their discrete counterparts. This process is usually carried out as the first step toward making them suitable for numerical evaluation and implementation on modern computers.

Let $\Omega \in \mathbb{R}^d$ be an open domain and $f \in L^2(\Omega)$. We consider the following model problem

$$\begin{cases} -\Delta u = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{cases}$$

Many discretization methods have been developed, such as finite difference (FD) and the finite element (FE) methods, each with specific approaches to discretization. After discretization, we usually end up with a linear algebraic system of equations

$$A\vec{u} = \vec{f}. \tag{1.26}$$

eqn:linear0

Finite difference method

In one-dimensional case, without loss of generality, we can assume $\Omega = (0,1)$ and the domain is sub-divided into $N+1$ equally spaced pieces. So we get a uniform mesh with meshsize $h = \frac{1}{N+1}$; see the following figure for illustration.

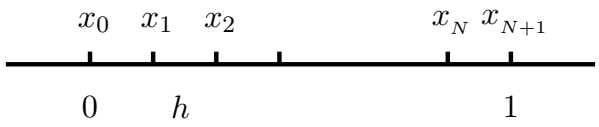


Figure 1.2: Uniform mesh in 1D.

fig:mesh1d

Using the Taylor's expansion, we can easily obtain that

$$\begin{aligned} u''(x_i) &= \frac{1}{h}\Big[u'(x_{i+\frac{1}{2}}) - u'(x_{i-\frac{1}{2}})\Big] + O(h^2) \\ &= \frac{1}{h^2}\Big[u(x_{i-1}) - 2u(x_i) + u(x_{i+1})\Big] + O(h^2). \end{aligned}$$

Let $u_i \approx u(x_i)$ be an approximate solution. Then the FD discretization of the Poisson's equation is

$$\frac{1}{h^2}\begin{pmatrix} 2 & -1 & & & \\ -1 & 2 & -1 & & \\ & \ddots & \ddots & \ddots & \\ & & \ddots & \ddots & -1 \\ & & & -1 & 2 \end{pmatrix}\begin{pmatrix} u_1 \\ u_2 \\ \vdots \\ u_{N-1} \\ u_N \end{pmatrix} = \begin{pmatrix} f(x_1) \\ f(x_2) \\ \vdots \\ f(x_{N-1}) \\ f(x_N) \end{pmatrix}. \tag{1.27}$$

eqn:1DFDA

That is to say,

$$A := \frac{1}{h^2}\,\text{tridiag}(-1, 2, -1) \qquad \text{and} \qquad \vec{f} := \Big(f_i\Big)_{i=1}^{N} = \Big(f(x_i)\Big)_{i=1}^{N}.$$

We need to solve the linear system $A\vec{u} = \vec{f}$ in order to obtain an approximate solution to the Poisson's equation. It is worth noticing that the coefficient matrix A is symmetric positive definite (SPD), sparse, as well as Toeplitz.

Remark 1.9 (An alternative form of the linear system)**.** Sometimes, it is more convenient (for implementation) to also include the boundary values in $\vec{u}$ and write the linear system as

$$\frac{1}{h^2}\begin{pmatrix} 1 & & & & & \\ -1 & 2 & -1 & & & \\ & \ddots & \ddots & \ddots & & \\ & & -1 & 2 & -1 & \\ & & & & 1 \end{pmatrix}\begin{pmatrix} u_0 \\ u_1 \\ \vdots \\ u_N \\ u_{N+1} \end{pmatrix} = \begin{pmatrix} 0 \\ f_1 \\ \vdots \\ f_N \\ 0 \end{pmatrix}.$$

Apparently this form is equivalent to the discrete problem above.

rem:FDevalue

Remark 1.10 (Eigenvalues of 1D FD problem)**.** For simplicity we now assume $h \equiv 1$. It is well-known (see HW 1.2) that the eigenvalues of $A := \text{tridiag}(-1, 2, -1)$ are

$$\lambda_k(A) = 2 - 2\cos\left(\frac{k\pi}{N+1}\right) = 4\sin^2\left(\frac{k\pi}{2(N+1)}\right)$$

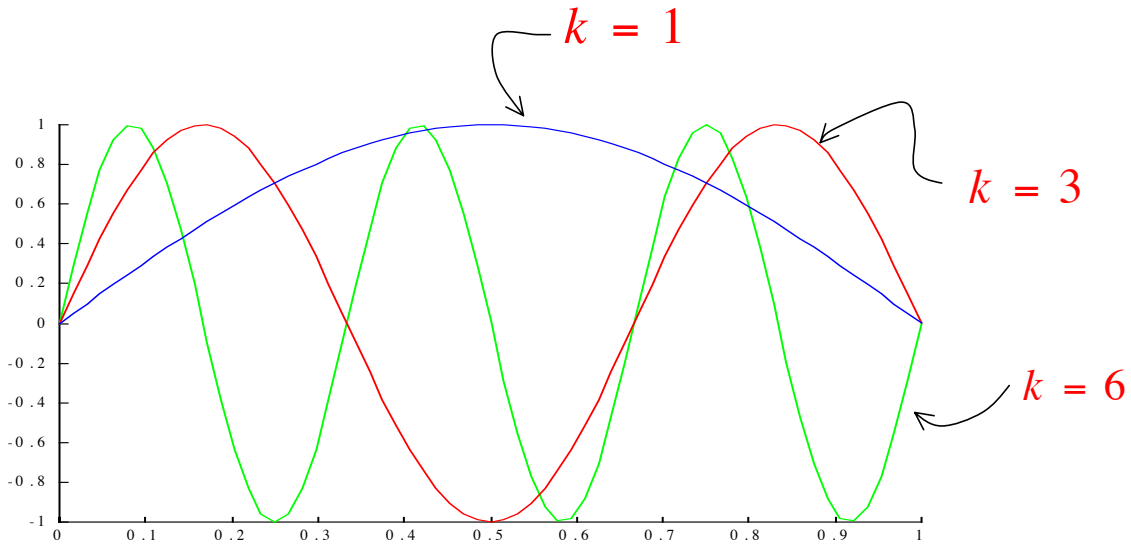
and the corresponding eigenvectors are

$$\vec{\xi}^k = \left(\xi_i^k\right)_{i=1}^N \in \mathbb{R}^N, \quad \text{with } \xi_i^k := \sin\left(\frac{ik\pi}{N+1}\right).$$

We note that the set of eigenvectors of A, $\vec{\xi}^k = \left(\xi_i^k\right)_{i=1}^N$, forms an orthogonal basis of $\mathbb{R}^N$. Therefore, any $\vec{\xi} \in \mathbb{R}^N$ can be expanded in terms of these eigenvectors:

$$\vec{\xi} = \sum_{k=1}^{N} \alpha_k \vec{\xi}^k.$$

This type of expansion is often called the discrete Fourier expansion. From Figure 1.3, we can easily see that the eigenvectors are "smooth" with small k and are "oscillatory" with large k. Hence the smoothness of $\vec{\xi}$ has a lot to do with the relative size of the coefficients α_k.



ig:eigenfct1D

Figure 1.3: Eigenvectors of 1D finite difference system for the Poisson's equation.

For two-dimensional problems, we can partition the domain uniformly in both x and y-directions into $n+1$ pieces ($N = n^2$). We denote $(x_i, y_j) = \left(\frac{i}{n+1}, \frac{j}{n+1}\right)$ and the Poisson's equation is discretize using the five-point stencil

$$\frac{1}{h^2}\Big[4u_{i,j} - \big(u_{i-1,j} + u_{i+1,j} + u_{i,j-1} + u_{i,j+1}\big)\Big] = f(x_i, y_j), \quad i,j = 1,\ldots,n.$$

Then we need to assign an order to the grid points in order to write the unknowns as a vector. There are many ways to order the unknowns for practical purposes. For simplicity, we use the

Lexicographic ordering, i.e., $p_{(j-1)n+i} := (x_i, y_j)$. Then we have

$$\frac{1}{h^2}\begin{pmatrix} A_1 & -I & & & \\ -I & A_2 & -I & & \\ & \ddots & \ddots & \ddots & \\ & & -I & A_{n-1} & -I \\ & & & -I & A_n \end{pmatrix}\begin{pmatrix} u_1 \\ u_2 \\ \vdots \\ \vdots \\ u_{N-1} \\ u_N \end{pmatrix} = \begin{pmatrix} f_1 \\ f_2 \\ \vdots \\ \vdots \\ f_{N-1} \\ f_N \end{pmatrix},$$

where the block diagonal matrices $A_i := \text{tridiag}(-1, 4, -1)$, $(i = 1, \ldots, n)$ are tridiagonal. Define $C := \text{tridiag}(-1, 0, -1)$. Then it is clear that

$$A = \frac{1}{h^2}\,\text{tridiag}(-I, A_1, -I) = \frac{1}{h^2} I \otimes A_1 + \frac{1}{h^2} C \otimes I.$$

Remark 1.11 (Eigenvalues of the 2D FD problem)**.** Again we assume $h \equiv 1$. Similar to the 1D problem, we can get the eigenvalues

$$\lambda_{i,j}(A) = 4 - 2\cos\frac{i\pi}{n+1} - 2\cos\frac{j\pi}{n+1} = 4\sin^2\frac{i\pi}{2(n+1)} + 4\sin^2\frac{j\pi}{2(n+1)},$$

with eigenvectors

$$\vec{\xi}_{i,j} = \left(\sin\frac{ki\pi}{n+1}\sin\frac{lj\pi}{n+1}\right)_{k,l=1,\ldots,n}.$$

Remark 1.12 (Ordering)**.** The shape of the above coefficient matrix A depends on the ordering of degrees of freedom (DOFs). We will see that the ordering also affects the smoothing properties of smoothers and parallelization. Finding minimal bandwidth ordering is important for some methods, like direct solvers. But it is NP-hard.

Finite element method

Finite element method (FEM) is a Galerkin method that uses piecewise polynomial spaces for approximate test and trial function spaces. The readers are referred to [25, 37, 10, 20] for more detailed discussion on construction and error analysis of the standard finite element method.

The weak formulation of the model equation can be written as (see Example 1.1): Find $u \in H_0^1(\Omega)$, such that

$$\int_\Omega \nabla u \cdot \nabla v dx = \int_\Omega f v \; dx, \quad \forall v \in H_0^1(\Omega).$$

In 1D, it is easy to explain the main idea of finite element method. Let $\mathcal{P}_k(\tau)$ be the space of all polynomials of degree less than or equal to k on τ. Let

$$V = V_h := \big\{v \in C(\overline{\Omega}) \,:\, v \in \mathcal{P}_1(x_{i-1}, x_i),\; v(0) = v(1) = 0\big\}.$$

Now we can write the discrete variational problem as: Find $u_h \in V_h$, such that

$$a[u_h, v_h] = (f, v_h), \quad \forall v_h \in V_h.$$

Furthermore, we use nodal basis functions $\phi_i \in V_h$, i.e. $\phi_i(x_j) = \delta_{i,j}$. In this way, we can express a given function $u_h \in V_h$ as $u_h(x) = \sum_{j=1}^N u_j \phi_j(x)$. Hence we arrive at the following equation: For any $i = 1, \ldots, N$,

$$\sum_{j=1}^{N} a[\phi_j, \phi_i] u_j = (f, \phi_i) \qquad \text{or} \qquad \sum_j A_{i,j} u_j = f_i.$$

This is a system of algebraic linear equations

$$A\vec{u} = \vec{f}, \tag{1.28}$$ eqn:FEM1D

with $\big(A\big)_{i,j} = a_{i,j} := a[\phi_i, \phi_j]$, $\vec{u} := \big(u_i\big)_{i=1}^N$, and $\vec{f} = \big(f_i\big)_{i=1}^N := \Big(\langle f, \phi_i \rangle\Big)_{i=1}^N$.

If we use the uniform mesh in Figure 1.2, then we have (see HW 1.3) that

$$A := \frac{1}{h}\,\text{tridiag}(-1, 2, -1) \qquad \text{and} \qquad \vec{f} := \big(hf(x_i)\big)_{i=1}^N.$$

Upon solving this finite-dimensional problem, we obtain a discrete approximation u_h. The finite element method has several appealing properties and it will be the main underlying discretization used in this lecture; see §3.1 for more details.

rem:ill **Remark 1.13** (Discrete Poisson's equation is ill-conditioned)**.** Remark 1.7 has shown that the Poisson's equation has a bounded condition number. On the other hand, the discrete problems from FD and FE are both ill-conditioned if meshsize h is small. Later on, we will see that this will cause problems for many iterative methods. The convergence rate of these methods usually depend on the spectrum of the coefficient matrix A.

rem:sim **Remark 1.14** (A useful notation)**.** We use some notations introduced by Xu [55]. The notation $a \lesssim b$ means: there is a generic constant C independent of meshsize h, such that $a \leqslant Cb$. Similarly, we can define "$\gtrsim$" and "$\cong$". This is important because, in our future discussions, we would like to construct solvers/preconditioners that yield convergence rate independent of meshsize h.

1.3 Simple iterative solvers

sec:simple

There are many different approaches for solving the linear algebraic equations results from the finite difference, finite element, and other discretizations for the Poisson's equation. For example, sparse direct solvers, FFT, and iterative methods. We only discuss iterative solvers in this lecture.

Some examples

Now we give a few well-known examples of simple iterative methods. Consider the linear system $A\vec{u} = \vec{f}$. Assume the coefficient matrix $A \in \mathbb{R}^{N\times N}$ can be partitioned as $A = L + D + U$, where $L, D, U \in \mathbb{R}^{N\times N}$ are the lower triangular, diagonal, and upper triangular parts of A, respectively (the rest is set to be zero).

example:R **Example 1.2** (Richardson method)**.** The simplest iterative method for solving $A\vec{u} = \vec{f}$ might be the Richardson method

$$\vec{u}^{\,\text{new}} = \vec{u}^{\,\text{old}} + \omega\Big(\vec{f} - A\vec{u}^{\,\text{old}}\Big). \tag{1.29}$$ eqn:richardso

We can choose an optimal weight ω to improve performance of this method. □

example:J **Example 1.3** (Weighted Jacobi method)**.** The weighted or damped Jacobi method can be written as

$$\vec{u}^{\,\text{new}} = \vec{u}^{\,\text{old}} + \omega D^{-1}(\vec{f} - A\vec{u}^{\,\text{old}}). \tag{1.30}$$ eqn:dJacobi

This method solves one equation for one variable at a time, simultaneously. Apparently, it is a generalization of the above Richardson method. If $\omega = 1$, then we arrive at the standard Jacobi method. □

example:GS **Example 1.4** (Gauss–Seidel method)**.** The Gauss–Seidel (G-S) method can be written as

$$\vec{u}^{\,\text{new}} = \vec{u}^{\,\text{old}} + (D + L)^{-1}(\vec{f} - A\vec{u}^{\,\text{old}}).$$

We rewrite this method as

$$(D + L)\vec{u}^{\,\text{new}} = (D + L)\vec{u}^{\,\text{old}} + (\vec{f} - A\vec{u}^{\,\text{old}}) = \vec{f} - U\vec{u}^{\,\text{old}}.$$

Thus we have

$$\vec{u}^{\,\text{new}} = \vec{u}^{\,\text{old}} + D^{-1}\Big(\vec{f} - L\vec{u}^{\,\text{new}} - (D + U)\vec{u}^{\,\text{old}}\Big). \tag{1.31}$$ eqn:GS

Compared with the Jacobi method (1.30) ($\omega = 1$), the G-S method uses the most updated solution in each iteration instead of the previous iteration. □

Example 1.5 (Successive over-relaxation method)**.** The successive over-relaxation (SOR) method can be written as

$$(D + \omega L)\vec{u}^{\,\text{new}} = \omega\vec{f} - \Big(\omega U + (\omega - 1)D\Big)\vec{u}^{\,\text{old}}. \tag{1.32}$$ eqn:SOR

The weight ω is usually in $(1, 2)$. If $\omega = 1$, then it is the G-S method. □

A simple observation

Many simple iterative methods exhibit different rates of convergence for short and long wavelength error components, suggesting these different scales should be treated differently. Let $\lambda_{\max}$ and $\lambda_{\min}$ be the largest eigenvalue and the smallest eigenvalue of A, respectively, and $\vec{\xi}_{\max}$ and $\vec{\xi}_{\min}$ be the corresponding eigenvectors. One interesting observation many people made is: When we use the weighted Jacobi method (1.30) with weight $\omega = 2/3$ to solve the problem $A\vec{u} = \vec{0}$ with the initial guess just equal to $\vec{\xi}_{\max}$, the convergence is very fast. On the other hand, if the weighted Jacobi iteration is used to solve the same equation but with a different initial guess $\vec{\xi}_{\min}$, the convergence becomes slow. See Figure 1.4 for a demonstration.

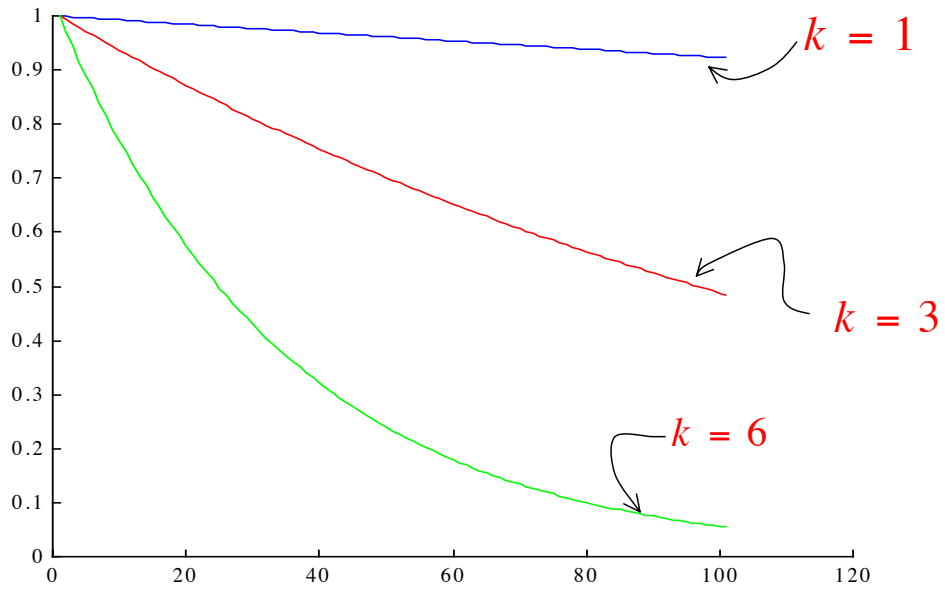


fig:errorJ

Figure 1.4: Error decay in $\|\cdot\|_\infty$-norm for weighted Jacobi method with initial guess $\vec{\xi}^k$.

Note that the reason which causes this difference mainly relies on the fact that the error in the first problem (corresponding to $\vec{\xi}_{\max}$) is oscillatory or of high frequency but the error in the second problem (corresponding to $\vec{\xi}_{\min}$) is smooth or of low frequency. This makes one speculate that the weighted Jacobi method can damp the high frequency part of the error rather quickly, but slowly for the low frequency part; see Remark 1.8.

Smoothing effect ⋆

In view of Remark 1.10, based on the understanding of the relation between the smoothness and the size of Fourier coefficients, we can analyze the smoothing property using the discrete Fourier expansion. Let $\vec{u}$ be the exact solution of the 1D FD problem on uniform grids and $\vec{u}^{(m)}$

the result of m-th iteration from the damped Jacobi method (or equivalently in this case, the Richardson method). Then

$$\vec{u} - \vec{u}^{(m)} = (I - \omega A)(\vec{u} - \vec{u}^{(m-1)}) = \cdots = (I - \omega A)^m(\vec{u} - \vec{u}^{(0)}).$$

It is straightforward to see that

$$\lambda_k(I - \omega A) = 1 - \omega\lambda_k(A) = 1 - 4\omega \sin^2\left(\frac{k\pi}{2(N+1)}\right).$$

Notice that $\lambda_k(I - \omega A)$ can be viewed as the damping factor for error components corresponding to Fourier mode k; see Remark 1.10. We would like to choose ω such that λ_k's are small.

Consider the Fourier expansion of the initial error:

$$\vec{u} - \vec{u}^{(0)} = \sum_{k=1}^{N} \alpha_k \vec{\xi}^k.$$

Then

$$\vec{u} - \vec{u}^{(m)} = \sum_{k=1}^{N} \alpha_k (I - \omega A)^m \vec{\xi}^k.$$

Note that, for any polynomial p, we have $p(A)\vec{\xi}^k = p(\lambda_k)\vec{\xi}^k$. By choosing $\omega = \frac{1}{4} \approx \frac{1}{\lambda_{\max}(A)}$, we obtain

$$\vec{u} - \vec{u}^{(m)} = \sum_{k=1}^{N} \alpha_k (1 - \omega\lambda_k)^m \vec{\xi}^k = \sum_{k=1}^{N} \alpha_k^{(m)} \vec{\xi}^k,$$

where

$$\alpha_k^{(m)} = \left(1 - \sin^2 \frac{k\pi}{2(N+1)}\right)^m \alpha_k.$$

The above equation implies

$$\alpha_k^{(m)} = \alpha_k \sin^{2m}\left(\frac{N-k+1}{N+1}\frac{\pi}{2}\right) \leqslant \alpha_k \left(\frac{N-k+1}{N+1}\frac{\pi}{2}\right)^{2m},$$

which approaches to 0 very rapidly as $m \to \infty$, if k is close to N (high-frequencies). This means that high frequency error can be damped very quickly. This simple analysis justifies the smoothing property we observed in the beginning of this section.

We can apply the same analysis to the Jacobi method as well and the Fourier coefficient in front of the highest frequency is as follows:

$$\alpha_N^{(m)} = \left(1 - 2\sin^2 \frac{N\pi}{2(N+1)}\right)^m \alpha_N = \cos^m\left(\frac{N\pi}{N+1}\right)\alpha_N \sim (-1)^m \left(1 - \frac{\pi^2}{2(N+1)^2}\right)^m \alpha_N.$$

This suggests that the regular Jacobi method might not have a smoothing property and should not be used a smoother in general.

1.4 Multigrid method in 1D

sec:GMG1D

In this section, we first give a simple motivation and sneak-peak of the well-known multigrid method, which is a representing example of multilevel iterative methods. The observations of this section will be helpful for our later discussions; see the famous tutorial by Briggs et al. [23] for a quick introduction to the multigrid methods. Consider the finite difference scheme (1.27) for the Poisson's equation in 1D, namely

$$A\vec{u} = \vec{f} \qquad \text{with } A = \frac{1}{h^2}\,\text{tridiag}(-1, 2, -1), \quad f_i = f(x_i).$$

Nested grids

Multigrid (MG) methods are a group of algorithms for solving partial differential equations using a hierarchy of discretizations. They are very useful in problems exhibiting multiple scales of behavior. In this section, we introduce the simplest multigrid method in 1D.

Suppose there are a hierarchy of $L+1$ grids with mesh sizes $h_l = (\frac{1}{2})^{l+1}$ $(l = 0, 1, \ldots, L)$; see Figure 1.5. It is clear that

$$h_0 > h_1 > h_2 > \cdots > h_L =: h$$

and $N = 2^{L+1} - 1$. We call level L the finest level and level 0 the coarsest level.

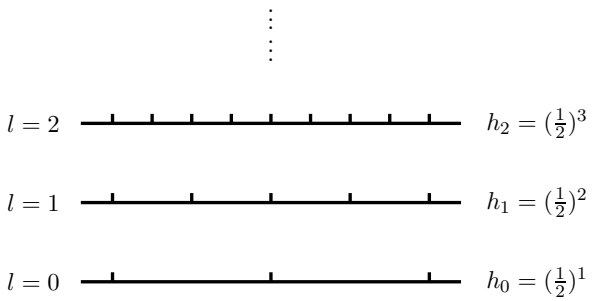


Figure 1.5: Hierarchical grids for 1D multigrid method.

fig:hiergrid

Smoothers

We consider how to approximate the solution on each level using some local relaxation method. Assume the 1D Poisson's equation is discretized using the finite difference scheme discussed in the previous section. Then, on each level, we have a linear system of equations

$$A_l\vec{u}_l = \vec{f}_l \qquad \text{with } A_l = h_l^{-2}\,\text{tridiag}(-1, 2, -1).$$

For each of these equations, we can apply a damped Jacobi method (with damping factor 1/2)

$$\vec{u}_l^{(m+1)} = \vec{u}_l^{(m)} + \frac{1}{2} D_l^{-1}\left(\vec{f}_l - A_l \vec{u}_l^{(m)}\right) \tag{1.33}$$ eqn:djacobi

to obtain approximate solutions. This method is usually referred as a local relaxation or smoother, which will be discussed later in this lecture note.

Prolongation and restriction

Another important component of a multigrid method is to define the transfer operators between different levels. In the 1D case, the transfer operators can be easily given; see Figure 1.6. In another word, we can also write the transfer operators in the matrix form, i.e.,

$$R_{l,l-1} := \frac{1}{4}\begin{pmatrix} 1 & 2 & 1 & & & \\ & & 1 & 2 & 1 & \\ & & & & & \ddots \end{pmatrix} \quad \text{and} \quad P_{l-1,l} := \frac{1}{2}\begin{pmatrix} 1 & & \\ 2 & & \\ 1 & 1 & \\ & 2 & \\ & 1 & \\ & & \ddots \end{pmatrix}. \tag{1.34}$$ eqn:RandP

We notice that $R = \frac{1}{2}P^T$. It is straight-forward to check that the coefficient matrices of two

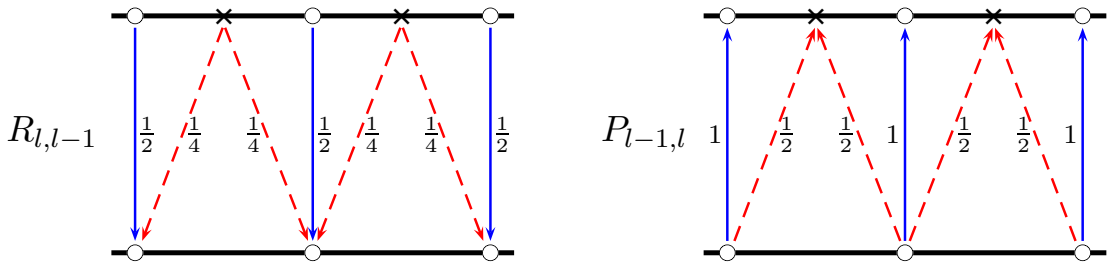


Figure 1.6: Transfer operators between two consecutive levels (Left: restriction operator; right: prolongation operator). ig:transfer1d

consecutive levels satisfy

$$A_{l-1} = R_{l,l-1} A_l P_{l-1,l}.$$

Multigrid algorithm

Now we are ready to give one step of the multigrid algorithm.

alg:1DGMG **Algorithm 1.1** (One iteration of multigrid method)**.** $\vec{u}_l = MG(l, \vec{f}_l, \vec{u}_l)$

(i) **Pre-smoothing**: $\vec{u}_l \leftarrow \vec{u}_l + \frac{1}{2} D_l^{-1}\left(\vec{f}_l - A_l \vec{u}_l\right)$

(ii) **Restriction**: $\vec{r}_{l-1} \leftarrow R_{l,l-1}\big(\vec{f}_l - A_l\vec{u}_l\big)$

(iii) **Coarse-grid correction**: $\vec{e}_{l-1} \leftarrow MG(l-1, \vec{r}_{l-1}, \vec{0}_{l-1})$

(iv) **Prolongation**: $\vec{u}_l \leftarrow \vec{u}_l + P_{l-1,l}\vec{e}_{l-1}$

(v) **Post-smoothing**: $\vec{u}_l \leftarrow \vec{u}_l + \frac{1}{2}D_l^{-1}\big(\vec{f}_l - A_l\vec{u}_l\big)$

Remark 1.15 (Coarse-grid correction)**.** Suppose that there is an approximate solution $\vec{u}^{(m)}$. Then we have

$$A\big(\vec{u} - \vec{u}^{(m)}\big) = \vec{r}^{(m)} := \vec{f} - A\vec{u}^{(m)}$$

and the *error equation* can be written

$$A\vec{e}^{(m)} = \vec{r}^{(m)}. \tag{1.35}$$ eqn:error

If we get $\vec{e}^{(m)}$ or its approximation, we can just update the iterative solution by $\vec{u}^{(m+1)} = \vec{u}^{(m)} + \vec{e}^{(m)}$ to obtain a better approximation of $\vec{u}$. This explains the steps (iii) and (iv) in the above algorithm.

Remark 1.16 (Coarsest-level solver)**.** It is clear that, in our setting, the solution at level $l = 0$ is trivial to obtain. In general, we can apply a direct or iterative solver to solve the coarsest-level problem, which is relatively cheap. Sometimes, we have singular problems at the coarsest level, which need to be handled with care.

Algorithm 1.1 is one iteration of the multigrid method. We can iterate until the approximation is "satisfactory". For example, we iterate until the relative residual $\|\vec{r}\|_0/\ |\vec{f}\|_0$ is less than 10^{-6}; we will discuss stopping criteria later in this lecture. This multigrid algorithm is easy to implement; see HW 1.5. In Table 1.1, we give the numerical results of Algorithm 1.1 for the 1D Poisson's equation. From the table, we find that, unlike the classical Jacobi and G-S methods, this multigrid method converges uniformly with respect to the meshsize h. This is, of course, a very desirable feature of the multilevel iterative methods, which will be investigated in more details in this lecture.

Now it is natural to have a few questions on such multilevel methods:

- When does the multigrid method converge?
- How fast the method converges?
- How to find a smoother?
- Why the matrices R and P are given as (1.34)?
- How to generalize the method to other problems?

#Levels	#DOF	#Iter	Contract factor
5	31	4	0.0257
6	63	4	0.0259
7	127	4	0.0260
8	255	4	0.0260
9	511	4	0.0261
10	1023	4	0.0262

tab:1DGMG

Table 1.1: Convergence behavior of 1D geometric multigrid method.

1.5 Tutorial of FASP ⋆

All the numerical examples in this lecture are done using the Fast Auxiliary Space Preconditioning (FASP) package. The FASP package provides C source files[1] to build a library of iterative solvers and preconditioners for the solution of large-scale linear systems of equations. The components of the FASP basic library include several ready-to-use, modern, and efficient iterative solvers used in applications ranging from simple examples of discretized scalar partial differential equations (PDEs) to numerical simulations of complex, multicomponent physical systems.

The main components of the FASP basic library are:

- Basic linear iterative methods;
- Standard Krylov subspace methods;
- Geometric and Algebraic Multigrid (G/AMG) methods;
- Incomplete factorization methods.

The FASP distribution also includes several examples for solving simple benchmark problems. The basic (kernel) FASP distribution is open-source and is licensed under GNU Lesser General Public License or LGPL. Other distributions may have different licensing (contact the developer team for details on this). The most updated version of FASP can be downloaded directly from

http://fasp.sf.net/download/faspsolver.zip

To build the FASP library for these operating systems. Open a terminal window, where you can issue commands from the command line and do the following: (1) go to the main FASP directory (we will refer to it as `$(faspsolver)` from now on); (2) modify the "*FASP.mk.example*" file to math your system and save it as "FASP.mk"; (3) then execute:

```
> make config
> make install
```

[1]The code is C99 (ISO/IEC 9899:1999) compatible.

These two commands build the FASP library/header files. By default, it installs the library in `$(faspsolver)/lib` and the header files in `$(faspsolver)/include`. It also creates a file `$(faspsolver)/Config.mk` which contains few of the configuration variables and can be loaded by external project Makefiles. If you do not have "FASP.mk" present in the current directory, default settings will be used for building and installation FASP.

Now, if you would like to try some of the examples that come with FASP, you can build the "tutorial" target and try out the tutorial examples:

```
> make tutorial
```

Equivalently, you may also build the test suite and the tutorial examples by using the "local" Makefile in `$(faspsolver)/tutorial`.

```
> make -C tutorial
```

For more information, we refer to the user's guide and reference manual of FASP[2] for technical details on the usage and implementation of FASP. Since FASP is under heavy development, please use this guide with caution because the code might have been changed before this document is updated.

1.6 Homework problems

lutionPoisson **HW 1.1.** Prove the uniqueness of the Poisson's equation. Hint: You can argue by the maximum principle or the energy method.

hw:Evalue1DFD **HW 1.2.** Prove the eigenvalues and eigenvectors of tridiag$(b, a, b) \in \mathbb{R}^{N\times N}$ are

$$\lambda_k = a - 2b\cos\Big(\frac{k\pi}{N+1}\Big) \quad \text{and} \quad \vec{\xi}^k = \Big(\sin\big(\frac{k\pi}{N+1}\big), \ldots, \sin\big(\frac{Nk\pi}{N+1}\big)\Big)^T,$$

respectively. Apply this result to give eigenvalues of the 1D FD matrix A. What are the eigenvalues of tridiag$(b, a, c) \in \mathbb{R}^{N\times N}$?

w:StiffnessFE **HW 1.3.** Derive the finite element stiffness matrix for 1D Poisson's equation with homogenous Dirichlet boundary condition using a uniform mesh.

hw:FDFEHeat **HW 1.4.** Derive 1D FD and FE discretizations for the heat equation (1.6) using the backward Euler method for time discretization.

hw:GMGcode **HW 1.5.** Implementation of multigrid in 1D using Matlab, C, Fortran, or Python. Try to study the efficiency of your implementation.

[2] Available online at http://fasp.sourceforge.net. It is also available in "faspsolver/doc/".

Chapter 2

Iterative Solvers and Preconditioners

ch:iterative

The term "iterative method" refers to a wide range of numerical techniques that use successive approximations $\{u^{(m)}\}$ for the exact solution u of a certain problem. In this chapter, we will discuss two types of iterative methods: (1) Stationary iterative method, which performs in each iteration the same operations on the current iteration; (2) Nonstationary iterative method, which has iteration-dependent operations. Stationary methods are simple to understand and implement, but usually not very effective. On the other hand, nonstationary methods are a relatively recent development; their analysis is usually more difficult.

2.1 Stationary linear iterative methods

ec:stationary

In this section, we discuss stationary iterative methods; typical examples include the Jacobi method and the Gauss–Seidel method. We will discuss why they are not efficient in general but still widely used. Let V be a finite-dimensional linear vector space, $\mathcal{A} : V \mapsto V$ be a non-singular linear operator, and $f \in V$. We would like to find a $u \in V$, such that

$$\mathcal{A}u = f. \tag{2.1}$$

eqn:linear

For example, in the finite difference context discussed in §1.2, $V = \mathbb{R}^N$ and the linear operator $\mathcal{A}$ becomes a matrix A. We just need to solve a system of linear equations: Find $\vec{u} \in \mathbb{R}^N$, such that

$$A\vec{u} = \vec{f}. \tag{2.2}$$

eqn:linear2

We will discuss the linear systems in both operator and matrix representations. We note that the notation here in (2.2) might be slightly different when we discuss the discrete problems arising from the finite element discretizations.

Remark 2.1 (More general setting)**.** In fact, we can consider iterative methods in a more general setting. For example, let V be a finite-dimensional Hilbert space, V' be its dual, and

$\mathcal{A} : V \mapsto V'$ be a linear operator and $f \in V'$. A significant part of this lecture can be generalized to such a setting easily.

Preliminaries

The most-used inner product in this lecture is the Euclidian inner product $(u, v) := \int_\Omega uv\, dx$; and if $V = \mathbb{R}^N$, $(u, v) := \sum_{i=1}^N u_i v_i$. Once we have the inner product, we can define the concept of transpose and symmetry on the Hilbert space V. Define the *adjoint* operator (transpose) of the linear operator $\mathcal{A}$ as $\mathcal{A}^T : V \mapsto V$, such that

$$(\mathcal{A}^T u, v) := (u, \mathcal{A}v), \quad \forall u, v \in V.$$

m:nonsingular **Remark 2.2** (Non-singularity)**.** We denote the null space and the range of $\mathcal{A}$ as

$$\text{null}(\mathcal{A}) := \{v \in V \,:\, \mathcal{A}v = 0\}, \tag{2.3}$$

$$\text{range}(\mathcal{A}) := \{u = \mathcal{A}v \,:\, v \in V\}. \tag{2.4}$$

We have $\text{null}(\mathcal{A}^T)^\perp = \overline{\text{range}(\mathcal{A})}$ and $\text{null}(\mathcal{A}^T) = \text{range}(\mathcal{A})^\perp$. If $\text{null}(\mathcal{A}) = \{0\}$, then $\mathcal{A}$ is *injective* or one-to-one. Apparently, $\mathcal{A} : V \mapsto \text{range}(\mathcal{A})$ is *surjective* or onto. If we consider a symmetric operator $\mathcal{A} : \text{null}(\mathcal{A})^\perp \mapsto \text{range}(\mathcal{A})$, then $\mathcal{A}$ is always *non-singular*.

A linear bounded operator $\mathcal{A}$ on V is *symmetric* if and only if

$$(\mathcal{A}u, v) = (u, \mathcal{A}v), \quad \forall u, v \in \text{domain}(\mathcal{A}) \subseteq V.$$

If $\mathcal{A}$ is densely defined and $\mathcal{A}^T = \mathcal{A}$, then $\mathcal{A}$ is called *self-adjoint*.

Remark 2.3 (Symmetric and self-adjoint operators)**.** A symmetric operator $\mathcal{A}$ is self-adjoint if $\text{domain}(\mathcal{A}) = V$. The difference between symmetric and self-adjoint operators is technical; see [61] for details.

The set of eigenvalues of $\mathcal{A}$ is called the spectrum, denoted as $\sigma(\mathcal{A})$. The spectrum of any bounded symmetric matrix is real, i.e., all eigenvalues are real, although a symmetric operator may have no eigenvalues[1]. We define the spectral radius $\rho(\mathcal{A}) := \sup\{|\lambda| \,:\, \lambda \in \sigma(\mathcal{A})\}$.

An important class of operators for this chapter is symmetric positive definite (SPD) operators. An operator $\mathcal{A}$ is called SPD if and only if $\mathcal{A}$ is symmetric and $(\mathcal{A}v, v) > 0$, for any $v \in V \backslash \{0\}$. Since $\mathcal{A}$ is SPD, all of its eigenvalues are positive; furthermore,

$$\lambda_{\min}(\mathcal{A}) = \min_{v \in V\backslash\{0\}} \frac{(\mathcal{A}v, v)}{\|v\|^2} \qquad \text{and} \qquad \lambda_{\max}(\mathcal{A}) = \max_{v \in V\backslash\{0\}} \frac{(\mathcal{A}v, v)}{\|v\|^2}.$$

[1] A bounded linear operator on an infinite-dimensional Hilbert space might not have any eigenvalues.

We define the condition number $\kappa(\mathcal{A}) := \frac{\lambda_{\max}(\mathcal{A})}{\lambda_{\min}(\mathcal{A})}$, which is more convenient, compared with spectrum, to characterize convergence rate of iterative methods. For the indefinite case, we can use

$$\kappa(\mathcal{A}) := \frac{\sup_{\lambda\in\sigma(\mathcal{A})}|\lambda|}{\inf_{\lambda\in\sigma(\mathcal{A})}|\lambda|}.$$

More generally, for an isomorphic mapping $\mathcal{A}\in\mathscr{L}(V;V)$, we can define

$$\kappa(\mathcal{A}) := \|\mathcal{A}\|_{\mathscr{L}(V;V)}\|\mathcal{A}^{-1}\|_{\mathscr{L}(V;V)}.$$

And all these definitions are consistent for symmetric positive definite problems. If $\mathcal{A}$ is a SPD operator, it induces a new inner product, which will be used heavily in our later discussions

$$(u,v)_{\mathcal{A}} := (\mathcal{A}u,v) \qquad \forall u,v\in V. \tag{2.5}$$ eqn:A-inner

It is easy to check $(\cdot,\cdot)_{\mathcal{A}}$ is an inner product on V. For any linear bounded operator $\mathcal{B}: V\mapsto V$, we can define two transposes with respect to the inner products $(\cdot,\cdot)$ and $(\cdot,\cdot)_{\mathcal{A}}$, respectively; namely,

$$\begin{aligned}(\mathcal{B}^T u,v) &= (u,\mathcal{B}v),\\ (\mathcal{B}^* u,v)_{\mathcal{A}} &= (u,\mathcal{B}v)_{\mathcal{A}}.\end{aligned}$$

By the above definitions, it is clear (see HW 2.1) that

$$\mathcal{B}^* = \mathcal{A}^{-1}\mathcal{B}^T\mathcal{A}. \tag{2.6}$$ eqn:BtBstar

Remark 2.4 (Induced norms)**.** The inner products define above also induce norms on V by $\|v\| := (v,v)^{\frac{1}{2}}$ and $\|v\|_{\mathcal{A}} := (v,v)_{\mathcal{A}}^{\frac{1}{2}}$. These, in turn, define the operator norms for $\mathcal{B}: V\mapsto V$, i.e.,

$$\|\mathcal{B}\| := \sup_{v\in V\setminus\{0\}}\frac{\|\mathcal{B}v\|}{\|v\|} \qquad \text{and} \qquad \|\mathcal{B}\|_{\mathcal{A}} := \sup_{v\in V\setminus\{0\}}\frac{\|\mathcal{B}v\|_{\mathcal{A}}}{\|v\|_{\mathcal{A}}}.$$

It is well-known that, for any consistent norm $\|\cdot\|$, we have $\rho(\mathcal{B})\leqslant\|\mathcal{B}\|$. Furthermore, we have the following results:

prop:rho **Proposition 2.1** (Spectral radius and norm)**.** *Suppose V is Hilbert space with an inner product $(\cdot,\cdot)$ and induced norm $\|\cdot\|$. If $\mathcal{A}: V\mapsto V$ is a bounded linear operator, then*

$$\rho(\mathcal{A}) = \lim_{m\to+\infty}\|\mathcal{A}^m\|^{\frac{1}{m}}.$$

Moreover, if $\mathcal{A}$ is self-adjoint, then $\rho(\mathcal{A}) = \|\mathcal{A}\|$.

From this general functional analysis result, we can immediately obtain the following relations:

lem:sym **Lemma 2.1** (Spectral radius of self-adjoint operators)**.** *If $\mathcal{B}^T = \mathcal{B}$, then $\rho(\mathcal{B}) = \|\mathcal{B}\|$. Similarly, if $\mathcal{B}^* = \mathcal{B}$, then $\rho(\mathcal{B}) = \|\mathcal{B}\|_{\mathcal{A}}$.*

Symmetry is a concept with respect to the underlying inner product. In this chapter, we always refers to the $(\cdot,\cdot)$-inner product for symmetry. By definition, $(\mathcal{B}\mathcal{A})^* = \mathcal{B}^T\mathcal{A}$. If $\mathcal{B}^T = \mathcal{B}$, we do not necessarily have $(\mathcal{B}\mathcal{A})^T = \mathcal{B}\mathcal{A}$. However, we have a key identity:

$$(\mathcal{B}\mathcal{A})^* = \mathcal{B}^T\mathcal{A} = \mathcal{B}\mathcal{A};$$

see HW 2.2 for the first equality. So Lemma 2.1 implies the following identity:

$$\rho(\mathcal{I} - \mathcal{B}\mathcal{A}) = \|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}},$$

where $\mathcal{I} : V \mapsto V$ is the identity operator.

Stationary iterative methods

A linear stationary iterative method (one iteration) can be expressed in the following general form:

alg:iter **Algorithm 2.1** (Stationary iterative method)**.** $u^{\text{new}} = ITER(u^{\text{old}})$

(i) **Form residual**: $r = f - \mathcal{A}u^{\text{old}}$

(ii) **Solve error equation**: $\mathcal{A}e = r$ by $\hat{e} = \mathcal{B}r$

(iii) **Correct iteration**: $u^{\text{new}} = u^{\text{old}} + \hat{e}$

That is to say, the new iteration is obtained by computing

$$u^{\text{new}} = u^{\text{old}} + \mathcal{B}(f - \mathcal{A}u^{\text{old}}). \tag{2.7}$$ eqn:iter

Notice that each iteration only depends on the previous approximate solution u^{old} and does not involve any information of the older iterations; in each iteration, it basically performs the same operations over and over again. It is easy to see that

$$u - u^{(m)} = (\mathcal{I} - \mathcal{B}\mathcal{A})\big(u - u^{(m-1)}\big) = \cdots = (\mathcal{I} - \mathcal{B}\mathcal{A})^m\big(u - u^{(0)}\big) = \mathcal{E}^m\big(u - u^{(0)}\big),$$

where the operator $\mathcal{E} := \mathcal{I} - \mathcal{B}\mathcal{A}$ is called the *error propagation operator* and $\mathcal{B}$ is called the *iterator*.

We can get the following simple convergence theorem.

thm:rho **Theorem 2.1** (Convergence of Algorithm 2.1)**.** *The Algorithm 2.1 converges for any initial guess if the spectral radius $\rho(\mathcal{I} - \mathcal{B}\mathcal{A}) < 1$, which is equivalent to $\lim_{m\to+\infty}(\mathcal{I} - \mathcal{B}\mathcal{A})^m = 0$. The converse direction is also true.*

This convergence result is simple but difficult to apply. More importantly, it does not provide any information on how fast the convergence could be if the algorithm converges; see the following example for further explanation.

Example 2.1 (Spectral radius and convergence speed)**.** Suppose we have an iterative method with an error propagation matrix

$$E := \begin{pmatrix} 0 & 1 & \cdots & 0 \\ \vdots & \ddots & \ddots & \vdots \\ \vdots & & \ddots & 1 \\ 0 & \cdots & \cdots & 0 \end{pmatrix} \in \mathbb{R}^{N\times N}$$

and the initial error is $\vec{e}^{(0)} := \vec{u} - \vec{u}^{(0)} = (0, \ldots, 0, 1)^T \in \mathbb{R}^N$. Notice that, in this example, $\rho(E) \equiv 0$. However, if applying this error propagation matrix to form a sequence of approximations, we will get

$$\|\vec{e}^{(0)}\|_2 = \|\vec{e}^{(1)}\|_2 = \cdots = \|\vec{e}^{(N-1)}\|_2 = 1 \qquad \text{and} \qquad \|\vec{e}^{(N)}\|_2 = 0.$$

The main question now is to find out whether there is a constant $\delta \in [0, 1)$ and a convenient norm $\|\cdot\|$ on $\mathbb{R}^N$, such that $\|\vec{e}^{(m+1)}\| \leqslant \delta\|\vec{e}^{(m)}\|$ for any $\vec{e}^{(0)} \in \mathbb{R}^N$. □

Remark 2.5 (Convergence rate of the Richardson method)**.** The simplest iterative method for solving $A\vec{u} = \vec{f}$ might be $B = \omega I$, which is the well-known Richardson method in Example 1.2. In this case, the iteration converges if and only if $\rho(I - \omega A) < 1$, i.e., all eigenvalues of matrix A are in $(0, \frac{2}{\omega})$. Since A is SPD, the iteration converges if $\omega < 2\lambda_{\max}^{-1}(A)$. If we take $\omega = \lambda_{\max}^{-1}(A)$, then

$$\rho(I - \lambda_{\max}^{-1}(A)A) = 1 - \frac{\lambda_{\min}(A)}{\lambda_{\max}(A)} = 1 - \frac{1}{\kappa(A)}.$$

In fact, the optimal weight is $\omega_{\text{opt}} = \dfrac{2}{\lambda_{\max}(A) + \lambda_{\min}(A)}$ and

$$\rho(I - \omega_{\text{opt}}A) = \|I - \omega_{\text{opt}}A\| = 1 - \frac{2\lambda_{\min}(A)}{\lambda_{\max}(A) + \lambda_{\min}(A)} = \frac{\kappa(A) - 1}{\kappa(A) + 1}.$$

We can see that the convergence is very slow if A is ill-conditioned.

Remark 2.6 (Richardson and steepest descent method)**.** If A is a SPD matrix, then $A\vec{u} = \vec{f}$ is equivalent to the unconstrained quadratic minimization problem

$$\operatorname*{argmin}_{\vec{u}\in\mathbb{R}^N} \frac{1}{2}\vec{u}^T A\vec{u} - \vec{f}^T\vec{u}.$$

We immediately notice that the search direction in the Richardson method is exactly the same as the steepest decent method for the above minimization problem; see §2.2.

Symmetrization

In general, the iterator $\mathcal{B}$ might not be symmetric and it is more convenient to work with symmetric problems. We can apply a simple symmetrization step:

$$u^{(m+\frac{1}{2})} = u^{(m)} + \mathcal{B}\Big(f - \mathcal{A}u^{(m)}\Big), \tag{2.8}$$

$$u^{(m+1)} = u^{(m+\frac{1}{2})} + \mathcal{B}^T\Big(f - \mathcal{A}u^{(m+\frac{1}{2})}\Big), \tag{2.9}$$

to obtain a new iterative method

$$u - u^{(m+1)} = (\mathcal{I} - \mathcal{B}^T\mathcal{A})(\mathcal{I} - \mathcal{B}\mathcal{A})(u - u^{(m)}) = (\mathcal{I} - \mathcal{B}\mathcal{A})^*(\mathcal{I} - \mathcal{B}\mathcal{A})(u - u^{(m)}).$$

This new method satisfies the relation

$$u - u^{(m+1)} = (\mathcal{I} - \overline{\mathcal{B}}\mathcal{A})(u - u^{(m)})$$

and it has a symmetric iteration operator

$$\overline{\mathcal{B}} := \mathcal{B}^T + \mathcal{B} - \mathcal{B}^T\mathcal{A}\mathcal{B} = \mathcal{B}^T(\mathcal{B}^{-T} + \mathcal{B}^{-1} - \mathcal{A})\mathcal{B} =: \mathcal{B}^T\mathcal{K}\mathcal{B}. \tag{2.10}$$ eqn:Bsym

lem:decay **Lemma 2.2** (Error decay property)**.** *We have, for any $v \in V$, that*

$$\|v\|_{\mathcal{A}}^2 - \|(\mathcal{I} - \mathcal{B}\mathcal{A})v\|_{\mathcal{A}}^2 = \big(\overline{\mathcal{B}}\mathcal{A}v, v\big)_{\mathcal{A}},$$

or equivalently,

$$\big((\mathcal{I} - \overline{\mathcal{B}}\mathcal{A})v, v\big)_{\mathcal{A}} = \|(\mathcal{I} - \mathcal{B}\mathcal{A})v\|_{\mathcal{A}}^2.$$

Proof. Notice that, by the definition of symmetrization,

$$\overline{\mathcal{B}}\mathcal{A} = \mathcal{B}^T(\mathcal{B}^{-T} + \mathcal{B}^{-1} - \mathcal{A})\mathcal{B}\mathcal{A}.$$

This immediately gives

$$\begin{aligned}(\overline{\mathcal{B}}\mathcal{A}v, v)_{\mathcal{A}} &= \big((\mathcal{B}^{-T} + \mathcal{B}^{-1} - \mathcal{A})\mathcal{B}\mathcal{A}v, \mathcal{B}\mathcal{A}v\big) = (\mathcal{B}\mathcal{A}v, \mathcal{A}v) + (\mathcal{A}v, \mathcal{B}\mathcal{A}v) - (\mathcal{A}\mathcal{B}\mathcal{A}v, \mathcal{B}\mathcal{A}v)\\ &= \big((2\mathcal{I} - \mathcal{B}\mathcal{A})v, \mathcal{B}\mathcal{A}v\big)_{\mathcal{A}}\end{aligned}$$

and the first equality follows. The second equality is trivial. □

BAcontraction **Remark 2.7** (Contraction property)**.** The stationary iterative method defined by $\mathcal{B}$ is a contraction if $\|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}} \leqslant \delta_0 < 1$. Apparently, it is equivalent to say

$$\|e\|_{\mathcal{A}}^2 - \|(\mathcal{I} - \mathcal{B}\mathcal{A})e\|_{\mathcal{A}}^2 \geqslant (1 - \delta_0^2)\|e\|_{\mathcal{A}}^2 > 0, \qquad \forall e \neq 0.$$

Lemma 2.2 indicates that $\delta := \|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}} < 1$ if and only if $\overline{\mathcal{B}}$ is SPD. The constant δ is called the contraction factor of the iterative method. In the later discussions, we can assume the iterator $\mathcal{B}$ is SPD; in fact, if the iterator is not symmetric, we can consider its symmetrization.

Remark 2.8 (Effect of symmetrization)**.** We notice that $\overline{\mathcal{B}}^T = \overline{\mathcal{B}}$ and $(\mathcal{I} - \overline{\mathcal{B}}\mathcal{A})^* = \mathcal{I} - \overline{\mathcal{B}}\mathcal{A}$. Furthermore, Lemma 2.2 implies that $\big((\mathcal{I} - \overline{\mathcal{B}}\mathcal{A})v, v\big)_{\mathcal{A}} = \|(\mathcal{I} - \mathcal{B}\mathcal{A})v\|_{\mathcal{A}}^2$, $\forall v \in V$. Since $\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}$ is self-adjoint w.r.t. $(\cdot,\cdot)_{\mathcal{A}}$, we have $\|\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}\|_{\mathcal{A}} = \rho(\mathcal{I} - \overline{\mathcal{B}}\mathcal{A})$. And as a consequence,

$$\|\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}\|_{\mathcal{A}} = \sup_{\|v\|_{\mathcal{A}}=1} \big((\mathcal{I} - \overline{\mathcal{B}}\mathcal{A})v, v\big)_{\mathcal{A}} = \sup_{\|v\|_{\mathcal{A}}=1} \|(\mathcal{I} - \mathcal{B}\mathcal{A})v\|_{\mathcal{A}}^2 = \|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}}^2. \tag{2.11}$$ eqn:Bsym_B

This immediately gives

$$\rho(\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}) = \|\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}\|_{\mathcal{A}} = \|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}}^2 \geqslant \rho(\mathcal{I} - \mathcal{B}\mathcal{A})^2.$$

Hence, if the symmetrized method (2.8)–(2.9) converges, then the original method (2.7) also converges. Furthermore, we can obtain the following identity:

$$\|\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}\|_{\mathcal{A}} = \rho(\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}) = \sup_{v \in V\setminus\{0\}} \frac{\big((\mathcal{I} - \overline{\mathcal{B}}\mathcal{A})v, v\big)_{\mathcal{A}}}{\|v\|_{\mathcal{A}}^2}. \tag{2.12}$$ eqn:rate1

Convergence rate of stationary iterative methods

Since $\sigma(\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}) = \{1 - \lambda : \lambda \in \sigma(\overline{\mathcal{B}}\mathcal{A})\}$, the convergence of the symmetrized method (2.8)–(2.9) is equivalent to $\rho(\mathcal{I} - \overline{\mathcal{B}}\mathcal{A}) < 1$. Up to a constant scaling, it is equivalent to $\sigma(\overline{\mathcal{B}}\mathcal{A}) \subset (0, 1]$, i.e., $\overline{\mathcal{B}}\mathcal{A}$ is SPD w.r.t. $(\cdot,\cdot)_{\mathcal{A}}$. It is equivalent to say that $\overline{\mathcal{B}}$ is SPD w.r.t. $(\cdot,\cdot)$. We summarize the convergence conditions in the following theorem.

thm:symconv **Theorem 2.2** (Convergence of Symmetrized Algorithm)**.** *The symmetrized iteration* (2.8)–(2.9) *(with appropriate scaling) converges if and only if* $\overline{\mathcal{B}}$ *is SPD.*

ex:dJacobi **Example 2.2** (Jacobi and weighted Jacobi methods)**.** If $A \in \mathbb{R}^{N\times N}$ is SPD and it can be partitioned as $A = L + D + U$, where $L, D, U \in \mathbb{R}^{N\times N}$ are lower triangular, diagonal, upper triangular parts of A, respectively. We can immediately see that $B = D^{-1}$ yields the Jacobi method. In this case, we have

$$\overline{B} = B^T(B^{-T} + B^{-1} - A)B = D^{-T}(D - L - U)D^{-1}.$$

If $K_{\text{Jacobi}} := D - L - U = 2D - A$ is SPD, the Jacobi method converges. In general, it might not converge, but we can apply an appropriate scaling (i.e., the damped Jacobi method) $B_\omega = \omega D^{-1}$. We then derive

$$B_\omega^{-T} + B_\omega^{-1} - A = 2\omega^{-1}D - A.$$

The damping factor should satisfy that $\omega < \frac{2}{\rho(D^{-1}A)}$ in order to guarantee convergence. For the 1D finite difference problem of the Poisson's equation, we should use a *damping* factor $0 < \omega < 1$. □

Based on the identity (2.12), we can prove the convergence rate estimate:

thm:rate **Theorem 2.3** (Convergence rate)**.** *The convergence rate of the stationary iterative method (or its symmetrization) is*

$$\|\mathcal{I}-\mathcal{B}\mathcal{A}\|_{\mathcal{A}}^2=\|\mathcal{I}-\overline{\mathcal{B}}\mathcal{A}\|_{\mathcal{A}}=1-\frac{1}{c_1},\quad \textit{with}\ \ c_1:=\sup_{\|v\|_{\mathcal{A}}=1}(\overline{\mathcal{B}}^{-1}v,v).$$

Proof. The first equality is directly from (2.11). Since $((\mathcal{I}-\overline{\mathcal{B}}\mathcal{A})v,v)_{\mathcal{A}}=\|v\|_{\mathcal{A}}^2-(\overline{\mathcal{B}}\mathcal{A}v,v)_{\mathcal{A}}$, the identity (2.12) yields

$$\|\mathcal{I}-\mathcal{B}\mathcal{A}\|_{\mathcal{A}}^2=1-\inf_{\|v\|_{\mathcal{A}}=1}(\overline{\mathcal{B}}\mathcal{A}v,v)_{\mathcal{A}}=1-\lambda_{\min}(\overline{\mathcal{B}}\mathcal{A})=1-\frac{1}{c_1},$$

where

$$c_1=\lambda_{\max}((\overline{\mathcal{B}}\mathcal{A})^{-1})=\sup_{\|v\|_{\mathcal{A}}=1}\left((\overline{\mathcal{B}}\mathcal{A})^{-1}v,v\right)_{\mathcal{A}}=\sup_{\|v\|_{\mathcal{A}}=1}\left(\overline{\mathcal{B}}^{-1}v,v\right).$$

□

An example: modified G-S method ⋆

Similar to the weighted Jacobi method (see Example 2.2), we define the weighted G-S method $B_\omega=(\omega^{-1}D+L)^{-1}$. We have

$$B_\omega^{-T}+B_\omega^{-1}-A=(\omega^{-1}D+L)^T+(\omega^{-1}D+L)-(D+L+U)=(2\omega^{-1}-1)D.$$

The weighted G-S method converges if and only if $0<\omega<2$. In fact, $\omega=1$ yields the standard G-S method; $0<\omega<1$ yields the *SUR* method; $1<\omega<2$ yields the *SOR* method. One can select optimal weights for different problems to achieve good convergence result, which is beyond the scope of this lecture.

Motived by the weighted G-S methods, we assume there is an invertible smoother or a local relaxation method S for the equation $A\vec{u}=\vec{f}$, like the damped Jacobi smoother $S=\omega D^{-1}$ $(0<\omega<1)$. We can define a general or modified G-S method:

$$B:=\left(S^{-1}+L\right)^{-1}. \tag{2.13}$$ eqn:MGS

This method seems abstract and not very interesting now; but we will employ this idea on block matrices for multilevel iterative methods later on.

We can analyze the convergence rate of this modified G-S method using the same technique discussed above. Since $K=B^{-T}+B^{-1}-A$ is a symmetric operator and we can write (2.10) as $\overline{B}=B^TKB$. If B is the iteration operator defined by (2.13), we have

$$K=(S^{-T}+U)+(S^{-1}+L)-(D+L+U)=S^{-T}+S^{-1}-D.$$

Furthermore, from the definition of K, we find that $B^{-1} = K + A - B^{-T}$. Hence we get an explicit form of $\overline{B}^{-1}$ by simple calculations:

$$\overline{B}^{-1} = (K + A - B^{-T})K^{-1}(K + A - B^{-1}) = A + (A - B^{-T})K^{-1}(A - B^{-1}).$$

This identity and the definition of B yield:

$$\left(\overline{B}^{-1}\vec{v}, \vec{v}\right) = (A\vec{v}, \vec{v}) + \left(K^{-1}(D + U - S^{-1})\vec{v},\, (D + U - S^{-1})\vec{v}\right), \quad \forall \vec{v} \in \mathbb{R}^N.$$

Now we apply Theorem 2.3 and get the following identity for the convergence rate:

thm:rateGS **Corollary 2.1** (Convergence rate of Modified G-S)**.** *If $K = S^{-T} + S^{-1} - D$ is SPD, then the modified G-S method converges and*

$$\|I - BA\|_A^2 = \|I - \overline{B}A\|_A = 1 - \frac{1}{1 + c_0}, \quad \textit{with } \; c_0 := \sup_{\|\vec{v}\|_A = 1} \left\| K^{-\frac{1}{2}}\big(D + U - S^{-1}\big)\vec{v} \right\|^2.$$

This simple result will motivate our later analysis for subspace correction methods in Chapter 4.

ex:FDGS **Example 2.3** (Solving 1D FD equation using G-S)**.** If we apply the G-S method to the 1D FD system (1.27) for the Poisson's equation discussion in §1.2. For simplicity, we first multiply the both sides of the equation by h^2, namely, $A := \text{tridiag}(-1, 2, -1)$ and $\vec{f} := \big(h^2 f(x_i)\big)_{i=1}^N$. In this case, $S = D^{-1}$ and $K = D$. Corollary 2.1 shows that the convergence rate of the G-S iteration satisfies that

$$\|I - BA\|_A^2 = 1 - \frac{1}{1 + c_0}, \quad \text{with } \; c_0 = \sup_{\vec{v} \in \mathbb{R}^N \setminus \{0\}} \frac{(LD^{-1}U\vec{v}, \vec{v})}{\|\vec{v}\|_A^2}.$$

The positive constant can be further written

$$c_0 = \sup_{\vec{v} \in \mathbb{R}^N \setminus \{0\}} \frac{\big(D^{-1}U\vec{v}, U\vec{v}\big)}{(A\vec{v}, \vec{v})} = \sup_{\vec{v} \in \mathbb{R}^N \setminus \{0\}} \frac{\frac{1}{2}\big(U\vec{v}, U\vec{v}\big)}{(A\vec{v}, \vec{v})} = \sup_{\vec{v} \in \mathbb{R}^N \setminus \{0\}} \frac{\frac{1}{2}\sum_{i=2}^N v_i^2}{(A\vec{v}, \vec{v})}.$$

Because we have the eigenvalues of this discrete coefficient matrix A of FD (see Remark 1.10), we can estimate the denominator

$$(A\vec{v}, \vec{v}) \geqslant \lambda_{\min}(A)\|\vec{v}\|^2 = 4\sin^2\left(\frac{\pi}{2(N+1)}\right)\|\vec{v}\|^2.$$

Hence, asymptotically, we have the following estimate

$$c_0 \leqslant \sup_{\vec{v} \in \mathbb{R}^N \setminus \{0\}} \frac{\frac{1}{2}\|\vec{v}\|^2}{4\sin^2\left(\frac{\pi}{2(N+1)}\right)\|\vec{v}\|^2} \sim (N+1)^2 = h^{-2}.$$

Hence

$$\|I - BA\|_A \sim \sqrt{1 - \tilde{C}h^2} \sim 1 - Ch^2.$$

Similarly, for the FE equation, the condition number also likes $O(h^{-2})$ and convergence rate will deteriorate as the meshsize decreases. □

2.2 Krylov subspace methods

sec:KSM

Nonstationary iterative methods are more popular for standard-alone usage. Krylov subspace method (KSM) is a well-known class of nonstationary methods [35]. Let $\mathcal{A} : V \mapsto V$ be an invertible operator. By the Cayley–Hamilton theorem (see HW 2.3), there exists a polynomial of degree no more than $N-1$, $q_{N-1}(\lambda) \in \mathcal{P}_{N-1}$, such that $\mathcal{A}^{-1} = q_{N-1}(\mathcal{A})$. Hence the solution of the linear system has the form $u = q_{N-1}(\mathcal{A})f$. Krylov subspace methods construct iterative approximations to u in

$$\mathcal{K}_m := \operatorname{span}\{f, \mathcal{A}f, \mathcal{A}^2 f, \ldots, \mathcal{A}^{m-1} f\}, \quad m = 1, 2, \ldots$$

Gradient descent method

Let $\mathcal{A} : V \mapsto V$ be an SPD operator. Consider the following convex minimization problem:

$$\min_{u \in V} \mathcal{F}(u) := \frac{1}{2}(\mathcal{A}u, u) - (f, u). \tag{2.14}$$

eqn:min

Suppose we have an initial approximation u^{old} and construct a new approximation

$$u^{\text{new}} = u^{\text{old}} + \alpha\, p$$

with a fixed search direction $p \in V$ and a stepsize α. In order to find the "best possible" stepsize, we can solve an one-dimensional problem (i.e., the exact line-search method):

$$\min_{\alpha \in \mathbb{R}} \mathcal{F}(\alpha) := \frac{1}{2}\big(u^{\text{old}} + \alpha p, u^{\text{old}} + \alpha p\big)_{\mathcal{A}} - (f, u^{\text{old}} + \alpha p).$$

By simple calculation (HW 2.4), we obtain

$$\mathcal{F}(\alpha) := \frac{1}{2}\alpha^2(\mathcal{A}p, p) - \alpha(f - \mathcal{A}u^{\text{old}}, p) + \frac{1}{2}\big(\mathcal{A}u^{\text{old}}, u^{\text{old}}\big) - (f, u^{\text{old}}),$$

and the optimal stepsize is

$$\alpha = \frac{(f - \mathcal{A}u^{\text{old}}, p)}{(\mathcal{A}p, p)} = \frac{(r^{\text{old}}, p)}{(\mathcal{A}p, p)}, \quad \text{with } r^{\text{old}} = f - \mathcal{A}u^{\text{old}}. \tag{2.15}$$

eqn:optimalst

In the previous chapter, we have discussed the Richardson method. A nonstationary version of the Richardson method can be given as:

$$u^{(m+1)} = u^{(m)} + \alpha_m\big(f - \mathcal{A}u^{(m)}\big),$$

which can be viewed as the gradient descent or steepest descent (SD) method with exact line-search for the above convex minimization problem. This method is easy to implement and cheap in computation (each step only requires 1 matrix-vector multiplication and 2 inner products). Unfortunately, the SD method usually converges very slowly. See the following algorithm description of the SD method:

Listing 2.1: Steepest descent method

```
%% Given an initial guess u and a tolerance ε;
r ← f − Au;
while ‖r‖ > ε
    α ← (r, r)/(Ar, r);
    u ← u + α r;
    r ← r − αAr;
end
```

ex:relax **Example 2.4** (Line-search and the G-S method)**.** Let $V = \mathbb{R}^N$, $A = (a_{i,j}) \in \mathbb{R}^{N\times N}$. Suppose we choose the natural basis as the search directions, i.e., $\vec{p} = \vec{e}_i := (0, \dots, 0, 1, 0, \dots, 0)^T \in V$. Let $\vec{u}^{\,\text{old}} = \vec{u}^{\,(0)}$ be an initial guess. Then the above method yields the iteration:

$$\vec{u}^{\,(i)} = \vec{u}^{\,(i-1)} + \alpha\vec{p} = \vec{u}^{\,(i-1)} + \frac{(\vec{r}^{\,(i-1)}, \vec{p})}{(A\vec{p}, \vec{p})}\vec{p} = \vec{u}^{\,(i-1)} + \frac{(\vec{r}^{\,(i-1)}, \vec{e}_i)}{(A\vec{e}_i, \vec{e}_i)}\vec{e}_i.$$

So we get

$$\vec{u}^{\,(i)} = \vec{u}^{\,(i-1)} + \frac{f_i - \sum_{j=1}^N a_{i,j}\, u_j^{(i-1)}}{a_{i,i}}\vec{e}_i.$$

This means that only one entry is updated in each iteration:

$$u_i^{\text{new}} = u_i^{(i-1)} + \frac{f_i - \sum_{j=1}^N a_{i,j}\, u_j^{(i-1)}}{a_{i,i}} = \frac{1}{a_{i,i}}\Big(f_i - \sum_{j<i} a_{i,j}u_j^{\text{new}} - \sum_{j>i} a_{i,j}u_j^{\text{old}}\Big). \tag{2.16}$$ eqn:GSiter

After N steps ($i = 1, 2, \dots, N$), we obtain a new iteration $\vec{u}^{\,\text{new}}$, which is exactly the G-S iteration. Based on (2.16), we can write the G-S error propagation matrix in a different form

$$I - BA = (I - I_N a_{N,N}^{-1} I_N^T A)\cdots(I - I_1 a_{1,1}^{-1} I_1^T A) = (I - \Pi_N)\cdots(I - \Pi_1), \tag{2.17}$$ eqn:GSoper

where I_i is the natural embedding from $\text{span}\{\vec{e}_i\}$ to $\mathbb{R}^N$ and $\Pi_i = I_i A_i^{-1} I_i^T A$. This form of G-S will be further discussed later in the framework of Schwarz method and subspace correction method.

Theorem 2.4 (Convergence rate of steepest descent method)**.** *If we apply the exact line-search using the stepsize*

$$\alpha_m := \frac{\big(r^{(m)}, r^{(m)}\big)}{\big(r^{(m)}, r^{(m)}\big)_{\mathcal{A}}},$$

then the convergence rate of the SD method satisfies that

$$\big\|u - u^{(m)}\big\|_{\mathcal{A}} \leqslant \left(\frac{\kappa(\mathcal{A}) - 1}{\kappa(\mathcal{A}) + 1}\right)^m \big\|u - u^{(0)}\big\|_{\mathcal{A}}. \tag{2.18}$$ eqn:rateSD

Proof. The exact line-search stepsize is easy to obtain by 1D quadratic programming. At the m-th iteration, the energy satisfies that

$$\mathcal{F}(u^{(m+1)}) = \mathcal{F}(u^{(m)} + \alpha_m r^{(m)}) = \mathcal{F}(u^{(m)}) - \alpha_m (r^{(m)}, r^{(m)}) + \frac{1}{2}\alpha_m^2 (\mathcal{A} r^{(m)}, r^{(m)}).$$

By plugging the expression of α_m into the right-had side of the above equality, we obtain that

$$\mathcal{F}(u^{(m+1)}) = \mathcal{F}(u^{(m)}) - \frac{1}{2}\frac{(r^{(m)}, r^{(m)})^2}{(\mathcal{A} r^{(m)}, r^{(m)})}.$$

This implies that

$$\begin{aligned} \frac{\mathcal{F}(u^{(m+1)}) - \mathcal{F}(u)}{\mathcal{F}(u^{(m)}) - \mathcal{F}(u)} &= \frac{\mathcal{F}(u^{(m)}) - \dfrac{(r^{(m)}, r^{(m)})^2}{2(\mathcal{A} r^{(m)}, r^{(m)})} - \mathcal{F}(u)}{\mathcal{F}(u^{(m)}) - \mathcal{F}(u)} \\ &= 1 - \frac{(r^{(m)}, r^{(m)})^2}{(\mathcal{A} r^{(m)}, r^{(m)})(\mathcal{A}^{-1} r^{(m)}, r^{(m)})} =: 1 - \frac{1}{\beta} \end{aligned}$$

By the Kantorovich inequality, we know $\beta \leqslant \frac{(\lambda_{\max}+\lambda_{\min})^2}{4\lambda_{\max}\lambda_{\min}}$. So it follows

$$\frac{\mathcal{F}(u^{(m+1)}) - \mathcal{F}(u)}{\mathcal{F}(u^{(m)}) - \mathcal{F}(u)} = 1 - \frac{1}{\beta} \leqslant 1 - \frac{4\lambda_{\max}\lambda_{\min}}{(\lambda_{\max}+\lambda_{\min})^2} = \frac{(\lambda_{\max}-\lambda_{\min})^2}{(\lambda_{\max}+\lambda_{\min})^2} = \left(\frac{\kappa(\mathcal{A})-1}{\kappa(\mathcal{A})+1}\right)^2.$$

Hence the result. □

Conjugate gradient method

Now we consider a descent direction method with search direction $p^{(m)}$, i.e.

$$u^{(m+1)} = u^{(m)} + \alpha_m p^{(m)}. \tag{2.19}$$ eqn:descent

In this case, the "optimal" stepsize from the exact line-search is

$$\alpha_m := \frac{\left(r^{(m)}, p^{(m)}\right)}{\left(p^{(m)}, p^{(m)}\right)_{\mathcal{A}}}. \tag{2.20}$$ eqn:alpha

We notice that the residual after one iteration is

$$r^{(m+1)} = r^{(m)} - \alpha_m \mathcal{A} p^{(m)}.$$

In order to keep the iteration going, we wish to construct a new search direction which is orthogonal to the previous search directions. This motives us to define

$$p^{(m+1)} := r^{(m+1)} + \beta_m p^{(m)}, \quad \text{such that } \left(p^{(m)}, p^{(m+1)}\right)_{\mathcal{A}} = 0.$$

By simple calculations, we get the weight

$$\beta_m := -\frac{\left(\mathcal{A} r^{(m+1)}, p^{(m)}\right)}{\left(\mathcal{A} p^{(m)}, p^{(m)}\right)}. \tag{2.21}$$ eqn:beta

This is basically the so-called conjugate gradient (CG) method.

lem:cg1 **Lemma 2.3** (Properties of conjugate directions)**.** For any conjugate gradient step i, we have following identities:

1. $\left(r^{(i)}, p^{(i)}\right) = \left(r^{(i)}, r^{(i)}\right)$;
2. $\left(r^{(j)}, p^{(i)}\right) = 0, \qquad j > i$;
3. $\left(p^{(j)}, p^{(i)}\right)_{\mathcal{A}} = 0, \qquad j \neq i$;
4. $\left(r^{(j)}, r^{(i)}\right) = 0, \qquad j \neq i$.

This lemma is very simple but important; see HW 2.5. It guarantees we can apply a short recurrence iteration procedure while keep all directions are orthogonal to each other.

lem:cg2 **Lemma 2.4** (Stepsizes for CG)**.** For the conjugate gradient method, we have following identities:

1. $\alpha_m = \dfrac{\left(r^{(m)}, r^{(m)}\right)}{\left(\mathcal{A}p^{(m)}, p^{(m)}\right)}$;
2. $\beta_m = \dfrac{\left(r^{(m+1)}, r^{(m+1)}\right)}{\left(r^{(m)}, r^{(m)}\right)}$.

The previous lemma may look like some trivial transformations, but it is essential for CG implementation, which is described as follows:

Listing 2.2: Conjugate gradient method

```
%% Given an initial guess u and a tolerance ε;
r ← f − Au,  p ← r;
while ‖r‖ > ε
    α ← (r, r)/(Ap, p);
    ũ ← u + α p;
    r̃ ← r − α A p;
    β ← (r̃, r̃)/(r, r);
    p̃ ← r̃ + β p;
    Update:  u ← ũ,  r ← r̃,  p ← p̃;
end
```

The CG method converges much faster than the steepest descent. In fact, we have the following theorem

Theorem 2.5 (Convergence rate of CG)**.** *The convergence rate of the CG iteration satisfies the following estimate:*

$$\left\|u - u^{(m)}\right\|_{\mathcal{A}} \leqslant 2\left(\frac{\sqrt{\kappa(\mathcal{A})} - 1}{\sqrt{\kappa(\mathcal{A})} + 1}\right)^{m} \left\|u - u^{(0)}\right\|_{\mathcal{A}}. \tag{2.22}$$

eqn:rateCG

Proof. We only give a sketch of proof here. From Lemma 2.3, the residual $r^{(m)}$ is orthogonal to

$$\mathcal{K}_m = \text{span}\{r^{(0)}, \mathcal{A}r^{(0)}, \ldots, \mathcal{A}^{m-1}r^{(0)}\},$$

namely

$$(\mathcal{A}(u - u^{(m)}), v) = (r^{(m)}, v) = 0, \quad \forall v \in \mathcal{K}_m.$$

This implies

$$\big((u - u^{(0)}) - (u^{(m)} - u^{(0)}), v\big)_{\mathcal{A}} = 0, \quad \forall v \in \mathcal{K}_m.$$

The above $\mathcal{A}$-orthogonality gives

$$\begin{aligned}\left\|u - u^{(m)}\right\|_{\mathcal{A}} &= \min_{w\in\mathcal{K}_m} \left\|u - u^{(0)} - w\right\|_{\mathcal{A}} = \min_{q_{m-1}} \left\|u - u^{(0)} - q_{m-1}(\mathcal{A})r^{(0)}\right\|_{\mathcal{A}} \\ &= \min_{q_{m-1}} \left\|\big(I - q_{m-1}(\mathcal{A})\mathcal{A}\big)(u - u^{(0)})\right\|_{\mathcal{A}} = \min_{q_m(0)=1} \left\|q_m(\mathcal{A})(u - u^{(0)})\right\|_{\mathcal{A}}.\end{aligned}$$

The desired estimate can then be obtained by choosing appropriate Chebyshev polynomials; see HW 2.6 as a guideline to complete the proof. □

If the spectrum of $\mathcal{A}$ is uniformly distributed in the interval $[\lambda_{\min}, \lambda_{\max}]$, then the upper bound in (2.22) is sharp. In fact a few "bad eigenvalues" have almost no effect on the asymptotic convergence of the method. In this case, the bound is not sharp. The asymptotic convergence rate can be estimated by the *effective condition number* [2, 3].

Remark 2.9 (Computational complexity of CG)**.** We find that, in each iteration of the CG method, the complexity is only 1 matrix-vector multiplication and 2 inner products, with a few vector additions.

Remark 2.10 (MINRES method)**.** If $\mathcal{A} : V \mapsto V$ is a symmetric isomorphism mapping and it is indefinite, we can apply the minimum residual (MINRES) method characterized by

$$u^{(m)} = \operatorname*{argmin}_{v\in\mathcal{K}_m} \|f - \mathcal{A}v\|_0^2.$$

We can derive analytically that (see, for example, [35])

$$\left\|r^{(m)}\right\|_0 \leqslant \min_{q_m(0)=1} \max_{\lambda\in\sigma(\mathcal{A})} |q_m(\lambda)| \left\|r^{(0)}\right\|_0$$

In this case, the following crude convergence estimate holds

$$\left\|r^{(m)}\right\|_0 = \left\|\mathcal{A}(u - u^{(m)})\right\|_0 \leqslant 2\left(\frac{\kappa(\mathcal{A}) - 1}{\kappa(\mathcal{A}) + 1}\right)^m \left\|\mathcal{A}(u - u^{(0)})\right\|_0 = 2\left(\frac{\kappa(\mathcal{A}) - 1}{\kappa(\mathcal{A}) + 1}\right)^m \left\|r^{(0)}\right\|_0. \quad (2.23)$$ eqn:rateMINRES

If all the eigenvalues are positive, we can get sharp convergence estimate using Chebyshev polynomials. Unfortunately, it is not easy to get a general yet sharp estimate for indefinite problems.

Generalization to Hilbert spaces

It is important to note that the above convergence estimates (2.22) and (2.23) do not depend on the finite dimensionality N. Hence the Krylov subspace methods (KSMs) can be applied for operators $\mathcal{A} : \mathscr{V} \mapsto \mathscr{V}$. In view of Remark 1.7, we have

$$\|\mathcal{A}\|_{\mathscr{L}(\mathscr{V};\mathscr{V})} = \sup_{v\in\mathscr{V}} \frac{(\mathcal{A}v, v)}{\|v\|^2_{\mathscr{V}}} = \sup_{v\in\mathscr{V}} \frac{a[v,v]}{\|v\|^2_{\mathscr{V}}} \leqslant C_a$$

and the inf-sup condition (1.17) gives

$$\|\mathcal{A}^{-1}\|^{-1}_{\mathscr{L}(\mathscr{V};\mathscr{V})} = \inf_{v\in\mathscr{V}} \frac{\|\mathcal{A}v\|_{\mathscr{V}}}{\|v\|_{\mathscr{V}}} = \inf_{v\in\mathscr{V}} \sup_{u\in\mathscr{V}} \frac{(\mathcal{A}v, u)}{\|v\|_{\mathscr{V}}\|u\|_{\mathscr{V}}} = \inf_{v\in\mathscr{V}} \sup_{u\in\mathscr{V}} \frac{a[v,u]}{\|v\|_{\mathscr{V}}\|u\|_{\mathscr{V}}} \geqslant \alpha.$$

Hence the condition number $\kappa(\mathcal{A}) \leqslant C_a/\alpha$, which is bounded.

However, in order to employ KSMs for the continuous equations that we are interested in, like the Poisson's equation, we have to consider $\mathcal{A} : \mathscr{V} \mapsto \mathscr{W}$, where $\mathscr{V}$ and $\mathscr{W}$ are separable Hilbert spaces. Typically, $\mathscr{W} \supset \mathscr{V}$ and most likely $\mathscr{W} = \mathscr{V}'$. The following discussion directly follows the work by Mardal and Winther [45]. For simplicity, we consider a symmetric isomorphism $\mathcal{A} \in \mathscr{L}(\mathscr{V}; \mathscr{V}')$, i.e.,

$$\langle \mathcal{A}u, v\rangle = \langle \mathcal{A}v, u\rangle, \quad u, v \in \mathscr{V},$$

where $\langle\cdot,\cdot\rangle$ is the duality pair. Since $\mathscr{V}' \not\subset \mathscr{V}$, KSMs are not well-defined in this case.

We need to construct an isomorphism $\mathcal{B}$ mapping $\mathscr{V}'$ back to $\mathscr{V}$. We assume that the map $\mathcal{B}$ is symmetric and positive definite, namely $\langle\cdot, \mathcal{B}\cdot\rangle$ defines an inner product in $\mathscr{V}'$. We immediately notice that $\mathcal{B}$ could be a Riesz operator[2]: For any given $f \in \mathscr{V}'$,

$$(\mathcal{B}f, v)_{\mathscr{V}} = \langle f, v\rangle, \quad \forall v \in \mathscr{V}.$$

As a consequence, $\langle \mathcal{B}^{-1}\cdot,\cdot\rangle$ is an inner product on $\mathscr{V}$, with associated norm equivalent to $\|\cdot\|_{\mathscr{V}}$. This leads to a so-called preconditioned system

$$\mathcal{B}\mathcal{A}u = \mathcal{B}f$$

and $\mathcal{B}\mathcal{A}$ is an isomorphism from $\mathscr{V}$ to itself. The Krylov subspace methods can be applied to this preconditioned system and $\mathcal{B}$ is called a preconditioner.

Note that $\mathcal{B}\mathcal{A} : \mathscr{V} \mapsto \mathscr{V}$ is symmetric with respect to $(\cdot,\cdot)_{\mathscr{V}}$, i.e.,

$$(\mathcal{B}\mathcal{A}u, v)_{\mathscr{V}} = \langle \mathcal{A}u, v\rangle = a[u,v] = (u, \mathcal{B}\mathcal{A}v)_{\mathscr{V}}, \quad u, v \in \mathscr{V}.$$

The last equality follows from the symmetry of the bilinear form $a[\cdot,\cdot]$. Furthermore, due to the continuity of $a[\cdot,\cdot]$ (1.15), we obtain

$$\|\mathcal{B}\mathcal{A}\|_{\mathscr{L}(\mathscr{V};\mathscr{V})} = \sup_{v\in\mathscr{V}} \frac{|(\mathcal{B}\mathcal{A}v, v)_{\mathscr{V}}|}{\|v\|^2_{\mathscr{V}}} = \sup_{v\in\mathscr{V}} \frac{a[v,v]}{\|v\|^2_{\mathscr{V}}} \leqslant C_a$$

[2]We note that $\mathcal{B}$ is inner product dependent.

and the inf-sup condition (1.17) gives

$$\|(\mathcal{B}\mathcal{A})^{-1}\|^{-1}_{\mathscr{L}(\mathscr{V};\mathscr{V})} = \inf_{v\in\mathscr{V}} \frac{\|\mathcal{B}\mathcal{A}v\|_{\mathscr{V}}}{\|v\|_{\mathscr{V}}} = \inf_{v\in\mathscr{V}} \sup_{u\in\mathscr{V}} \frac{(\mathcal{B}\mathcal{A}v,u)_{\mathscr{V}}}{\|v\|_{\mathscr{V}}\|u\|_{\mathscr{V}}} = \inf_{v\in\mathscr{V}} \sup_{u\in\mathscr{V}} \frac{a[v,u]}{\|v\|_{\mathscr{V}}\|u\|_{\mathscr{V}}} \geqslant \alpha.$$

Example 2.5 (Poisson solver as a preconditioner)**.** As an example, we consider a second-order elliptic operator $\mathcal{A} : H^1_0(\Omega) \mapsto H^{-1}(\Omega)$. We need to define

$$(\mathcal{B}f, v)_{H^1_0(\Omega)} := (\nabla\mathcal{B}f, \nabla v)_{0,\Omega} = \langle f, v\rangle.$$

In this sense, we can choose $\mathcal{B} = (-\Delta)^{-1}$ as a preconditioner. We note that other inner products can be used, which will yield different preconditioners. As long as the above continuity condition and the inf-sup condition hold, the preconditioned system is well-conditioned. □

Now we summarize the above discussion on how to construct a "natural" preconditioner:

1. Define an appropriate inner product $(\cdot,\cdot)_{\mathscr{V}}$;
2. Establish the inf-sup condition $\sup_{v\in\mathscr{V}} \frac{a[u,v]}{\|v\|_{\mathscr{V}}} \geqslant \alpha\|u\|_{\mathscr{V}}$ for any $u \in \mathscr{V}$;
3. Define $\mathcal{B}$ as the Reisz operator, i.e., $(\mathcal{B}f, v)_{\mathscr{V}} = \langle f, v\rangle$ for any $v \in \mathscr{V}$;
4. The preconditioned system $\mathcal{B}\mathcal{A}$ is symmetric with respect to $(\cdot,\cdot)_{\mathscr{V}}$ and well-conditioned;
5. Construct a discretization which satisfies the corresponding discrete inf-sup condition;
6. Define a spectrally equivalent $\mathcal{B}_h$ as a preconditioner.

2.3 Condition number and preconditioning

The convergence rate of an iterative method depends greatly on the spectrum of the coefficient matrix. Hence, iterative methods usually involve a second matrix that transforms the coefficient matrix into one with a more favorable spectrum. The transformation matrix is called a preconditioner. A good preconditioner $\mathcal{B}$ improves the convergence of the iterative method sufficiently and is relatively cheap to compute, in order to overcome the overhead (extra cost) of constructing and applying the preconditioner. There are a few ways to apply preconditioners, for example:

$\mathcal{B}\mathcal{A}u = \mathcal{B}f$		Left preconditioning
$\mathcal{A}\mathcal{B}v = f$	$u = \mathcal{B}v$	Right preconditioning
$\mathcal{B}_L\mathcal{A}\mathcal{B}_R v = \mathcal{B}_L f$	$u = \mathcal{B}_R v$	Split preconditioning

Construction of preconditioners

We first introduce a few simple facts that could be helpful when we need to estimate the condition number $\kappa(\mathcal{BA})$. Although convergence behavior of iterative methods is not governed by the condition number alone, it provides useful information for a variety of methods. For example, we would hope that $\kappa(\mathcal{BA}) \ll \kappa(\mathcal{A})$, if we apply the steepest descent method or the CG method to solve a linear system.

It is desirable to have an optimal preconditioner which satisfy most of, if not all, the following properties:

- The preconditioned linear systems have improved convergence behavior. Furthermore, the spectral condition number of $\mathcal{BA}$ should be bounded independently of the size of the problem.
- The preconditioner is relatively easy to setup and cheap to apply. The computational cost of $\mathcal{B}r$ should be proportional to the size of the problem.
- The preconditioner should be robust on different domain shapes, mesh types, jumps in coefficients, etc.
- The preconditioner can be parallelized easily and efficiently.

lem:condnum **Lemma 2.5** (Estimation of condition number)**.** If μ_0 and μ_1 are positive constants satisfying

$$\mu_0\big(\mathcal{A}u, u\big) \leqslant \big(\mathcal{B}^{-1}u, u\big) \leqslant \mu_1\big(\mathcal{A}u, u\big), \quad \forall u \in V, \tag{2.24}$$ eqn:condnum1

then the condition number

$$\kappa(\mathcal{BA}) \leqslant \frac{\mu_1}{\mu_0}.$$

Proof. By change of variable $u = \mathcal{A}^{-\frac{1}{2}}v$, we have $\sigma\big(\mathcal{A}^{-\frac{1}{2}}\mathcal{B}^{-1}\mathcal{A}^{-\frac{1}{2}}\big) \subseteq \big[\mu_0, \mu_1\big]$ and, hence, $\sigma\big((\mathcal{BA})^{-1}\big) \subseteq \big[\mu_0, \mu_1\big]$. □

Sometimes, it is more convenient to use the following equivalent conditions of (2.24):

equiv-condnum **Lemma 2.6** (Some equivalent conditions)**.** If $\mathcal{A}$ and $\mathcal{B}$ are symmetric positive definite operators on a finite-dimensional space V, then we have the inequalities (2.24) are equivalent to

$$\mu_0(\mathcal{B}u, u) \leqslant (\mathcal{A}^{-1}u, u) \leqslant \mu_1(\mathcal{B}u, u), \quad \forall u \in V, \tag{2.25}$$ eqn:condnum2

or

$$\mu_1^{-1}(\mathcal{A}u, u) \leqslant (\mathcal{ABA}u, u) \leqslant \mu_0^{-1}(\mathcal{A}u, u), \quad \forall u \in V, \tag{2.26}$$ eqn:condnum3

or

$$\mu_1^{-1}(\mathcal{B}u, u) \leqslant (\mathcal{BAB}u, u) \leqslant \mu_0^{-1}(\mathcal{B}u, u), \quad \forall u \in V. \tag{2.27}$$ eqn:condnum4

Proof of this lemma is left to the readers as an exercise; see HW 2.7.

rem:AmABA **Remark 2.11** (Another equivalent condition)**.** If $\mathcal{A}$ and $\mathcal{B}$ are symmetric positive definite operators on a finite-dimensional space V, $\alpha > 0$ and $0 < \delta < 1$, then it is easy to verify the following two conditions are equivalent:

$$-\alpha(\mathcal{A}u, u) \leqslant (\mathcal{A}(\mathcal{I} - \mathcal{B}\mathcal{A})u, u) \leqslant \delta(\mathcal{A}u, u), \quad \forall u \in V \tag{2.28}$$ eqn:condnum5

and

$$(1+\alpha)^{-1}(\mathcal{A}u, u) \leqslant (\mathcal{B}^{-1}u, u) \leqslant (1-\delta)^{-1}(\mathcal{A}u, u), \quad \forall u \in V. \tag{2.29}$$ eqn:condnum6

Let $\mathcal{B}$ be a symmetric iterator of the SPD operator $\mathcal{A}$. We have seen that a sufficient condition for the iterative method to be convergent is that

$$\rho(\mathcal{I} - \mathcal{B}\mathcal{A}) < 1.$$

In this case, $\rho := \|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}} < 1$. The method is not only converging but also a contraction, i.e., $\|u - u^{(m)}\|_{\mathcal{A}} \leqslant \rho^m \|u - u^{(0)}\|_{\mathcal{A}} \to 0$ as $m \to +\infty$. Similar argument as Theorem 2.2 shows that $\mathcal{B}$ must be SPD. Furthermore, by definition, we have

$$\Big(\big(\mathcal{A} - 2\mathcal{A}\mathcal{B}\mathcal{A} + \mathcal{A}\mathcal{B}\mathcal{A}\mathcal{B}\mathcal{A}\big)u, u \Big) \leqslant \rho^2 (u, u)_{\mathcal{A}}.$$

Changing variable $v = \mathcal{A}^{1/2}u$, we obtain

$$\begin{aligned} \Big(\big(\mathcal{I} - \mathcal{A}^{1/2}\mathcal{B}\mathcal{A}^{1/2}\big)^2 v, v \Big) \leqslant \rho^2 (v, v) \quad &\Longrightarrow \quad \Big| \big((\mathcal{I} - \mathcal{A}^{1/2}\mathcal{B}\mathcal{A}^{1/2})v, v\big) \Big| \leqslant \rho (v, v) \\ &\Longrightarrow \quad \Big| \big((\mathcal{A} - \mathcal{A}\mathcal{B}\mathcal{A})u, u\big) \Big| \leqslant \rho (\mathcal{A}u, u), \quad \forall u \in V. \end{aligned}$$

Hence Remark 2.11 shows (see HW 2.9) that the condition number is uniformly bounded, i.e.,

$$\kappa(\mathcal{B}\mathcal{A}) \leqslant \frac{1+\rho}{1-\rho}.$$

In fact, the above estimate can also be easily obtained from $\rho(\mathcal{I} - \mathcal{B}\mathcal{A}) = \rho < 1$.

Preconditioned conjugate gradient method

The first question to answer is whether and how CG can still be applied to the preconditioned system $\mathcal{B}\mathcal{A}u = \mathcal{B}f$. As we mentioned $\mathcal{B}\mathcal{A}$ is not symmetric w.r.t. $(\cdot,\cdot)$ but symmetric w.r.t. $(\cdot,\cdot)_{\mathcal{A}}$. Similarly, we can define a new inner product $(\cdot,\cdot)_{\mathcal{B}^{-1}} := (\mathcal{B}^{-1}\cdot, \cdot)$. Then

$$(\mathcal{B}\mathcal{A}\cdot, \cdot)_{\mathcal{B}^{-1}} = (\mathcal{A}\cdot, \cdot) \quad \Longrightarrow \quad \mathcal{B}\mathcal{A} \text{ is SPD w.r.t. } (\cdot,\cdot)_{\mathcal{B}^{-1}},$$

which means CG can be applied to $\mathcal{B}\mathcal{A}u = \mathcal{B}f$ with the new inner product.

lem:pcg2 **Lemma 2.7** (Stepsizes of PCG)**.** For the preconditioned conjugate gradient method, we have the following identities:

1. $\alpha_m = \dfrac{\left(\mathcal{B}r^{(m)}, r^{(m)}\right)}{\left(\mathcal{A}p^{(m)}, p^{(m)}\right)}$,
2. $\beta_m = \dfrac{\left(\mathcal{B}r^{(m+1)}, r^{(m+1)}\right)}{\left(\mathcal{B}r^{(m)}, r^{(m)}\right)}$.

We notice that $\mathcal{B}^{-1}$ is cancelled out in the above inner products. With the help of this lemma, we can write the pseudo-code of PCG with left preconditioner (compared with regular CG, it just requires one more matrix-vector multiplication):

Listing 2.3: Preconditioned conjugate gradient method

```
%% Given an initial guess u and a tolerance ε;
r ← f − Au,  p ← Br;
while ‖r‖ > ε
    α ← (Br, r)/(Ap, p);
    ũ ← u + α p;
    r̃ ← r − αAp;
    β ← (Br̃, r̃)/(Br, r);
    p̃ ← Br̃ + β p;
    Update:  u ← ũ,  r ← r̃,  p ← p̃;
end
```

Remark 2.12 (Iterator and preconditioner)**.** We use the same notation $\mathcal{B}$ for the preconditioner and the iterator, apparently for a reason. Indeed, the convergence rate of the preconditioned CG method (2.22) is equal to

$$\delta_{\mathrm{CG}} = \frac{\sqrt{\kappa(\mathcal{BA})}-1}{\sqrt{\kappa(\mathcal{BA})}+1} \leqslant \frac{\sqrt{\frac{1+\rho}{1-\rho}}-1}{\sqrt{\frac{1+\rho}{1-\rho}}+1} = \frac{1-\sqrt{1-\rho^2}}{\rho} < \rho.$$

The last inequality holds true when $0 < \rho < 1$. Hence, for any convergent stationary linear iterative method, a preconditioner can be found and its convergence can be accelerated by PCG. Of course, it comes with the extra cost of applying the preconditioners. Preconditioning is so important for practical problems and KSMs are sometimes referred as accelerators.

Stopping criteria ⋆

When an iterative method is employed, sometimes it is hard to determine when to stop the iteration process. Ultimately we would like to have the error $e^{(m)} = u - u^{(m)}$ in certain norm (e.g. the energy norm) to be small enough, i.e., $\left(e^{(m)}, e^{(m)}\right)_{\mathcal{A}}^{\frac{1}{2}} < \epsilon$. However, the error is not usually computable. Norms of the residual $r^{(m)} = f - \mathcal{A}u^{(m)}$, which is not only computable but

also naturally available in the iterative process, are used instead. According to the standard perturbation analysis, we have

$$\frac{\|u-u^{(m)}\|}{\|u\|} \leqslant \kappa(\mathcal{A})\frac{\|r^{(m)}\|}{\|f\|}. \tag{2.30}$$ eqn:ResErr

In fact, $\mathcal{A}(u-u^{(m)}) = f - \mathcal{A}u^{(m)} = r^{(m)}$. Hence $\|u-u^{(m)}\| \leqslant \|\mathcal{A}^{-1}\|\|r^{(m)}\|$. On the other hand, it is easy to see that $\|f\| \leqslant \|\mathcal{A}\|\|u\|$. By combining the last two inequalities, we can obtain the desired estimate (2.30). We notice that the right-hand side of (2.30) is the relative residual (with initial guess equals zero) and the left-hand side is just the relative error. Hence this inequality shows that, even if the relative residual is small, the relative error could be still very large, especially for the ill-conditioned problems.

Although L^2-norm of $r^{(m)}$ is usually used in practice, $\left(r^{(m)}, r^{(m)}\right)_{\mathcal{B}}^{\frac{1}{2}}$ is a better quantity to monitor for convergence. We notice that

$$(r^{(m)}, r^{(m)})_{\mathcal{B}} = \left(\mathcal{A}e^{(m)}, \mathcal{A}e^{(m)}\right)_{\mathcal{B}} = \left(\mathcal{A}\mathcal{B}\mathcal{A}e^{(m)}, e^{(m)}\right).$$

According to Lemma 2.6, $\left(r^{(m)}, r^{(m)}\right)_{\mathcal{B}}^{\frac{1}{2}}$ is equivalent to $\left(e^{(m)}, e^{(m)}\right)_{\mathcal{A}}^{\frac{1}{2}}$, if $\mathcal{B}$ is a *good* preconditioner. In practice, there might be situations that left part of the preconditioner changes the residual of the equation a lot, which will cause trouble for users to design stopping criteria. The preconditioned equation has a residual $r_{\mathcal{B}} = \mathcal{B}r = \mathcal{B}(f - \mathcal{A}u)$ and $\|r_{\mathcal{B}}\|$ might be a lot different than $\|r\|$. Thus it is usually not good to use $r_{\mathcal{B}}$ instead of r.

2.4 Domain decomposition methods

sec:DDM

In the field of numerical methods for partial differential equations, domain decomposition methods (DDMs) make use of *divide and conquer* techniques by iteratively solving subproblems defined on smaller subdomains. It is a convenient framework for the solution and, more importantly, preconditioning of heterogeneous or multiphysics problems. It can be used in the framework of many discretization methods (e.g., FD and FE) to make their algebraic solution efficient, especially on parallel computers. Roughly speaking, there are two ways of subdividing the computational domain, overlapping and non-overlapping. We will only discuss overlapping domain decomposition methods here.

Divide and conquer

We consider the model boundary value problem

$$\begin{cases} \mathcal{A}u &= f & \text{in } \Omega \\ u &= 0 & \text{on } \partial\Omega \end{cases}$$

Overlapping domain decomposition algorithms are based on a decomposition of the domain Ω into a number of overlapping subdomains. To introduce the main ideas of DDMs, we consider the case of two overlapping subdomains Ω_1 and Ω_2, which form a covering of Ω and $\Omega_1 \bigcap \Omega_2 \neq \emptyset$; see Figure 1. We let Γ_i $(i = 1, 2)$ denote the part of the boundary of Ω_i, which is in the interior of Ω.

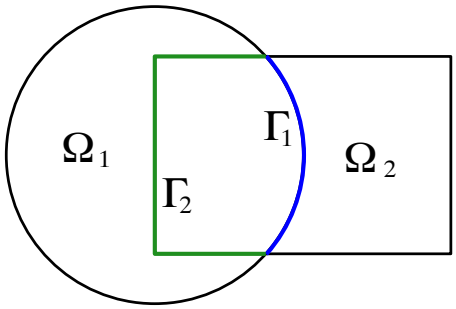


fig:ddm

Figure 2.1: Overlapping domain partition with two sub-domains.

If we already have an approximate solution $u^{(m)}$, we can construct a new approximation by solving the following two equations:

$$\begin{cases} \mathcal{A}u_1^{(m+1)} &= f & \text{in } \Omega_1, \\ u_1^{(m+1)} &= u^{(m)} & \text{on } \Gamma_1, \\ u_1^{(m+1)} &= 0 & \text{on } \partial\Omega_1 \backslash \Gamma_1, \end{cases}$$

and

$$\begin{cases} \mathcal{A}u_2^{(m+1)} &= f & \text{in } \Omega_2, \\ u_2^{(m+1)} &= g^{(m)} & \text{on } \Gamma_2, \\ u_2^{(m+1)} &= 0 & \text{on } \partial\Omega_2 \backslash \Gamma_2. \end{cases}$$

There are two approaches to apply these two subdomain corrections—the additive approach and the multiplicative approach. In the additive approach, we take $g^{(m)} = u^{(m)}$ and carry out the two corrections simultaneously. In the multiplicative approach, we take $g^{(m)} = u_1^{(m+1)}$ and use the most up-to-date iterative information. We then define the new iteration as

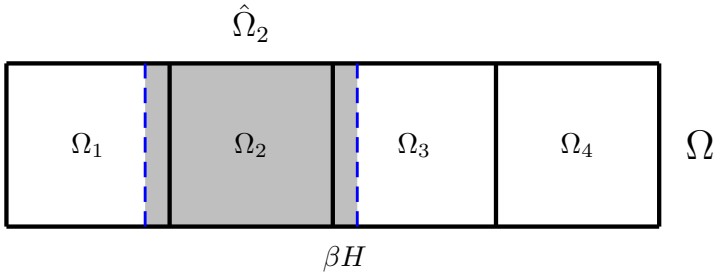
$$u^{(m+1)}(x) := \begin{cases} u_2^{(m+1)}, & \text{if } x \in \Omega_2; \\ u_1^{(m+1)}, & \text{if } x \in \Omega \backslash \Omega_2. \end{cases}$$

Overlapping DD methods

With the above motivation in mind, we are ready to introduce the standard overlapping domain decomposition method in matrix form:

$$A\vec{u} = \vec{f}, \qquad V = \mathbb{R}^N.$$

Suppose we have an one-dimensional domain partitioning of Ω; see Figure 2.2. Of course, we can use more general partitioning strategies as well.



g:overlapping

Figure 2.2: Overlapping domain partition with four sub-domains.

Denote the set of grid point indices as $G := \{1, 2, \ldots, N\}$ and it is partitioned into n subdomains. Let $\hat{G}_i$ be the index set of the interior points of $\hat{\Omega}_i$, and $N_i := |\hat{G}_i|$ be the cardinality of $\hat{G}_i$. Apparently, we have

$$G = \hat{G}_1 \bigcup \hat{G}_2 \bigcup \cdots \bigcup \hat{G}_n \qquad \text{and} \qquad N < N_1 + N_2 + \cdots + N_n.$$

In the matrix form, the injection matrix (natural embedding) $I_i \in \mathbb{R}^{N \times N_i}$ is defined as

$$(I_i \vec{v}_i)_k = \begin{cases} (\vec{v}_i)_k, & \text{if } k \in \hat{G}_i; \\ 0, & \text{if } k \in G \backslash \hat{G}_i. \end{cases} \tag{2.31}$$

eqn:restricti

It is natural to define sub-problems as $A_i := I_i^T A I_i$ $(i = 1, \ldots, n)$. If we solve each sub-problem exactly, then we have $B_i := I_i A_i^{-1} I_i^T$.

We can define an *additive Schwarz method* (ASM) as

$$B_{\text{as}} := \sum_{i=1}^{n} B_i = \sum_{i=1}^{n} I_i A_i^{-1} I_i^T, \tag{2.32}$$

eqn:ASM

which generalizes the block Jacobi method. Similarly, a *multiplicative Schwarz method* (MSM) is then defined by the following error propagation operator

$$I - B_{\text{ms}} A := (I - B_n A) \cdots (I - B_1 A) = \prod_{i=n}^{1} (I - B_i A). \tag{2.33}$$

eqn:MSM

This is a generalization of the block G-S method (with overlapping blocks). In practice, the sub-problem solver A_i^{-1} could be replaced by an approximation, like the ILU method.

Classical convergence results of overlapping DDMs ⋆

These DD methods, especially the ASM version, are usually applied as preconditioners for parallel computing. Its convergence has been analyzed in [27, 28] and we only show the results for the additive version here.

Theorem 2.6 (ASM DD preconditioner)**.** *The condition number of ASM domain decomposition method is independent of the mesh size h and satisfies*

$$\kappa(B_{as}A) \lesssim H^{-2}(1+\beta^{-2}),$$

where H is size of domain partitions and βH characterizes size of the overlaps.

The DD preconditioner (2.32) performs very well in practice. But the convergence rate still depends on H and the condition number could be large if H is very small. A simple approach to get rid of this dependence on H is to introduce a coarse space $V_0 \subset V$ and a corresponding coarse-level solver, i.e.

$$B_{\text{as},2} := I_0 A_0^{-1} I_0^T + \sum_{i=1}^{n} I_i A_i^{-1} I_i^T,$$

where $I_0 : V_0 \mapsto V$ is the injection matrix and A_0 is the coarse space problem. We then have the following estimate on the condition number:

Theorem 2.7 (Two-level ASM DD preconditioner)**.** *The condition number of ASM domain decomposition method is independent of the mesh size h and satisfies*

$$\kappa(B_{as,2}A) \lesssim 1+\beta^{-1}.$$

The above theorem shows the dependence on meshsize can be removed by introducing an appropriate coarse-level correction. We will construct and analyze two-level and, more generally, multilevel iterative methods in the following chapters.

2.5 Homework problems

hw:BtBstar **HW 2.1.** Show the identity (2.6).

w:A-transpose **HW 2.2.** If $\mathcal{B}^T = \mathcal{B}$, show that $(\mathcal{B}\mathcal{A})^* = \mathcal{B}^T\mathcal{A} = \mathcal{B}\mathcal{A}$.

hw:CHThm **HW 2.3.** Let $A \in \mathbb{R}^{N\times N}$ and $q(\lambda) := |\lambda I - A|$ be the characteristic polynomial of A. Show the Cayley-Hamilton theorem, i.e., $q(A) = 0$.

hw:SDalpha **HW 2.4.** Show the optimal stepsize (2.15) for general descent direction method.

ateDirections **HW 2.5.** Prove Lemmas 2.3 and 2.4.

hw:Cheb **HW 2.6.** The Chebyshev (or Tchebycheff) polynomial of first kind on $[-1,1]$ can be defined recursively as

$$T_0(x) = 1, \quad T_1(x) = x, \quad T_{n+1}(x) = 2xT_n(x) - T_{n-1}(x).$$

Show that

$$T_n(x) = \frac{1}{2}\Big(\big(x + \sqrt{x^2-1}\big)^n + \big(x - \sqrt{x^2-1}\big)^n \Big)$$

and $\big|T_n(x)\big| \leqslant 1$ for any $x \in [-1,1]$. Let $0 < \lambda_{\min} \leqslant \lambda_{\max}$. Define

$$S_n(\lambda) := \left[T_n\Big(\frac{\lambda_{\max}+\lambda_{\min}}{\lambda_{\max}-\lambda_{\min}} \Big) \right]^{-1} T_n\Big(\frac{\lambda_{\max}+\lambda_{\min}-2\lambda}{\lambda_{\max}-\lambda_{\min}} \Big)$$

and we have

$$\left| T_n\Big(\frac{\lambda_{\max}+\lambda_{\min}}{\lambda_{\max}-\lambda_{\min}} \Big) \right|^{-1} = \big\| S_n \big\|_{\infty,[\lambda_{\min},\lambda_{\max}]} = \min_{p\in\mathcal{P}_n;\, p(0)=1} \big\| p \big\|_{\infty,[\lambda_{\min},\lambda_{\max}]},$$

where $\mathcal{P}_n$ is the set of polynomials of degree less than or equal to n.

equiv-condnum **HW 2.7.** Prove Lemma 2.6.

equiv-condnum2 **HW 2.8.** Show that (2.28) and (2.29) are equivalent to each.

hw:condBA **HW 2.9.** Let $\mathcal{A}$ be SPD and $\mathcal{B}$ be a symmetric iterator. If $\rho = \|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}} < 1$, then $\mathcal{B}$ is also SPD and

$$\kappa(\mathcal{B}\mathcal{A}) \leqslant \frac{1+\rho}{1-\rho}.$$

Chapter 3

Two-grid Methods

ch:twogrid

In the previous chapter, we have seen several simple iterative solvers and preconditioners for solving the linear algebraic system (2.1). However, the convergence rate of most methods discussed in the previous chapter deteriorates when meshsize h goes to zero, except the overlapping domain decomposition method with coarse-grid correction. In this chapter, we discuss the two-grid (or more generally, two-level) method for the discrete Poisson's equation:

$$\begin{cases} -\Delta u = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega. \end{cases} \quad \Longrightarrow \quad A\vec{u} = \vec{f}.$$

In Chapter 1, we have briefly discussed the finite element approximation for this model problem. From now on, we will mainly discuss in the context of finite element discretizations.

Throughout this chapter, we use the standard notations for Sobolev spaces introduced in Chapter 1: $H^k(\Omega)$ denotes the classical Sobolev space of scalar functions on a bounded domain $\Omega \subset \mathbb{R}^d$ whose derivatives up to order k are square integrable, with the full norm $\|\cdot\|_k$ and the corresponding semi-norm $|\cdot|_k$. The symbol $H_0^1(\Omega)$ denotes the subspace of $H^1(\Omega)$ whose trace vanishes on the boundary $\partial\Omega$. We will also discuss the corresponding spaces restricted to the subdomain of Ω.

3.1 Finite element methods

sec:FEM

We now take a little detour and say a few more words about finite element discretizations; see [20] for details. The linear operator $\mathcal{A} : \mathscr{V} \mapsto \mathscr{V}'$ is defined by

$$(\mathcal{A}u, v) := a[u, v] = \int_\Omega \nabla u \nabla v \, dx, \quad \forall v \in \mathscr{V}$$

and $f \in \mathscr{V}'$ is a function or distribution. Suppose that $\mathcal{A}$ is bounded (1.15), i.e.,

$$a[u, v] \leqslant C_a \|u\|_{\mathscr{V}} \|v\|_{\mathscr{V}}, \quad \forall\, u, v \in \mathscr{V}$$

and coercive (1.24), i.e.,

$$a[v,v] \geqslant \alpha \|v\|_{\mathscr{V}}^2, \quad \forall v \in \mathscr{V}.$$

We would like to find $u \in \mathscr{V}$ such that $\mathcal{A}u = f$, which is well-posed; see Remark 1.7.

Galerkin approximation

The *Galerkin method* exploits the weak formulation and replaces the underlying function space by appropriate finite dimensional subspaces. We choose a finite dimensional space V_N (trial/test space), which is an approximation to the space $\mathscr{V}$ with $\dim(V_N) = N$. When it causes no confusion, we shall just drop the subscript and denote $V = V_N$. Then we arrive at the Galerkin discretization:

$$\text{Find } u_N \in V: \quad a[u_N, v_N] = \langle f, v_N \rangle, \quad \forall\, v_N \in V. \tag{3.1}$$

eq:ell_d

Equation (3.1) yields the so-called *Galerkin discretization*. If the bilinear form $a[\cdot,\cdot]$ is symmetric and coercive, it is called the Ritz–Galerkin discretization. In the finite-dimensional setting, we can identify the dual space V' and V; this way, the duality pair $\langle \cdot, \cdot \rangle$ becomes the l^2-inner product $(\cdot,\cdot)$.

For conforming discretizations, the bilinear form $a[\cdot,\cdot]$ is well-defined on $V \times V$. If the bilinear form $a[\cdot,\cdot]$ is coercive, then we have

$$a[v_N, v_N] \geqslant \alpha_{_N} \|v_N\|_{\mathscr{V}}^2, \quad \forall\, v_N \in V.$$

Since coercivity is inherited from $\mathscr{V}$ to V, we can see that the constants α_N is bounded from below, i.e.,

$$\alpha_{_N} \geqslant \alpha, \quad \forall N$$

As a consequence, the inf-sup condition holds[1]. It is easy to show the following simple optimality approximation properties.

rem:GO **Remark 3.1** (Galerkin Orthogonality)**.** Assume $V \subset \mathscr{V}$. The weak formulations of the exact and discrete solutions satisfy

$$\begin{cases} a[u,v] = \langle f, v \rangle, & \forall\ v \in \mathscr{V}; \\ a[u_N, v_N] = \langle f, v_N \rangle, & \forall\ v_N \in V. \end{cases}$$

Taking $v = v_N$ in the first equation and simply subtracting the two equations gives the Galerkin orthogonality, i.e.,

$$a[u - u_N, v_N] = 0, \quad \forall\, v_N \in V. \tag{3.2}$$

eqn:GO

[1] In general, the continuous inf-sup condition does not imply the discrete one.

If $a[\cdot,\cdot]$ is symmetric and coercive, then (3.2) means the error $u-u_N$ is orthogonal to V in the induced inner product from the bilinear form $a[\cdot,\cdot]$. Apparently, $\Pi_N u := u_N$ is a projection from $\mathscr{V}$ to V with respect to $(\cdot,\cdot)_{\mathcal{A}}$-inner product. It is oftentimes called the *Ritz projection*.

lem:Cea **Lemma 3.1** (Céa's Lemma)**.** *If the bilinear form $a[\cdot,\cdot]$ is continuous and coercive, then the Galerkin approximation u_N satisfies*

$$\|u-u_N\|_{\mathscr{V}} \leqslant \frac{C_a}{\alpha}\|u-v_N\|_{\mathscr{V}}, \quad \forall\, v_N \in V.$$

More generally, we have the following *quasi-optimality* or *quasi-best-approximation* of the finite-dimensional Galerkin approximation.

prop:optimal **Proposition 3.1** (Quasi-Optimality)**.** *Suppose $a[\cdot,\cdot]: \mathscr{V}\times\mathscr{V}\mapsto\mathbb{R}$ is continuous. The finite dimensional subspace V in the Galerkin approximation satisfies the discrete inf-sup condition (1.23) with $\alpha_N>0$. Let u and u_N be the exact solution of (1.16) and the Galerkin solution of (3.1), respectively. Then the error*

$$\|u-u_N\|_{\mathscr{V}} \leqslant \frac{\|\mathcal{A}\|}{\alpha_{_N}} \min_{w_N\in V}\|u-w_N\|_{\mathscr{V}}.$$

Proof. For all $w_N\in V$, applying (1.21) and (3.2), we have

$$\alpha_{_N}\|u_N-w_N\|_{\mathscr{V}} \leqslant \sup_{v_N\in V}\frac{a[u_N-w_N,v_N]}{\|v_N\|_{\mathscr{V}}} = \sup_{v_N\in V}\frac{a[u-w_N,v_N]}{\|v_N\|_{\mathscr{V}}} \leqslant \|\mathcal{A}\|\,\|u-w_N\|_{\mathscr{V}}.$$

Then simply applying the triangular inequality gives the estimate.

$$\|u-u_N\|_{\mathscr{V}} \leqslant \frac{\|\mathcal{A}\|+\alpha_{_N}}{\alpha_{_N}} \min_{w_N\in V}\|u-w_N\|_{\mathscr{V}}.$$

This constant in the upper bound is not sharp. The desired constant is obtained by Xu and Zikatanov [59]. □

rem:PGstab **Remark 3.2** (Stability)**.** In view of Theorem 1.4, we can see that the Galerkin solution depends on the data continuously, i.e.,

$$\|u_N\|_{\mathscr{V}} \leqslant \frac{1}{\alpha}\|f\|_{\mathscr{V}'}.$$

Finite element ⋆

The finite element method (FEM) has a long history in practical use and is widely applied to lots of problems in physics and engineering. It has been proved to be very successful in many areas, like structural mechanics. After decades of extensive development, the subject of classical (conforming) finite element method has become a well-understood and successful area in scientific computation. The most attractive feature of the FEM is its ability to handle complex geometries, boundaries, and operators with relative ease.

def:fe **Definition 3.1** (Finite element)**.** *A triple $(K,\mathcal{P},\mathcal{N})$ is called a finite element if and only if*

(i) $K\subseteq\mathbb{R}^d$ be a bounded closed set with nonempty interior and piecewise smooth boundary;

(ii) $\mathcal{P}$ be a finite-dimensional space of functions on K;

(iii) $\mathcal{N}=\{\mathcal{N}_1,\dots,\mathcal{N}_k\}$ be a basis of $\mathcal{P}'$.

We usually call K the *element domain*, $\mathcal{P}$ the space of *shape functions*, and $\mathcal{N}$ the set of *nodal variables*.

f:nodal_basis **Definition 3.2** (Nodal basis)**.** *Let $(K,\mathcal{P},\mathcal{N})$ be a finite element. The basis $\{\phi_j\}_{j=1,\dots,k}$ of $\mathcal{P}$ dual to $\mathcal{N}$, i.e., $\mathcal{N}_i(\phi_j)=\delta_{i,j}$ is called the nodal basis of $\mathcal{P}$.*

Example 3.1 (1D Lagrange element)**.** Let $K=[0,1]$, $\mathcal{P}$ be the set of linear polynomials, and $\mathcal{N}=\{\mathcal{N}_1,\mathcal{N}_2\}$ where $\mathcal{N}_1(v)=v(0)$ and $\mathcal{N}_2(v)=v(1)$. Then $(K,\mathcal{P},\mathcal{N})$ is a finite element and it is the well-known $\mathcal{P}_1$ Lagrange finite element discussed in Chapter 1. The nodal basis functions are $\phi_1(x)=1-x$ and $\phi_2(x)=x$. □

Remark 3.3 (Set of nodal variables)**.** If $\mathcal{P}$ is a k-dimensional space and $\{\mathcal{N}_1,\dots,\mathcal{N}_k\}\subset\mathcal{P}'$. Then condition (iii) in Definition 3.1 is equivalent to the unisolvence: For any $v\in\mathcal{P}$,

$$\mathcal{N}_i(v)=0,\ \ i=1,\dots,k \quad\Longrightarrow\quad v\equiv 0.$$

Now we describe the main steps of discretization using the $(K,\mathcal{P},\mathcal{N})$-finite element:

Step 1. Domain partitioning: Suppose K is simplex in $\mathbb{R}^d$. We first partition the physical domain into small subdomains. We discretize a polygonal domain Ω into small triangles or tetrahedrons τ. Let $h_\tau:=|\tau|^{\frac{1}{d}}$ be the diameter of $\tau\in\mathcal{M}$ and $h(x)$ be the local meshsize, that is the piecewise constant function with $h|_\tau:=h_\tau$ for all $\tau\in\mathcal{M}$. The collection $\mathcal{M}$ of elements is called a mesh or triangulation. We call $\mathcal{M}_h:=\mathcal{M}$ *quasi-uniform* if there exists a constant h independent of τ such that

$$h\lesssim h_\tau\lesssim h,\quad\forall\ \tau\in\mathcal{M}.$$

We will only consider *conforming* meshes, i.e., the intersection of any two elements in $\mathcal{M}$ is either an edge ($d=2$) / a face ($d=3$), vertex, or empty (see Figure 3.1 for an example). We denote by $G(\mathcal{M})$ the set of all grid points (vertices) in the mesh $\mathcal{M}$. And $\mathring{G}(\mathcal{M})\subseteq G(\mathcal{M})$ is the set of vertices except those on the Dirichlet boundary. Here we use the subscript h to describe the discrete nature and this does not imply the underlying meshes are quasi-uniform with meshsize h. In the future discussions, we will focus on uniform conforming meshes only.

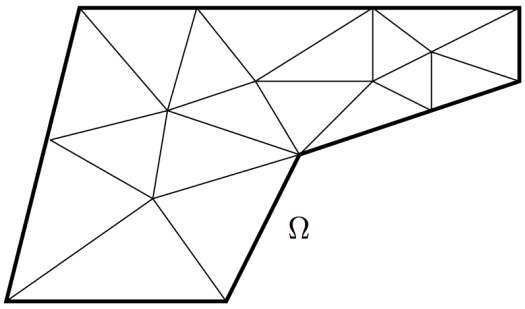


fig:domain

Figure 3.1: A polygonal domain Ω with conforming partition.

Step 2. Finite-dimensional approximation: Let $V_h \subset \mathscr{V}$ be the space of continuous piecewise polynomials over a quasi-uniform conforming mesh $\mathcal{M}_h$, which satisfies appropriate conditions on the boundary $\Gamma := \overline{\Omega}\backslash\Omega$, i.e.,

$$V_h := \left\{v \in C(\overline{\Omega}) \,:\, v|_\tau \in \mathcal{P}_\tau, \text{ for all } \tau \in \mathcal{M}_h\right\} \bigcap \mathscr{V}. \tag{3.3}$$ eq:u_h

We notice that there are many ways to approximate the continuous test function space. Different choices will then result in different numerical methods. In this section, we shall focus on the simplest case—linear finite element method on triangles or tetrahedrons, i.e., $v|_\tau$ is a linear polynomial on each $\tau \in \mathcal{M}_h$. The weak form of the finite element approximation reads: Find $u_h \in V_h$ such that

$$a[u_h, v_h] = \langle f, v_h\rangle, \quad \forall\ v_h \in V_h, \tag{3.4}$$ eq:ell_w

or, equivalently,

$$\mathcal{A}_h u_h = f_h. \tag{3.5}$$ eq:ell_s

Step 3. Assembling the finite-dimensional problem: Using the finite element definition $(K, \mathcal{P}, \mathcal{N})$, we can give a basis of the finite dimensional approximation space V_h. Suppose $\{\phi_i\}_{i=1}^N$ be a basis of the N-dimensional space V_h. Then (3.5) can be written as an linear algebraic equation

$$\hat{\mathcal{A}}_h \underline{u_h} = \vec{f}_h. \tag{3.6}$$ eq:ell_dis

We are going to discuss this notation later in §3.2.

Some properties of finite element methods ⋆

There are a few important properties of finite element space and method that will become crucial for our later analysis for multilevel iterative methods.

prop:interp **Proposition 3.2** (Interpolation error). *Let $\mathcal{M}_h$ be a uniform mesh and V_h be a C^α ($\alpha \geqslant 0$) finite element space on $\mathcal{M}_h$. The interpolant $\mathcal{J}_h : W_p^m(\Omega) \mapsto V_h$ satisfies*

$$\|v - \mathcal{J}_h v\|_{W_p^k(\Omega)} \lesssim h^{m-k}\|v\|_{W_p^m(\Omega)}, \quad \forall\, v \in W_p^m(\Omega),\ 0 \leqslant k \leqslant \min\{m, \alpha+1\}.$$

prop:inverse **Proposition 3.3** (Inverse estimate). *Let $\mathcal{M}_h$ be a uniform mesh and $\mathcal{P} \subseteq W_p^k(K) \bigcap W_q^m(K)$ and $0 \leqslant m \leqslant k$. If V_h is a finite element space for $(K, \mathcal{P}, \mathcal{N})$ on $\mathcal{M}_h$, then we have*

$$\Big(\sum_{\tau\in\mathcal{M}_h} \|v\|^p_{W_p^k(\tau)}\Big)^{\frac{1}{p}} \lesssim h^{m-k+\min\{0,\frac{d}{p}-\frac{d}{q}\}}\Big(\sum_{\tau\in\mathcal{M}_h} \|v\|^q_{W_q^m(\tau)}\Big)^{\frac{1}{q}}, \quad \forall\, v \in V_h.$$

rem:rho-FE **Remark 3.4** (Spectral radius and condition number of $\mathcal{A}_h$). Suppose that we have a uniform partition with meshsize h. It is clear, from the Poincaré inequality and the inverse inequality, that

$$\|v\|_0^2 \lesssim (\mathcal{A}_h v, v) = (\nabla v, \nabla v) = |\!|\!| v |\!|\!|^2 \lesssim h^{-2}\|v\|_0^2, \quad \forall v \in V_h.$$

In fact, we have $\rho(\mathcal{A}_h) \cong h^{-2}$ and $\kappa(\mathcal{A}_h) \cong h^{-2}$.

Using Proposition 3.3, we can easily see that, for any $v \in V_h$,

$$\begin{aligned}
&\|v\|_{L^\infty(\Omega)} \lesssim h^{-\frac{d}{p}}\|v\|_{L^p(\Omega)}, && p \in [1,\infty);\\
&\|v\|_{H^s(\Omega)} \lesssim h^{-s}\|v\|_{L^2(\Omega)}, && s \in [0,1];\\
&\|v\|_{H^{1+\sigma}(\Omega)} \lesssim h^{-\sigma}\|v\|_{H^1(\Omega)}, && \sigma \in (0,\tfrac{1}{2}).
\end{aligned}$$

Moreover, there is a discrete Sobolev inequality at the bottom-line case (when $d = 2$) which is worthy for special attention.

nfty-estimate **Proposition 3.4** (Discrete Sobolev inequality [15]). *The following inequality holds*

$$\|v\|_{L^\infty(\Omega)} \lesssim C_d(h)\|v\|_{H^1(\Omega)}, \quad \forall\, v \in V_h,$$

where $C_1(h) \equiv 1$, $C_2(h) = |\log h|^{1/2}$, and $C_3(h) = h^{-\frac{1}{2}}$.

prop:L2proj **Proposition 3.5** (L^2-projection [15]). *Define $\mathcal{Q}_h : L^2(\Omega) \mapsto V_h$ by*

$$(\mathcal{Q}_h v, w) = (v, w), \quad \forall\, v \in L^2(\Omega), w \in V_h.$$

Then we have the following weighted L^2-estimate

$$\|v - \mathcal{Q}_h v\|_0 + h|\!|\!|\mathcal{Q}_h v|\!|\!|_1 \lesssim h|\!|\!| v |\!|\!|_1.$$

Error analysis ⋆

We now briefly introduce standard error estimates for the continuous linear finite element; see [25, 20] for details. For standard finite element approximation of elliptic equations, the most important property is the following Galerkin orthogonality property (see Remark 3.1)

$$a[u - u_h, v_h] = 0, \quad \forall\ v_h \in V.$$

Using the definition of the energy norm $\|\!|\cdot|\!\| := a[\cdot,\cdot]^{1/2}$, the Galerkin orthogonality (3.2), and the Cauchy-Schwarz inequality, we have

$$\|\!|u - u_h|\!\|^2 = a[u - u_h, u - u_h] = a[u - u_h, u - v_h] \leqslant \|\!|u - u_h|\!\|\, \|\!|u - v_h|\!\|, \quad \forall v_h \in V.$$

Hence, we obtain the *optimality* of the finite element approximation, i.e.,

$$\|\!|u - u_h|\!\| \leqslant \inf_{v_h \in V} \|\!|u - v_h|\!\|. \tag{3.7}$$ eqn:optimal

This means u_h is the best approximation of u in the subspace V. In general, it is not true for finite element approximations.

thm:H1error **Theorem 3.1** (H^1-error estimate)**.** *If $u \in H_0^m(\Omega)$ $(1 < m \leqslant 2)$, its $\mathcal{P}_1$ Lagrange finite element approximation $u_h \in V_h \subset \mathscr{V} = H_0^1(\Omega)$ satisfies*

$$\|u - u_h\|_{1,\Omega} \lesssim h^{m-1}|u|_{m,\Omega}.$$

If $m = 2$, then we have $\|u - u_h\|_{1,\Omega} \lesssim h\|f\|_{0,\Omega}$.

thm:L2error **Theorem 3.2** (L^2-error estimate)**.** *If $u \in H_0^2(\Omega)$, its $\mathcal{P}_1$ Lagrange finite element approximation $u_h \in V_h \subset \mathscr{V} = H_0^1(\Omega)$ satisfies*

$$\|u - u_h\|_{0,\Omega} \lesssim h|u - u_h|_{1,\Omega} \lesssim h^2|u|_{2,\Omega} \lesssim h^2\|f\|_{0,\Omega}.$$

Remark 3.5 (A posteriori error analysis)**.** A posteriori error estimation relies on the following error equation (or residual equation):

$$a[u - u_h, v] = a[u, v] - a[u_h, v] = \langle f, v\rangle - a[u_h, v] = \langle f - \mathcal{A}u_h, v\rangle, \quad \forall v \in \mathscr{V}.$$

Hence, by the Cauchy-Schwarz inequality, we obtain (see HW 3.1)

$$\|\!|f - \mathcal{A}u_h|\!\|_* \lesssim \|\!|u - u_h|\!\| \lesssim \|\!|f - \mathcal{A}u_h|\!\|_*. \tag{3.8}$$ eqn:err-res-e

Here $\|\!|\cdot|\!\|_*$ is the dual norm of $\|\!|\cdot|\!\|$. Notice that, on the right-hand side, we only have the data f and the discrete solution u_h. This upper bound does not depend on the unknown solution u. Of course, to make the upper bound useful in adaptive algorithms, we need it to be local and computable.

3.2 Matrix representations

sec:matrixform

In the previous chapters, we have written the discrete problem simply as

$$A\vec{u} = \vec{f},$$

which is actually abuse of notation. Now we would like to clarify (especially for finite element methods) the relation between the general operator form $\mathcal{A}_h u_h = f_h$ and its often-used matrix form (3.6), i.e., $\hat{\mathcal{A}}_h \underline{u_h} = \vec{f_h}$.

Vector and matrix representations

Assume that $\{\phi_i\}_{i=1,\dots,N}$ is a basis of V. Any function $v \in V$ can be represented as

$$v = \sum_{i=1}^{N} \underline{v}_i \phi_i$$

and the vector representation (coefficient vector) of v is defined as

$$\underline{v} := \begin{pmatrix} \underline{v}_1 \\ \vdots \\ \underline{v}_N \end{pmatrix} \in \mathbb{R}^N. \tag{3.9}$$

eqn:v-matrix

It is not hard to notice that there is another natural and easier-to-compute vector representation

$$\vec{v} := \begin{pmatrix} (v,\phi_1) \\ (v,\phi_2) \\ \vdots \\ (v,\phi_N) \end{pmatrix} \quad \text{and} \quad \vec{v} = M\underline{v}, \tag{3.10}$$

eqn:v-matrix2

where $M \in \mathbb{R}^{N\times N}$ with $M_{i,j} := (\phi_j,\phi_i) = (\phi_i,\phi_j)$ is the *mass matrix*. $\underline{v}$ and $\vec{v}$ are sometimes called the *primal and dual vector representations* of v, respectively. Apparently, we have

$$(\underline{u},\vec{v})_{l^2} = \underline{u}^T M \underline{v} = (u,v)_V.$$

Suppose W is another finite-dimensional linear space with a basis $\{\psi_i\}_{i=1,\dots,N'}$. In general, W could be of different dimension than V, namely, $N' \neq N$. For any linear operator $\mathcal{A} : V \mapsto W$, we give a matrix representation (the so-called primal representation), $\underline{\mathcal{A}} \in \mathbb{R}^{N'\times N}$, such that it satisfies that $\sum_{i=1}^{N'} \left(\underline{\mathcal{A}}\right)_{i,j} \psi_i = \mathcal{A}\phi_j \ \ (j = 1,\dots,N)$, i.e.,

$$\mathcal{A}(\phi_1,\dots,\phi_N) = (\psi_1,\dots,\psi_{N'})\underline{\mathcal{A}}. \tag{3.11}$$

eqn:A-matrix

On the other hand, the dual representation (the *stiffness matrix*) corresponding to $\mathcal{A}$ is denoted by $\hat{\mathcal{A}} \in \mathbb{R}^{N\times N}$ with entries $\left(\hat{\mathcal{A}}\right)_{i,j} := (\mathcal{A}\phi_j,\phi_i)$.

It is not difficult to check the statements in the following identities; see HW 3.2.

lem:mat-form **Lemma 3.2** (Matrix representations)**.** *If $\mathcal{A}, \mathcal{B} : V \mapsto V$ and $v, u \in V$, we have the following results:*

1. $\underline{\mathcal{A}\mathcal{B}} = \underline{\mathcal{A}}\,\underline{\mathcal{B}}$;
2. $\underline{\mathcal{A}v} = \underline{\mathcal{A}}\,\underline{v}$;
3. $\sigma(\mathcal{A}) = \sigma(\underline{\mathcal{A}})$, $\kappa(\mathcal{A}) = \kappa(\underline{\mathcal{A}})$;
4. $\overrightarrow{\mathcal{A}v} = \hat{\mathcal{A}}\,\underline{v}$, $\vec{v} = M\,\underline{v}$;
5. $\hat{\mathcal{A}} = M\,\underline{\mathcal{A}}$;
6. $(u, v) = (M\underline{u}, \underline{v})$.

ex:identity **Example 3.2** (Identity operator)**.** Let $\mathcal{I} : V \mapsto V$ be the identity operator. Its stiffness matrix and mass matrix are equal to each other, i.e., $\hat{\mathcal{I}} = M$. Hence $\underline{\mathcal{I}} = M^{-1}\hat{\mathcal{I}} = I$. As a consequence, we have

$$I = \underline{\mathcal{I}} = \underline{\mathcal{A}\mathcal{A}^{-1}} = \underline{\mathcal{A}}\,\underline{\mathcal{A}^{-1}},$$

which shows $\underline{\mathcal{A}}^{-1} = \underline{\mathcal{A}^{-1}}$. □

Example 3.3 (Finite difference matrices)**.** For the finite difference methods, we can simply let $\mathcal{A} : \mathbb{R}^N \mapsto \mathbb{R}^N$ be a matrix and the basis $\phi_i = \vec{e}_i := (0, \ldots, 1, \ldots, 0)^T \in \mathbb{R}^N$, then we have $\mathcal{A} = \underline{\mathcal{A}} = \hat{\mathcal{A}} = A$. □

Remark 3.6 (Matrix representation with orthonormal basis)**.** More generally speaking, if $\mathcal{A} : V \mapsto V$ and $\{\phi_i\}_{i=1}^N$ is an orthonormal basis of V, then we have $M = I$ and $\hat{\mathcal{A}} = \underline{\mathcal{A}}$.

Finite element matrices

We now use a few simple examples to demonstrate how to apply these notations. Suppose that $V = V_h$ is the piecewise linear finite element space and $\{\phi_i\}_{i=1,\ldots,N}$ are the canonical basis functions. Let A be the resulting coefficient matrix of (3.1) with $(A)_{i,j} = a_{i,j} := a[\phi_i, \phi_j]$. By definition, $A = \hat{\mathcal{A}} \in \mathbb{R}^{N\times N}$ is the stiffness matrix corresponding to $\mathcal{A}$. Let $\underline{u} = \big(u_i\big)_{i=1}^N \in \mathbb{R}^N$ be the vector of coefficients of u_h. Let $\vec{f} = \big(f_i\big)_{i=1}^N := \big\{\langle f, \phi_i\rangle\big\}_{i=1}^N$. Then $\underline{u}$ satisfies the linear system of equations:

$$\hat{\mathcal{A}}\underline{u} = \vec{f} \qquad \text{or} \qquad A\underline{u} = \vec{f}.$$

Since we are going to focus on the finite element discretization from now on, we denote $A := \hat{\mathcal{A}}$, when there is no ambiguity arises. Upon solving this finite-dimensional problem, we obtain a discrete approximation $u_h = \sum_{i=1}^N \underline{u}_i\phi_i$.

The main algebraic properties for the stiffness matrix is that: A is *sparse* with $O(N)$ nonzeros, symmetric positive definite (for Dirichlet or mixed boundary condition problems) or symmetric positive semi-definite (for Neumann boundary condition problems). We now summarize this brief introduction of finite element methods with a few comments. The following results are valid for a large class of finite elements for second-order elliptic boundary value problems in general domains.

rem:SpecMass **Remark 3.7** (Spectrum of mass matrix)**.** Another often used matrix is the mass matrix $M \in \mathbb{R}^{N\times N}$, in which $M_{i,j} = (\phi_i, \phi_j)$. It is well-known that the mass matrix is also SPD and well-conditioned, i.e.,

$$h^d \|\xi\|_0^2 \lesssim \xi^T M\, \xi \lesssim h^d \|\xi\|_0^2, \quad \forall\, \xi \in \mathbb{R}^N.$$

In fact, we know that

$$(M\underline{v}, \underline{v}) = \sum_{i,j} \underline{v}_i\, \underline{v}_j\, (\phi_i, \phi_j) = (v, v) = \int_\Omega v^2(x)\, dx \cong h^d \sum_i \underline{v}_i^2 \cong h^d (\underline{v}, \underline{v}). \tag{3.12}$$ eqn:spec-M

rem:SpecStiff **Remark 3.8** (Spectrum of stiffness matrix)**.** Suppose that we have a uniform partition with meshsize h. It is well-known that the stiffness matrix A is SPD and, from Remark 3.4,

$$h^d \|\xi\|_0^2 \lesssim \xi^T A\, \xi \lesssim h^{d-2} \|\xi\|_0^2, \quad \forall\, \xi \in \mathbb{R}^N.$$

Hence the condition number $\kappa(A) \cong h^{-2}$ and the CG method becomes very slow when h decreases.

Simple iterators in matrix form

Now we consider the solution of the standard $\mathcal{P}_1$ Lagrange finite element for the Poisson's equation, i.e., $\hat{\mathcal{A}}\underline{u} = \vec{f}$. The simplest iterative solver for this finite element equation is probably the well-known Richardson method:

$$\underline{u}^{\text{new}} = \underline{u}^{\text{old}} + \omega\Big(\vec{f} - \hat{\mathcal{A}}\, \underline{u}^{\text{old}}\Big). \tag{3.13}$$ eqn:Richardso

It is equivalent to

$$\underline{u}^{\text{new}} = \underline{u}^{\text{old}} + \omega\Big(M\underline{f} - M\underline{\mathcal{A}}\, \underline{u}^{\text{old}}\Big) = \underline{u}^{\text{old}} + \omega M\Big(\underline{f} - \underline{\mathcal{A}}\, \underline{u}^{\text{old}}\Big).$$

That is to say, the Richardson method, in the operator form, can be written as

$$u^{\text{new}} = u^{\text{old}} + \mathcal{B}_\omega\Big(f - \mathcal{A}\, u^{\text{old}}\Big)$$

with an iterator $\mathcal{B}_\omega$, whose matrix representation is $\underline{\mathcal{B}_\omega} = \omega M$. Therefore, it is easy to check (HW 3.3) that the operator form of the Richardson method is

$$\mathcal{B}_\omega v := \omega \sum_{i=1}^N (v, \phi_i)\phi_i, \quad \forall v \in V \quad \Longleftrightarrow \quad \underline{\mathcal{B}_\omega} = \omega M. \tag{3.14}$$ eqn:Richardso

If we choose $\omega = 1$ in the Richardson method, then we get $\mathcal{B}_1 v = \sum_{i=1}^{N}(v,\phi_i)\phi_i$. This implies

$$(\mathcal{B}_1 v, v) = \sum_{i=1}^{N}(v,\phi_i)^2 = \sum_{i=1}^{N}(M\underline{v})_i^2 = (M\underline{v}, M\underline{v}) = (M^2\underline{v}, \underline{v}).$$

Since M is SPD, we get

$$(M^2\underline{v}, \underline{v}) = (MM^{\frac{1}{2}}\underline{v},\, M^{\frac{1}{2}}\underline{v}) \cong h^d(M^{\frac{1}{2}}\underline{v},\, M^{\frac{1}{2}}\underline{v}) = h^d(M\underline{v},\, \underline{v}).$$

The estimate (3.12) implies that

$$(\mathcal{B}_1 v, v) \cong h^d(v, v). \tag{3.15}$$ eqn:Richardso

In general, a smoother or local relaxation is just a stationary iterative method

$$u^{\text{new}} = u^{\text{old}} + \mathcal{S}(f - \mathcal{A}u^{\text{old}})$$

and its matrix representation is

$$\underline{u}^{\text{new}} = \underline{u}^{\text{old}} + \underline{\mathcal{S}}(M^{-1}\vec{f} - M^{-1}\hat{\mathcal{A}}\underline{u}^{\text{old}}) = \underline{u}^{\text{old}} + \underline{\mathcal{S}}M^{-1}(\vec{f} - \hat{\mathcal{A}}\underline{u}^{\text{old}}). \tag{3.16}$$ eqn:MatS

The above equality indicates that, in the matrix form, we shall define a smoother as

$$S := \underline{\mathcal{S}}M^{-1}. \tag{3.17}$$ eqn:MatRepSmo

Example 3.4 (Matrix form of the Richardson iteration)**.** For example, if we consider the above Richardson method (3.14) as an example, then

$$S_{\text{R}} = \underline{\mathcal{S}_{\text{R}}}M^{-1} = \underline{\mathcal{B}_\omega}M^{-1} = \omega I.$$

This coincides with the algebraic form of the Richardson method (3.13). □

Let $w := \mathcal{S}^T u$. Then we have

$$\vec{w} = \Big((\mathcal{S}^T u, \phi_i)\Big)_{i=1}^{N} = \Big(\sum_j \underline{u}_j(\mathcal{S}^T\phi_j, \phi_i)\Big)_{i=1}^{N} = \Big(\sum_j \underline{u}_j(\phi_j, \mathcal{S}\phi_i)\Big)_{i=1}^{N} = \big(\hat{\mathcal{S}}\big)^T\underline{u}.$$

This immediately gives

$$\underline{\mathcal{S}^T}\,\underline{u} = \underline{\mathcal{S}^T}\underline{u} = \underline{w} = M^{-1}\vec{w} = M^{-1}\big(\hat{\mathcal{S}}\big)^T\underline{u} = M^{-1}\big(M\underline{\mathcal{S}}\big)^T\underline{u}.$$

In turn, it shows

$$\underline{\mathcal{S}^T} = M^{-1}\big(M\underline{\mathcal{S}}\big)^T = M^{-1}\underline{\mathcal{S}}^T M = S^T M. \tag{3.18}$$ eqn:MatST

By definition of the primal matrix representation of an operator, we have

$$\mathcal{S}(\phi_1,\ldots,\phi_N) = (\phi_1,\ldots,\phi_N)\underline{\mathcal{S}} \qquad \text{and} \qquad \mathcal{S}^{-1}(\phi_1,\ldots,\phi_N) = (\phi_1,\ldots,\phi_N)\underline{\mathcal{S}^{-1}}.$$

Using Example 3.2, it is easy to see that

$$\underline{\mathcal{S}^{-1}} = \big(\underline{\mathcal{S}}\big)^{-1} = \big(SM\big)^{-1} = M^{-1}S^{-1}. \tag{3.19}$$ eqn:MatSinv

Using (3.17)–(3.19) and the definition of symmetrized operator (2.10), we can obtain the matrix form of the symmetrization

$$\begin{aligned} \overline{S} &= \overline{\underline{\mathcal{S}}}M^{-1} = S^T M\big(M^{-1}S^{-T} + M^{-1}S^{-1} - M^{-1}\hat{A}\big)SMM^{-1} \\ &= S^T\big(S^{-T} + S^{-1} - A\big)S, \end{aligned} \tag{3.20}$$

which is consistently with the definition of symmetrization.

3.3 Smoothers and smoothing effect

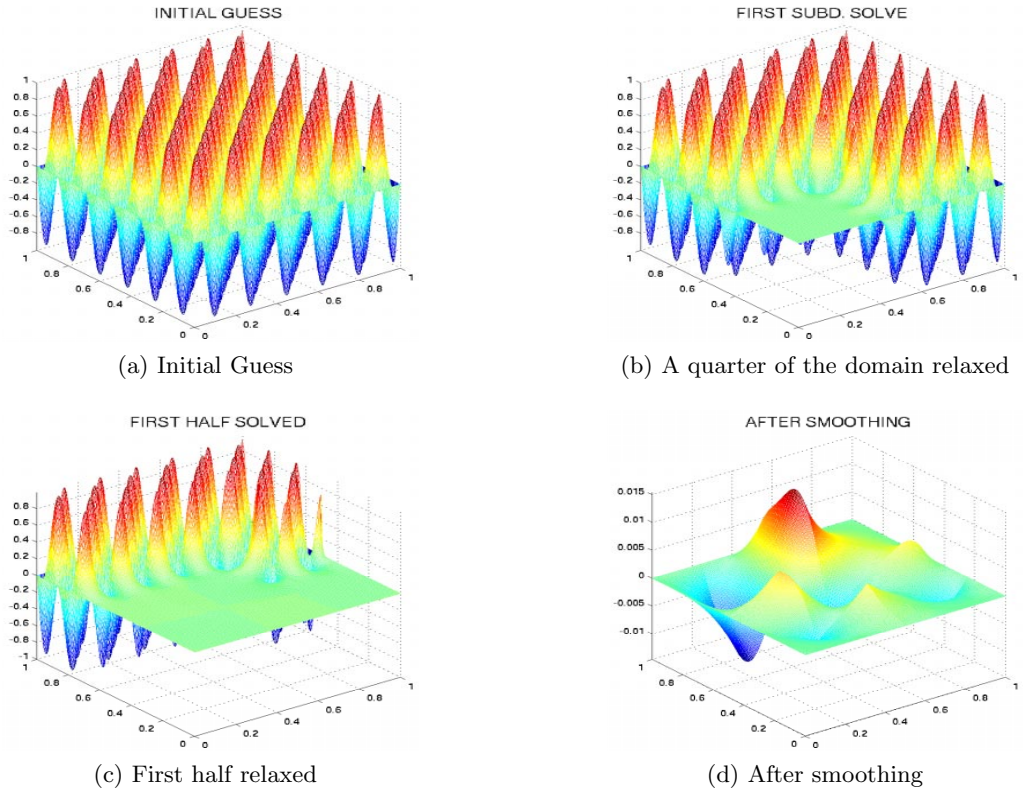
sec:smoother

The methods discussed by far, for example the damped Jacobi and Gauss–Seidel methods, are mostly *local relaxation* methods. Other similar methods include the SOR method and incomplete LU factorizations. The name "local relaxation" comes from the fact that these methods just correct the residual vector locally at one nodal point at a time; see Example 2.4. Although these methods are not very efficient as a solver by themselves, they are still key ingredients of modern multilevel iterative methods. These methods can be applied to reduce high-frequency error components. In this section, we analyze their smoothing effect in several different approaches.

A numerical example

The damped Jacobi and Gauss–Seidel methods are often called local relaxations. This relaxation procedure is effective to the error components that are local in nature. Therefore, it is not surprising both the damped Jacobi and the Gauss–Seidel methods can damp out non-smooth components more easily. These methods are very inefficient for relatively smoother components in the error since a smoother function is more globally related.

We have observed that the basic stationary linear iterative schemes converge rather fast in the very beginning but then slows down after a few step; see Figure 1.4 for the convergence behavior of the damped Jacobi method. Moreover, these methods not only converges fast in the first few steps, but also smooth out the error function very quickly. In other words, the error becomes a much smoother function after a few iterations. This property of the iterative scheme is naturally called smoothing property and an iterative scheme having this smoothing property is called a *smoother*. Figure 3.2 is a pictorial example for applying multiplicative overlapping domain decomposition method with four subdomains. We can see that, after one iteration, the method smoothes out the high frequency part and leaves the lower frequency part



(a) Initial Guess (b) A quarter of the domain relaxed

(c) First half relaxed (d) After smoothing

Figure 3.2: Iterative method in the viewpoint of subdomain relaxation.

basiciteration

behind. In fact, basic linear relaxation schemes, such as the Richardson, Jacobi, and Gauss–Seidel iterations, are local and can only capture high frequency (local) part of the error, but do not work well on low frequency (global) part.

Local Fourier analysis ⋆

In order to analyze the local behavior of iterative methods, we consider the 2D Poisson's equation with homogenous Dirichlet boundary condition on the unit square discretized with a uniform triangulation; see §1.2.

We first analyze the damped Jacobi method. Using the local Fourier analysis [17], we have the following steps:

1. The standard FD stencil can be written as

$$4u_{i,j} - \big(u_{i-1,j} + u_{i+1,j} + u_{i,j-1} + u_{i,j+1}\big) = h^2 f_{i,j}, \quad i,j = 1,\ldots,n$$

and the damped Jacobi (or Richardson) method for the above equation reads

$$u_{i,j}^{\text{new}} = (1-\omega)u_{i,j}^{\text{old}} + \frac{\omega}{4}\big(u_{i-1,j}^{\text{old}} + u_{i+1,j}^{\text{old}} + u_{i,j-1}^{\text{old}} + u_{i,j+1}^{\text{old}}\big) + \frac{\omega}{4}h^2 f_{i,j}, \quad i,j = 1,\ldots,n.$$

2. Define the discrete error function $e_{i,j}^{\text{new}} := u_{i,j} - u_{i,j}^{\text{new}}$ and $e_{i,j}^{\text{old}} := u_{i,j} - u_{i,j}^{\text{old}}$, for $i, j = 1, \ldots, n$. It is clear that the error function satisfies the local error equation

$$e_{i,j}^{\text{new}} = (1-\omega)e_{i,j}^{\text{old}} + \frac{\omega}{4}\big(e_{i-1,j}^{\text{old}} + e_{i+1,j}^{\text{old}} + e_{i,j-1}^{\text{old}} + e_{i,j+1}^{\text{old}}\big), \quad i, j = 1, \ldots, n.$$

3. Apply the discrete Fourier transformation:

$$e_{i,j} = \sum_{\theta \in \Theta_n} \alpha_\theta e^{\sqrt{-1}(i\theta_1 + j\theta_2)} \text{ and } \Theta_n := \Big\{(\theta_1, \theta_2) \,:\, \theta_1 = \frac{2k\pi}{n},\ \theta_2 = \frac{2l\pi}{n},\ \ k, l \in [-m_1, m_2]\Big\},$$

where $m_1 = n/2 - 1, m_2 = n/2$, if n is even and $m_1 = m_2 = (n-1)/2$, if n is odd. Plugging the discrete Fourier transforms of $e_{i,j}^{\text{new}}$ and $e_{i,j}^{\text{old}}$ to the above error equation, we get the amplification factor of the local mode $e^{\sqrt{-1}(i\theta_1 + j\theta_2)}$

$$\lambda(\theta) := \frac{\alpha_\theta^{\text{new}}}{\alpha_\theta^{\text{old}}} = 1 - \omega\Big(1 - \frac{\cos(\theta_1) + \cos(\theta_2)}{2}\Big) \leqslant 1.$$

Furthermore, $\lambda(\theta) \to 1$ when $|\theta| \to 0$ (low-frequency components).

4. Asymptoticly, $m_1 \approx m_2 \approx \frac{n}{2}$. So we can define a smoothing factor (i.e. maximal amplification factor corresponding to high-frequency local modes) by

$$\bar{\rho} := \sup_\theta \Big\{\big|\lambda(\theta)\big| \,:\, \frac{\pi}{2} \leqslant |\theta_k| \leqslant \pi,\ k = 1, 2\Big\}.$$

By plugging in the end points, we get the the smoothing factor for the damped Jacobi method is

$$\bar{\rho}_{\text{Jacobi}} := \max\Big\{\Big|1 - 2\omega\Big|, \Big|1 - \frac{1}{2}\omega\Big|, \Big|1 - \frac{3}{2}\omega\Big|\Big\}.$$

Remark 3.9 (Optimal damping factor for smoothing)**.** We notice that, if $\omega = 1$ (the Jacobi method), then $\bar{\rho}_{\text{Jacobi}} = 1$. This confirms the result we obtained in the previous subsection. Apparently, the "best" weight that minimizes the smoothing factor is $\omega = 4/5$, which leads to $\bar{\rho}_{\text{Jacobi}} = 3/5$.

It is not hard to imagine that the G-S method should have better smoothing property than the Jacobi method. Using the same steps as above, we have:

1. The G-S method in lexicographical order reads

$$u_{i,j}^{\text{new}} = \frac{1}{4}\big(u_{i-1,j}^{\text{new}} + u_{i+1,j}^{\text{old}} + u_{i,j-1}^{\text{new}} + u_{i,j+1}^{\text{old}}\big) + \frac{1}{4}h^2 f_{i,j}, \quad i, j = 1, \ldots, n.$$

2. The discrete error function satisfies

$$e_{i,j}^{\text{new}} = \frac{1}{4}\big(e_{i-1,j}^{\text{new}} + e_{i+1,j}^{\text{old}} + e_{i,j-1}^{\text{new}} + e_{i,j+1}^{\text{old}}\big), \quad i, j = 1, \ldots, n.$$

3. Apply the discrete Fourier transform and compute the amplification factor

$$\lambda(\theta) := \frac{\alpha_\theta^{\text{new}}}{\alpha_\theta^{\text{old}}} = \frac{e^{\sqrt{-1}\theta_1} + e^{\sqrt{-1}\theta_2}}{4 - e^{-\sqrt{-1}\theta_1} - e^{-\sqrt{-1}\theta_2}}.$$

4. One can show the smoothing factor for the G-S method is

$$\bar{\rho}_{\text{GS}} := \left|\lambda\Big(\frac{\pi}{2}, \arccos(4/5)\Big)\right| = \frac{1}{2}.$$

Remark 3.10 (Anisotropic problems and smoothing effect)**.** Notice that the above analysis only works for uniform partition and isotropic coefficients. When we solve an anisotropic problem, it is important to note that the G-S method (and other point relaxation methods) yields not as good smoothing factor as the isotropic case. In fact, the smoothing factor goes to 1 when ratio between small and large coefficients goes to 0; see Chapter 6 for details.

Remark 3.11 (Ordering and smoothing effect)**.** For the G-S method, ordering is important. When using the red-black ordering, one can show the smoothing factor is $\bar{\rho}_{\text{RBGS}} = \frac{1}{4}$. This means the smoothing effect of the red-black ordering is better.

Smoother analysis

Consider the Richardson method (3.14) with weight $\omega = h^{2-d}$, i.e.,

$$\mathcal{S}_{\text{R}} v := h^{2-d} \sum_{i=1}^{N} (v, \phi_i)\phi_i, \quad \forall v \in V. \tag{3.21}$$ eqn:Richardso

In view of (3.15) and using the fact that the spectral radius of the FE operator is $\rho(\mathcal{A}) \cong h^{-2}$ (see Remark 3.4), we find

$$(\mathcal{S}_{\text{R}} v, v) = h^{2-d}(\vec{v}, \vec{v}) \cong h^2(v, v) \cong \frac{1}{\rho(\mathcal{A})}(v, v). \tag{3.22}$$ eqn:smoothing

This is the property we will ask for a smoother later on. Roughly speaking, $\mathcal{S}_{\text{R}}$ behaves like $\mathcal{A}^{-1}$ in the high-frequency part.

Apparently, the damped Jacobi method also satisfies this condition. In fact, using the standard scaling argument on each element, we can see that

$$h^{d-2}(\xi, \xi) \lesssim (D\xi, \xi) \lesssim h^{d-2}(\xi, \xi).$$

Hence, we have the Jacobi smoother

$$(\mathcal{S}_{\text{J}} v, v) = (M\underline{\mathcal{S}_{\text{J}} v}, \underline{v}) = (MD^{-1}M\underline{v}, \underline{v}) \cong h^{d+2}(\underline{v}, \underline{v}) \cong h^2(v, v) \cong \frac{1}{\rho(\mathcal{A})}(v, v).$$

Remark 3.12 (Smoother v.s. preconditioner)**.** From the property (3.25) of these popular smoothers, we can easily see that

$$\rho_{\mathcal{A}}^{-1}(v,v) \lesssim (\mathcal{S}v,v) \lesssim \rho_{\mathcal{A}}^{-1}(v,v), \tag{3.23}$$ eqn:smoothing

where $\rho_{\mathcal{A}} := \rho(\mathcal{A})$. Thus we have a lower bound for the minimal eigenvalue $\rho_{\mathcal{A}}^{-1} \lesssim \lambda_{\min}(\mathcal{S})$. With simple manipulations, we also get

$$\rho_{\mathcal{A}}^{-1}(v,v)_{\mathcal{A}} \lesssim \rho_{\mathcal{A}}^{-1}(\mathcal{A}v,\mathcal{A}v) \lesssim (\mathcal{S}\mathcal{A}v,v)_{\mathcal{A}} \lesssim \rho_{\mathcal{A}}^{-1}(\mathcal{A}v,\mathcal{A}v) \leqslant (v,v)_{\mathcal{A}}. \tag{3.24}$$ eqn:smoothing

Due to Lemmas 2.5 and 2.6, (3.24) indicates that $\kappa(\mathcal{S}\mathcal{A}) \lesssim \rho(\mathcal{A}) \cong \kappa(\mathcal{A})$, which means these smoothers might not improve the condition number by themselves.

Next, we shall show that the G-S method behaves in a similar way.

:GS-algebraic **Lemma 3.3** (Smoothing property of G-S in matrix form)**.** *Let $\hat{\mathcal{A}}$ be the stiffness matrix and $\hat{\mathcal{A}} = A = D + L + U$. Then the G-S method satisfies*

$$\|(D+L)\xi\|_0 \cong \|D\xi\|_0 \cong h^{d-2}\|\xi\|_0, \quad \forall \xi \in \mathbb{R}^N.$$

Proof. Locality of the nodal basis functions leads to sparse matrix L; in turn, this gives

$$\|(D+L)\xi\|_0 \lesssim \|D\xi\|_0 \lesssim h^{d-2}\|\xi\|_0.$$

The other direction follows from

$$h^{d-2}\|\xi\|_0^2 \lesssim \big(D\xi,\xi\big) \leqslant \big((D+A)\xi,\xi\big) = 2\big((D+L)\xi,\xi\big) \lesssim \|(D+L)\xi\|_0\,\|\xi\|_0.$$

We then get the desired estimates with simple manipulations. □

:SGS-operator **Lemma 3.4** (Smoothing property of SGS)**.** *Let $\mathcal{S} : V \mapsto V$ be the symmetrized G-S (SGS) iterator. Then we have*

$$\big(\mathcal{S}v,v\big) \cong h^2(v,v) \cong \frac{1}{\rho(\mathcal{A})}(v,v). \tag{3.25}$$ eqn:smoother-

Proof. The matrix form of SGS is

$$\underline{\mathcal{S}} = SM = (D+U)^{-1}D(D+L)^{-1}M.$$

Let $\underline{v}$ be the vector representation of $v \in V$. Then we have

$$(\mathcal{S}v,v) = (M\underline{\mathcal{S}v},\underline{v}) = (M\underline{\mathcal{S}}\,\underline{v},\underline{v}) = \|D^{\frac{1}{2}}(D+L)^{-1}M\underline{v}\|_0^2.$$

Hence to show the lemma is equivalent to prove that

$$\|D^{\frac{1}{2}}(D+L)^{-1}M\underline{v}\|_0^2 \cong h^2(M\underline{v},\underline{v}).$$

By changing of variable $\xi := (D+L)^{-1}M\underline{v} \in \mathbb{R}^N$ and the fact $M \cong h^d$, we reduce the above equality to

$$h^{d-2}(D\xi,\xi) \cong h^{2(d-2)}\|\xi\|_0^2 \cong \|(D+L)\xi\|_0^2 = (M\underline{v}, M\underline{v}), \quad \forall\, \xi \in \mathbb{R}^N,$$

which is true due to Lemma 3.3. □

3.4 A two-grid method

rid-algorithm

From the analysis in §3.3, we have found that local relaxation methods (smoothers) can damp the oscillatory components of the error quickly. Motivated by the two-level DD method in §2.4, we can introduce coarser levels to take care of the smooth components. A natural idea is then, after a few smoothing steps, to approximate the resulting problem on a coarser grid and continue the iteration with a "coarse version" of the problem. The main idea is to resolve the high frequency part of the error with relaxation schemes and leave the low frequency part to the coarse levels. Before we discuss multilevel methods, we first investigate a much simpler case, the two-grid method.

General two-grid methods

Let V_h be fine grid finite element space and V_H be the coarse gird space (usually it is a subspace of V_h.) The two-grid method for equation (3.1) can be described as

alg:atwogrid **Algorithm 3.1** (General two-grid method)**.** Given an initial guess $u^{(0)} \in V_h$.

(i) **Pre-smoothing**: Apply a few relaxation steps to smooth $u^{(0)}$ on the fine grid to obtain a new approximation $u^{(1)} \in V_h$;

(ii) **Coarse-grid Correction**: Find $e_H \in V_H$ by solving or approximating the error equation

$$(\mathcal{A}e_H, v_H) = (f - \mathcal{A}u^{(1)}, v_H)$$

on the coarse gird, and then set $u^{(2)} = u^{(1)} + e_H$;

(iii) **Post-smoothing**: Apply relaxation to smooth $u^{(2)}$ on fine grid to obtain $u^{(3)} \in V_h$.

Remark 3.13 (Low frequency error)**.** A simple observation is that smooth functions can be represented on coarse grids rather accurately. For example, suppose u_h and u_H are the finite element solutions on V_h and $V_H \subset V_h$, respectively. Then we immediately have

$$a[u_h - u_H, v_H] = 0, \quad \forall\, v_H \in V_H.$$

Using the Aubin-Nitsche's argument, we consider a boundary value problem

$$-\Delta v = u_h - u_H \ \text{ in } \Omega \qquad \text{and} \qquad v = 0 \ \text{ on } \partial\Omega.$$

If we have full elliptic regularity, then $\|v\|_2 \leqslant C\|u_h - u_H\|_0$ is bounded. For any $v_H \in V_H$, we get

$$\|u_h - u_H\|_0^2 = a[v, u_h - u_H] = a[v - v_H, u_h - u_H] \leqslant |\!|\!| v - v_H |\!|\!| \, |\!|\!| u_h - u_H |\!|\!| \lesssim H|v|_2 \, |\!|\!| u_h - u_H |\!|\!| \,.$$

Hence the following inequality holds

$$\|u_h - u_H\|_0 \lesssim H \, |\!|\!| u_h - u_H |\!|\!| \lesssim H \, |\!|\!| u_h |\!|\!| \,. \tag{3.26}$$ eqn:SmoothPar

That is to say, if u_h is relatively smooth (small first derivatives), then u_h can be well approximated by u_H.

A more concrete algorithm based on the abstract algorithm above can be introduced. Let V be the fine space associated with meshsize h and $V_c \subset V$ be the coarse space associated with meshsize H. Let $\mathcal{I}_c : V_c \mapsto V$ be the natural embedding (injection), i.e., $\mathcal{I}_c v_c = v_c, \ \forall v_c \in V_c$.

Remark 3.14 (Embedding and projection)**.** By the definition of embedding $\mathcal{I}_c : V_c \mapsto V$ and the fact

$$(\mathcal{I}_c^T v, w_c) = (v, \mathcal{I}_c w_c) = (v, w_c), \quad \forall v \in V, w_c \in V_c,$$

it is easy to see that $\mathcal{I}_c^T = \mathcal{Q}_c$ is the $(\cdot,\cdot)$-projection from V to V_c. And the coarse-level operator can be defined by the Galerkin relation

$$\mathcal{A}_c = \mathcal{I}_c^T \mathcal{A} \mathcal{I}_c = \mathcal{Q}_c \mathcal{A} \mathcal{I}_c.$$

Suppose that $\mathcal{S}$ is a smoother and $\mathcal{B}_c$ is a solver (iterator) for the coarse-grid problem.

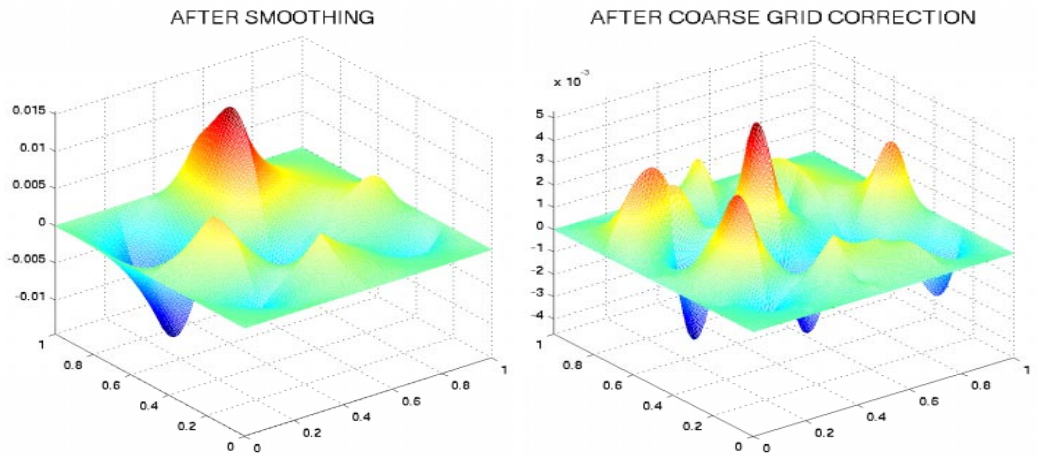
alg:twogrid **Algorithm 3.2** (Two-grid method)**.** Given an initial guess $u^{(0)} \in V$.

(i) **Pre-smoothing**: $u^{(1)} = u^{(0)} + \mathcal{S}(f - \mathcal{A}u^{(0)})$;

(ii) **Coarse-grid Correction**: $u^{(2)} = u^{(1)} + (\mathcal{I}_c \mathcal{B}_c \mathcal{I}_c^T)(f - \mathcal{A}u^{(1)})$;

(iii) **Post-smoothing**: $u^{(3)} = u^{(2)} + \mathcal{S}^T(f - \mathcal{A}u^{(2)})$.

We note that this algorithm is very similar to the multigrid algorithm discussed in Chapter 1. We can choose $\mathcal{S}$, V_c, and $\mathcal{B}_c$ to make the method efficient for the equation. The two-grid method is defined in the hope of capturing the high-frequency components of error on the fine grid, and leaving the low-frequency components to the coarser grid. The effect of coarse grid correction is illustrated in Figure 3.3. Note that these two pictures have different scales.



rsecorrection Figure 3.3: After coarse-grid correction, global low frequency is replaced by local high frequency.

Convergence analysis of TG

In this section, we will estimate convergence rate of two-grid methods. We first give two simple lemmas. Proof to these two simple lemmas are straightforward and left to the readers; see HW 3.4.

lem:TGiter **Lemma 3.5** (Iterator of two-grid method)**.** *The two-grid method has a corresponding iterator* $\mathcal{B}_{TG} : V' \to V$ *defined as*

$$\mathcal{B}_{TG} = \overline{\mathcal{S}} + \left(\mathcal{I} - \mathcal{S}^T\mathcal{A}\right)\mathcal{I}_c\mathcal{B}_c\mathcal{I}_c^T\left(\mathcal{I} - \mathcal{A}\mathcal{S}\right), \tag{3.27}$$ B_Inv

where $\overline{\mathcal{S}} = \mathcal{S}^T + \mathcal{S} - \mathcal{S}^T\mathcal{A}\mathcal{S}$ *is the symmetrization of the smoother* $\mathcal{S}$*.*

lem:TGerror **Lemma 3.6** (Error propagation of two-grid method)**.** *The error propagation operator* $\mathcal{E}_{TG} = \mathcal{I} - \mathcal{B}_{TG}\mathcal{A}$ *for two-grid method is*

$$\mathcal{E}_{TG} = (\mathcal{I} - \mathcal{S}^T\mathcal{A})(\mathcal{I} - \mathcal{B}_c\mathcal{A}_c\Pi_c)(\mathcal{I} - \mathcal{S}\mathcal{A}), \tag{3.28}$$ E_op

where Π_c *is the* $(\cdot,\cdot)_{\mathcal{A}}$*-orthogonal projection onto* V_c*. If the coarse-level solver is exact, namely,* $\mathcal{B}_c = \mathcal{A}_c^{-1}$*, then we have*

$$\mathcal{E}_{TG} = (\mathcal{I} - \mathcal{S}^T\mathcal{A})(\mathcal{I} - \Pi_c)(\mathcal{I} - \mathcal{S}\mathcal{A}). \tag{3.29}$$ E_op1

Notice that Π_c is the $\mathcal{A}$-projection from V to V_c. So there is an implicit natural embedding operator $\mathcal{I}_c$ in front of Π_c in the above equality.

We now present a theorem which gives the convergence rate of a simplified two-grid method (Algorithm 3.3) in terms of approximability of the coarser space V_c.

alg:twogrid2 **Algorithm 3.3** (Simplified two-grid method)**.** Given an initial guess $u^{(0)} \in V$.

(i) **Coarse-grid Correction**: $u^{(1)} = u^{(0)} + (\mathcal{I}_c \mathcal{B}_c \mathcal{I}_c^T)(f - \mathcal{A}u^{(0)})$;

(ii) **Post-smoothing**: $u^{(2)} = u^{(1)} + \mathcal{S}(f - \mathcal{A}u^{(1)})$.

Assume that $\overline{\mathcal{S}}$ is SPD. In the two-grid method analysis below, we need the following notation

$$\mathcal{T} = \mathcal{T}_{\overline{\mathcal{S}}} := \overline{\mathcal{S}}\mathcal{A} \,:\, V \mapsto V. \tag{3.30}$$ eqn:SA

With the above notation, the inner product

$$\big(u, v\big)_{\overline{\mathcal{S}}^{-1}} := \big(\mathcal{T}^{-1}u, v\big)_{\mathcal{A}},$$

the accompanying norm $\|\cdot\|_{\overline{\mathcal{S}}^{-1}}$, and $(\cdot,\cdot)_{\overline{\mathcal{S}}^{-1}}$-orthogonal projection $\mathcal{Q}_{\overline{\mathcal{S}}^{-1}} : V \mapsto V_c$. The convergence rate of the two-grid method is obtained in the following theorem; compare this result with the convergence rate of stationary iterative method in Theorem 2.3.

d-convergence **Theorem 3.3** (Convergence rate of the two-grid method)**.** *The convergence rate of the two-grid method* (3.27) *with the exact coarse-level solver is given by*

$$\|\mathcal{E}_{TG}\|_{\mathcal{A}}^2 = 1 - \frac{1}{c_1(V_c)}, \tag{3.31}$$ eq:two-grid-c

where

$$c_1(V_c) := \sup_{v\in V} \frac{\|(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})v\|_{\overline{\mathcal{S}}^{-1}}^2}{\|v\|_{\mathcal{A}}^2} = \sup_{v\in V} \inf_{v_c \in V_c} \frac{\|v - v_c\|_{\overline{\mathcal{S}}^{-1}}^2}{\|v\|_{\mathcal{A}}^2}. \tag{3.32}$$ KVc

Sketch of the proof. (1) It follows from (3.28) that

$$\mathcal{E}_{\mathrm{TG}} = (\mathcal{I} - \mathcal{S}\mathcal{A})(\mathcal{I} - \Pi_c).$$

Hence, we can immediately obtain

$$\|\mathcal{E}_{\mathrm{TG}}\|_{\mathcal{A}}^2 = \sup_{v\in V} \frac{\|(\mathcal{I} - \mathcal{S}\mathcal{A})(\mathcal{I} - \Pi_c)v\|_{\mathcal{A}}^2}{\|v\|_{\mathcal{A}}^2} = \sup_{v\in V_c^{\perp_{\mathcal{A}}}} \frac{\|(\mathcal{I} - \mathcal{S}\mathcal{A})v\|_{\mathcal{A}}^2}{\|v\|_{\mathcal{A}}^2}.$$

Using the definition of $(\cdot,\cdot)_{\mathcal{A}}$-projection Π_c, we can show that

$$\|\mathcal{E}_{\mathrm{TG}}\|_{\mathcal{A}}^2 = \sup_{v\in V_c^{\perp_{\mathcal{A}}}} \frac{\big((\mathcal{I} - \mathcal{T})v, v\big)_{\mathcal{A}}}{\|v\|_{\mathcal{A}}^2} = 1 - \inf_{v\in V_c^{\perp_{\mathcal{A}}}} \frac{\big(\mathcal{T}v, v\big)_{\mathcal{A}}}{(v,v)_{\mathcal{A}}} = 1 - \inf_{v\in V_c^{\perp_{\mathcal{A}}}} \frac{\big((\mathcal{I} - \Pi_c)\mathcal{T}v, v\big)_{\mathcal{A}}}{(v,v)_{\mathcal{A}}}.$$

(2) Define a new operator

$$\mathcal{X} := (\mathcal{I} - \Pi_c)\mathcal{T} \,:\, V_c^{\perp_{\mathcal{A}}} \mapsto V_c^{\perp_{\mathcal{A}}} \tag{3.33}$$ eqn:Xdef

and it is easy to check that $\mathcal{X}$ is self-adjoint with respect to $(\cdot,\cdot)_{\mathcal{A}}$. A key observation is that the inverse of $\mathcal{X}$ can be explicitly written as

$$\mathcal{Z} = \mathcal{T}^{-1}(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}}).$$

Since $\big(\Pi_c\mathcal{T}^{-1}(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})u, v\big)_{\mathcal{A}} = \big(\mathcal{T}^{-1}(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})u, v\big)_{\mathcal{A}} = \big((\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})u, v\big)_{\overline{\mathcal{S}}^{-1}} = 0$ for any $u \in V_c^{\perp_{\mathcal{A}}}$ and $v \in V_c$, we have $\Pi_c\mathcal{Z} = 0$, which implies that $\mathcal{Z} : V_c^{\perp_{\mathcal{A}}} \mapsto V_c^{\perp_{\mathcal{A}}}$. Furthermore, by the definition of projections, we get

$$\mathcal{X}\mathcal{Z} = (\mathcal{I} - \Pi_c)(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}}) = \mathcal{I} - \Pi_c = \mathcal{I} \quad \text{ on } V_c^{\perp_{\mathcal{A}}}.$$

(3) Consequently $\lambda_{\min}(\mathcal{X}) = \lambda_{\max}(\mathcal{Z})^{-1}$. Finally,

$$\begin{aligned}\lambda_{\max}(\mathcal{Z}) &= \sup_{v\in V_c^{\perp_{\mathcal{A}}}} \frac{(\mathcal{T}^{-1}(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})v, v)_{\mathcal{A}}}{(v,v)_{\mathcal{A}}} = \sup_{v\in V_c^{\perp_{\mathcal{A}}}} \frac{\big((\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})v, v\big)_{\overline{\mathcal{S}}^{-1}}}{(v,v)_{\mathcal{A}}} \\ &= \sup_{v\in V_c^{\perp_{\mathcal{A}}}} \frac{\|(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})v\|^2_{\overline{\mathcal{S}}^{-1}}}{(v,v)_{\mathcal{A}}} = \sup_{v\in V} \frac{\|(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})v\|^2_{\overline{\mathcal{S}}^{-1}}}{\|v\|^2_{\mathcal{A}}} =: c_1(V_c).\end{aligned}$$

The last identity holds because $\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}} = (\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})(\mathcal{I} - \Pi_c)$ and we can then take the supremum back over all $v \in V$ (similar to the argument in the very beginning of this proof). □

Optimal coarse space

Now we discuss how to choose the coarse space to maximize the convergence speed, which will become handy later for developing algebraic multigrid methods (AMGs). We will show that the space spanned by the eigenvectors of $\overline{\mathcal{S}}\mathcal{A}$ corresponding to small eigenvalues gives the "best" coarse space. Here "best" refers to this coarse space minimizes the convergence rate.

Theorem 3.3 provides an estimate on the convergence rate of a two-grid method in terms of $c_1(V_c)$. For a given method, a smaller bound on $c_1(V_c)$ means faster convergence. In particular, the two-grid method is uniformly convergent if $c_1(V_c)$ is uniformly bounded with respect to meshsize. However, one problem for applying Theorem 3.3 is that it is sometimes difficult to work with $\overline{\mathcal{S}}^{-1}$.

A natural approach to overcome such a difficulty is to introduce a simpler but spectrally equivalent SPD operator $\mathcal{D}$, such that

$$C_L\|v\|^2_{\mathcal{D}} \leqslant \|v\|^2_{\overline{\mathcal{S}}^{-1}} \leqslant C_U\|v\|^2_{\mathcal{D}}, \quad \forall\, v \in V.$$

Similar to the definition of $c_1(V_c)$, we can introduce the quantity

$$c_1(V_c, \mathcal{D}) = \sup_{v\in V} \frac{\|(\mathcal{I} - \mathcal{Q}_{\mathcal{D}})v\|^2_{\mathcal{D}}}{\|v\|^2_{\mathcal{A}}} = \sup_{v\in V} \inf_{v_c\in V_c} \frac{\|v - v_c\|^2_{\mathcal{D}}}{\|v\|^2_{\mathcal{A}}},$$

where $\mathcal{Q}_{\mathcal{D}} : V \mapsto V_c$ is the $(\cdot,\cdot)_{\mathcal{D}}$-orthogonal projection. Hence

$$C_L\, c_1(V_c, \mathcal{D}) \leqslant c_1(V_c) \leqslant C_U\, c_1(V_c, \mathcal{D}).$$

It is straight-forward to derive the following estimates:

-convergence2 **Theorem 3.4** (An estimate of convergence rate of TG)**.** *The convergence rate of the two-grid method* (3.27) *with exact coarse-level solver is given by*

$$1 - \frac{1}{C_L\, c_1(V_c, \mathcal{D})} \leqslant \|\mathcal{E}_{TG}\|_{\mathcal{A}} \leqslant 1 - \frac{1}{C_U\, c_1(V_c, \mathcal{D})} \leqslant 1 - \frac{1}{C_U C}, \tag{3.34}$$ eq:two-grid-c

where C is an upper bound of $c_1(V_c, \mathcal{D})$, i.e.,

$$\inf_{v_c \in V_c} \|v - v_c\|_{\mathcal{D}}^2 \leqslant C \|v\|_{\mathcal{A}}^2, \quad \forall\, v \in V. \tag{3.35}$$ KVc1

The following theorem characterizes the optimal choice of coarse space V_c with a fixed smoother $\mathcal{S}$:

-grid-optimal **Theorem 3.5** (Optimal coarse space)**.** *Given a smoother $\mathcal{S}$, the best coarse space of dimension N_c is given by*

$$V_c^{\text{opt}} := \underset{\dim V_c = N_c}{\operatorname{argmin}} \|\mathcal{E}_{TG}(V_c)\|_{\mathcal{A}} = \operatorname{span}\left\{\xi_k\right\}_{k=1}^{N_c}, \tag{3.36}$$ optVH

where $\left\{\xi_k\right\}_{k=1}^{N_c}$ are the eigenfunctions corresponding to the smallest eigenvalues λ_k of $\overline{\mathcal{S}}\mathcal{A}$.

Proof. Recall that $\mathcal{E}_{\text{TG}} = (\mathcal{I} - \mathcal{S}^T\mathcal{A})(\mathcal{I} - \Pi_c)(\mathcal{I} - \mathcal{S}\mathcal{A})$. Since $\mathcal{E}_{\text{TG}}$ depends on V_c we write $\mathcal{E}_{\text{TG}}(V_c)$ and using the same argument as in the proof of Theorem 3.3, we have

$$\|\mathcal{E}_{\text{TG}}(V_c)\|_{\mathcal{A}} \quad = \quad 1 - \min_{v \in V_c^{\perp_{\mathcal{A}}}} \frac{(\overline{\mathcal{S}}\mathcal{A}v, v)_{\mathcal{A}}}{\|v\|_{\mathcal{A}}^2}.$$

Thus,

$$\min_{\dim V_c = N_c} \|\mathcal{E}_{\text{TG}}(V_c)\|_{\mathcal{A}} = 1 - \max_{\dim V_c = N_c} \min_{v \in V_c^{\perp_{\mathcal{A}}}} \frac{(\overline{\mathcal{S}}\mathcal{A}v, v)_{\mathcal{A}}}{\|v\|_{\mathcal{A}}^2}.$$

By the well-known Courant minimax principle [26], we have

$$\max_{\dim V_c = N_c} \min_{v \in V_c^{\perp_{\mathcal{A}}}} \frac{(\overline{\mathcal{S}}\mathcal{A}v, v)_{\mathcal{A}}}{\|v\|_{\mathcal{A}}^2} = \lambda_{N_c+1}$$

and the equality holds if $V_c = V_c^{\text{opt}}$ as given in (3.36). □

Remark 3.15 (Lower bound of contraction factor)**.** Since the coarse space which minimizes the convergence rate is the coarse space which minimizes also $c_1(V_c)$, we have the following inequalities

$$c_1(V_c) = \frac{1}{1 - \|\mathcal{E}_{\text{TG}}\|_{\mathcal{A}}} \geqslant \frac{1}{\lambda_{N_c+1}} \qquad \text{or} \qquad \|\mathcal{E}_{\text{TG}}\|_{\mathcal{A}} \geqslant 1 - \lambda_{N_c+1},$$

which is a lower bound of the contraction factor in terms of size of the small eigenvalues (low frequencies) of $\overline{\mathcal{S}}\mathcal{A}$.

Since the eigenvalues of $\overline{\mathcal{S}}\mathcal{A}$ are expensive to compute, the practical value of Theorem 3.5 is limited. But it will provide useful guidance in the design practical algebraic multilevel methods in §6.3.

3.5 Matrix representation of the two-grid method

:MatrixFromTG

In practice, we have to understand the matrix representation of an abstract algorithm before we can actually implement it. We now explain the matrix representation of the two-grid method in the finite element context.

Grid transfer operators in matrix form

Let $\{\phi_i\}$ be the basis of a finite element space V on the fine-grid, then the stiffness matrix $\hat{\mathcal{A}}$ reads

$$\left(\hat{\mathcal{A}}\right)_{i,j} = a[\phi_i, \phi_j].$$

Let $\{\phi_l^c\}$ be the basis functions of the coarse-grid subspace $V_c \subset V$ and the stiffness matrix on the coarser space is denote by $\hat{\mathcal{A}}_c$ with $\left(\hat{\mathcal{A}}_c\right)_{k,l} = a[\phi_k^c, \phi_l^c]$. Then ϕ_l^c can be expressed as

$$\phi_l^c = \sum_{i=1}^{N} \left(P\right)_{i,l} \phi_i$$

or

$$(\phi_1^c, \ldots, \phi_{N_c}^c) = (\phi_1, \ldots, \phi_N) P,$$

which defines a *prolongation* matrix $P \in \mathbb{R}^{N \times N_c}$. By definition, this implies that $P = \underline{\mathcal{I}_c}$.

PreserveConst

Remark 3.16 (Cannonical prolongation operator)**.** Let $\mathbf{1}_N := (1, 1, \ldots, 1)^T$. Since the basis functions form the partition of unity, it follows that

$$(\phi_1, \cdots, \phi_N)\, \mathbf{1}_N = \sum_{i=1}^{N} \phi_i = 1 = \sum_{l=1}^{N_c} \phi_l^c = (\phi_1^c, \cdots, \phi_{N_c}^c)\mathbf{1}_{N_c} = (\phi_1, \cdots, \phi_N) P\, \mathbf{1}_{N_c}.$$

Hence we have that the prolongation matrix preserves constant away from the boundary, i.e.,

$$P\, \mathbf{1}_{N_c} = \mathbf{1}_N.$$

It is important to note that $\underline{\mathcal{I}_c^T} = \underline{\mathcal{Q}_c} \neq \underline{\mathcal{I}_c}^T$, i.e., the matrix representation of adjoint operator is not equal to the transpose of the matrix representation. If we take any $v \in V$, then we have

$$v_c := \mathcal{Q}_c v \qquad \text{and} \qquad v_c = (\phi_1^c, \ldots, \phi_{N_c}^c)\underline{v_c}.$$

On the other hand, with straightforward calculations, we obtain that

$$\vec{v}_c = \Big((v_c, \phi_k^c)\Big)_{k=1}^{N_c} = \Big((v, \phi_k^c)\Big)_{k=1}^{N_c} = \left(\sum_{j=1}^{N} \underline{v}_j (\phi_j, \phi_k^c)\right)_{k=1}^{N_c} = \left(\sum_{j=1}^{N} \underline{v}_j \left(\underline{\mathcal{I}_c}^T M\right)_{k,j}\right)_{k=1}^{N_c} = \underline{\mathcal{I}_c}^T M \underline{v}.$$

In turn, we can obtain the matrix representation of the L^2-projection

$$\underline{\mathcal{Q}_c} v = \underline{v_c} = M_c^{-1} \vec{v}_c = M_c^{-1} \underline{\mathcal{I}_c}^T M \underline{v} \quad \Longrightarrow \quad \underline{\mathcal{I}_c^T} = \underline{\mathcal{Q}_c} = M_c^{-1} \underline{\mathcal{I}_c}^T M = M_c^{-1} P^T M. \tag{3.37}$$ eqn:MatQ

Coarse problem in matrix form

Since the coarse-level operator is defined as $\mathcal{A}_c = \mathcal{I}_c^T \mathcal{A} \mathcal{I}_c$, we obtain its matrix representation

$$\underline{\mathcal{A}_c} = \underline{\mathcal{Q}_c}\, \underline{\mathcal{A}}\, \underline{\mathcal{I}_c} \quad \Longrightarrow \quad \hat{\mathcal{A}}_c = M_c \underline{\mathcal{A}_c} = M_c \underline{\mathcal{Q}_c}\, \underline{\mathcal{A}}\, \underline{\mathcal{I}_c} = P^T M \underline{\mathcal{A}} P = P^T \hat{\mathcal{A}} P. \tag{3.38}$$ eqn:Ac-PtAP

Then the coarse stiffness matrix satisfies

$$\hat{\mathcal{A}}_c = P^T \hat{\mathcal{A}} P. \tag{3.39}$$ eqn:PtAP

Therefore, the algebraic form (3.39) of the coarse level problem is equivalent to the matrix representation of the operator form.

In the above equality, we observe that, the L^2-projection $\mathcal{Q}_c$ is not needed for implementation. Instead, we only need to use a restriction matrix $R := P^T$.

Remark 3.17 (Finite difference case)**.** Notice that, here, for the finite element stiffness matrices, the restriction matrix is just $R = P^T$. However, we have already noticed that $R \neq P^T$ for the finite difference method in (1.34). In fact, many books (see [23] for example) states $R = cP^T$. This difference comes from the scaling effect caused by h. In the 1D FD example, the coefficient matrices on fine and coarse levels are $A = h^{-1}\hat{\mathcal{A}}$ and $A_c = H^{-1}\hat{\mathcal{A}}_c$, respectively. Hence we get

$$\hat{\mathcal{A}}_c = P^T \hat{\mathcal{A}} P \quad \Longrightarrow \quad A_c = \left(\frac{h}{H} P^T\right) A P =: RAP.$$

This remark explains how we can obtain such the constant c in general.

Two-grid iterator in matrix form

From (3.27), we have that the two-grid method with exact coarse solver is

$$\mathcal{B}_{\mathrm{TG}} = \overline{\mathcal{S}} + \left(\mathcal{I} - \mathcal{S}^T \mathcal{A}\right) \mathcal{I}_c \mathcal{A}_c^{-1} \mathcal{I}_c^T \left(\mathcal{I} - \mathcal{A}\mathcal{S}\right).$$

We can then write the above equation in matrix form

$$\underline{\mathcal{B}_{\mathrm{TG}}} = \overline{\mathcal{S}} + \left(\underline{\mathcal{I}} - \underline{\mathcal{S}^T \mathcal{A}}\right) \underline{\mathcal{I}_c \mathcal{A}_c^{-1} \mathcal{I}_c^T} \left(\underline{\mathcal{I}} - \underline{\mathcal{A}\mathcal{S}}\right).$$

So we define

$$B_{\mathrm{TG}} := \underline{\mathcal{B}_{\mathrm{TG}}} M^{-1} = \overline{\mathcal{S}} M^{-1} + (\underline{\mathcal{I}} - \underline{\mathcal{S}^T \mathcal{A}})\underline{\mathcal{I}_c \mathcal{A}_c^{-1} \mathcal{I}_c^T}(\underline{\mathcal{I}} - \underline{\mathcal{A}\mathcal{S}}) M^{-1}.$$

Using the matrix form the symmetrization, inversion, and transpose derived earlier, we can easily get

$$B_{\mathrm{TG}} = \overline{S} + (I - S^T A) P A_c^{-1} P^T (I - AS) = \overline{S} + (I - S^T A) P \big(P^T A P\big)^{-1} P^T (I - AS).$$

Now we are ready to introduce the matrix representation of the two-grid method for solving the linear system $A\underline{u} = \vec{f}$. We describe the two-grid method as a preconditioner action $B_{\mathrm{TG}}(\cdot)$. For any given vector (usually it is the residual vector) $\vec{r} \in \mathbb{R}^N$, we can compute $B_{\mathrm{TG}}(\vec{r})$ in the following steps:

Listing 3.1: A two-grid method

```
%% Given any vector r⃗;
Pre−smoothing: v⃗ ← Sr⃗;
Coarse−grid correction: w⃗ ← v⃗ + P(P^T AP)^{−1} P^T (r⃗ − Av⃗);
Post−smoothing: B_TG r⃗ ← w⃗ + S^T (r⃗ − Aw⃗);
```

3.6 Homework problems

hw:ApostBound **HW 3.1.** Show the a posteriori error bounds (3.8).

hw:matrix-rep **HW 3.2.** Prove the statements in Lemma 3.2.

hw:Richardson **HW 3.3.** Show the operator form and matrix form (3.14) of the Richardson method.

hw:TG **HW 3.4.** Prove Lemma 3.5 and Lemma 3.6.

hw:MGTG **HW 3.5.** Write the 1D multigrid method in §1.4 as a two-grid method (Algorithm 3.2) called recursively.

hw:ETG **HW 3.6.** Give the detailed proof of Theorem 3.3. Hint: First show that

$$\sup_{v \in V} \frac{\|(\mathcal{I} - \mathcal{S}\mathcal{A})(\mathcal{I} - \Pi_c)v\|_{\mathcal{A}}^2}{\|v\|_{\mathcal{A}}^2} = \sup_{v \in V} \frac{\|(\mathcal{I} - \mathcal{S}\mathcal{A})(\mathcal{I} - \Pi_c)v\|_{\mathcal{A}}^2}{\|(\mathcal{I} - \Pi_c)v\|_{\mathcal{A}}^2 + \|\Pi_c v\|_{\mathcal{A}}^2} = \sup_{v \in V_c^{\perp_{\mathcal{A}}}} \frac{\|(\mathcal{I} - \mathcal{S}\mathcal{A})v\|_{\mathcal{A}}^2}{\|v\|_{\mathcal{A}}^2};$$

Then prove that $\mathcal{X}$ defined in (3.33) is self-adjoint with respect to $(\cdot,\cdot)_{\mathcal{A}}$-inner product.

Chapter 4

Subspace Correction Methods

ch:subspace

In the previous chapters, we have been considering the linear equation

$$\mathcal{A}u = f, \tag{4.1}$$ eqn:original

where $\mathcal{A} : V \mapsto V$ is SPD. A linear stationary iterative method can be written as

$$u^{\text{new}} = u^{\text{old}} + \mathcal{B}(f - \mathcal{A}u^{\text{old}}). \tag{4.2}$$ eqn:original-

If $\mathcal{B}$ is an SPD operator, with proper scaling, the above iterative method converges. Furthermore, $\mathcal{B}$ can be applied as a preconditioner of Krylov subspace methods.

In this chapter, we present a framework for analyzing linear iterative methods and/or preconditioners obtained by the concept of space decomposition and subspace corrections. This general framework can be used to establish convergence theory for various methods, including the multigrid method and domain decomposition method discussed in the previous chapters.

4.1 Successive and parallel subspace corrections

Suppose we have a subspace decomposition of the solution space

$$V = \sum_{j=1}^{J} V_j \qquad \text{and} \qquad V_j \subset V \ \ (j = 1, \ldots, J).$$

For any $v \in V$, we can write it as $v = \sum_{j=1}^{J} v_j$ with $v_j \in V_j$. Notice that this representation is not unique as there could be redundancy in the subspace decomposition. Furthermore, such redundancy is crucial for optimality of multilevel methods.

Abstract framework for subspace corrections

We first define a few operators which have already been used at different places in the previous chapters.

def:subproblems **Definition 4.1.** *Let V be a Hilbert space with inner product $(\cdot,\cdot)$ and $V_j \subset V$ be a subspace. We define*

$$\begin{cases} \text{subspace problem} & \mathcal{A}_j : V_j \mapsto V_j, & (\mathcal{A}_j v_j, w_j) = (\mathcal{A} v_j, w_j), & \forall\, v_j, w_j \in V_j; \\ (\cdot,\cdot)\text{-projection} & \mathcal{Q}_j : V \mapsto V_j, & (\mathcal{Q}_j v, w_j) = (v, w_j), & \forall\, w_j \in V_j; \\ (\cdot,\cdot)_{\mathcal{A}}\text{-projection} & \Pi_j : V \mapsto V_j, & (\Pi_j v, w_j)_{\mathcal{A}} = (v, w_j)_{\mathcal{A}}, & \forall\, w_j \in V_j. \end{cases}$$

Remark 4.1 (Matrix representation of the $\mathcal{A}$-projection)**.** Let $u_c := \Pi_c u$. Since $\Pi_c : V \mapsto V_c \subset V$ is the $\mathcal{A}$-orthogonal projection operator, for any $u \in V$, we have

$$a[u_c, v_c] = a[\Pi_c u, v_c] = a[u, v_c], \quad \forall\, v_c \in V_c.$$

Using the matrix representation notations introduced in §3.2, we have, for any $v_c \in V_c$, that

$$a[u_c, v_c] = (\mathcal{A} u_c, v_c) = \underline{v_c}^T \hat{\mathcal{A}}_c \, \underline{u_c}, \qquad \forall\, u_c \in V_c; \tag{4.3}$$

$$a[u, v_c] = (\mathcal{A} u, v_c) = (\underline{\mathcal{I}_c v_c})^T \hat{\mathcal{A}} \, \underline{u} = \underline{v_c}^T P^T \hat{\mathcal{A}} \, \underline{u}, \qquad \forall\, u \in V. \tag{4.4}$$

From (4.3) and (4.4), we can derive the matrix representation of the Galerkin projection on the coarse grid

$$\hat{\mathcal{A}}_c \underline{u_c} = P^T \hat{\mathcal{A}} \underline{u} \quad \Longrightarrow \quad \underline{\Pi_c} \, \underline{u} = \underline{\Pi_c u} = \underline{u_c} = \hat{\mathcal{A}}_c^{-1} P^T \hat{\mathcal{A}} \underline{u}.$$

Hence, we obtain the matrix representation of the $\mathcal{A}$-projection operator

$$\underline{\Pi_c} = \hat{\mathcal{A}}_c^{-1} P^T \hat{\mathcal{A}}. \tag{4.5}$$

Using Definition 4.1, we have the following elementary results (the proof is left to the readers; see HW 4.2):

lem:projections **Lemma 4.1** (Relation between projections)**.** *The following equalities hold:*

1. $\mathcal{I}_j^T = \mathcal{Q}_j, \;\; \mathcal{I}_j^* = \Pi_j$;
2. $\mathcal{Q}_j \mathcal{A} = \mathcal{A}_j \Pi_j$.

Remark 4.2 (Subspace problems)**.** From the definition of $\mathcal{A}_j$, we get

$$\mathcal{A}_j = \mathcal{I}_j^T \mathcal{A} \mathcal{I}_j = \mathcal{Q}_j \mathcal{A} \mathcal{I}_j = \mathcal{Q}_j \mathcal{A} \mathcal{Q}_j^T.$$

With the help of Lemma 4.1 and simple calculations, we obtain the error equation on each subspace V_j:

$$\mathcal{A} e = r \quad \Longrightarrow \quad Q_j \mathcal{A} e = \mathcal{Q}_j r \quad \Longrightarrow \quad \mathcal{A}_j \Pi_j e = \mathcal{Q}_j r \quad \Longrightarrow \quad \mathcal{A}_j e_j = r_j,$$

where $r_j = \mathcal{Q}_j r$ and $e_j = \Pi_j e$.

The idea of method of subspace corrections (MSC) is motivated by Algorithm 2.1[1] and the idea of *divide and conquer*. We can describe the idea of subspace correction in the following abstract algorithm[2], which is a just generalization of Algorithm 2.1:

alg:msc **Algorithm 4.1** (Method of subspace corrections)**.** $u^{\text{new}} = SC(u^{\text{old}})$

(i) **Form residual**: $r = f - \mathcal{A}u^{\text{old}}$

(ii) **Solve error equation on** V_j: $\mathcal{A}_j e_j = r_j$ by $e_j \approx \hat{e}_j = \mathcal{S}_j r_j$

(iii) **Correct iteration**: $u^{\text{new}} = u^{\text{old}} + \hat{e}_j$

Notice that, instead of constructing an iterator for the whole system, Algorithm 4.1 only considers subproblems on subspaces.

rem:subsolvers **Remark 4.3** (Subspace solvers)**.** It is well-known that

$$u_j = \operatorname*{argmin}_{v \in V_j} \mathcal{F}(v) := \frac{1}{2}(\mathcal{A}v, v) - (f, v)$$

is equivalent to

$$u_j = \operatorname*{argmin}_{v \in V_j} \|u - v\|_{\mathcal{A}}.$$

We notice that the solution of the subspace problem $\mathcal{A}_j e_j = r_j = \mathcal{Q}_j r^{\text{old}}$ satisfies that

$$\mathcal{F}(u^{\text{old}} + e_j) = \min_{e \in V_j} \mathcal{F}(u^{\text{old}} + e).$$

SSC and PSC methods

Algorithm 4.1 does not specify how to combine the corrections $\hat{e}_j$'s from different subspaces. There are basically two approaches: the successive subspace correction (SSC) and the parallel subspace correction (PSC). SSC can be viewed as the multiplicative Schwarz method (2.33) and PSC can be viewed as the additive Schwarz method (2.32). We now give description of the SSC and PSC algorithms.

alg:ssc **Algorithm 4.2** (Successive subspace corrections)**.** $u^{\text{new}} = SSC(u^{\text{old}})$

(i) $v = u^{\text{old}}$

(ii) $v = v + \mathcal{S}_j \mathcal{Q}_j (f - \mathcal{A}v), \quad j = 1, \ldots, J$

[1] The corresponding operator form is just (4.2).

[2] Note that this procedure is not really an algorithm as it does not specify how to combine the corrections $\hat{e}_j$'s from different subspaces.

(iii) $u^{\text{new}} = v$

alg:psc **Algorithm 4.3** (Parallel subspace corrections)**.** $u^{\text{new}} = PSC(u^{\text{old}})$

(i) $r = f - \mathcal{A}u^{\text{old}}$

(ii) $\hat{e}_j = \mathcal{S}_j \mathcal{Q}_j r,\ j = 1, \ldots, J$

(iii) $u^{\text{new}} = u^{\text{old}} + \sum_{j=1}^{J} \hat{e}_j$

From the above algorithms, it is immediately clear why they are named as PSC and SSC, respectively. As in (3.30), we define an operator

$$\mathcal{T}_j = \mathcal{T}_{\mathcal{S}_j} := \mathcal{S}_j \mathcal{Q}_j \mathcal{A} = \mathcal{S}_j \mathcal{A}_j \Pi_j : V \mapsto V_j.$$

Apparently, if we restrict the domain to V_j, then we have

$$\mathcal{T}_j = \mathcal{T}_{\mathcal{S}_j} = \mathcal{S}_j \mathcal{A}_j : V_j \mapsto V_j.$$

We shall now assume all the subspace solvers (smoothers) $\mathcal{S}_j$ are SPD operators. As $\mathcal{S}_j^T = \mathcal{S}_j$, the operator $\mathcal{T}_j = \mathcal{S}_j \mathcal{A}_j : V_j \mapsto V_j$ is symmetric and positive definite with respect to $(\cdot, \cdot)_{\mathcal{A}}$. If $\mathcal{S}_j = \mathcal{A}_j^{-1}$, i.e., the smoother is the exact solver on each subspace, then we have $\mathcal{T}_j = \Pi_j$.

- The SSC method satisfies:

$$u - u^{\text{new}} = (\mathcal{I} - \mathcal{B}\mathcal{A})(u - u^{\text{old}}) = (\mathcal{I} - \mathcal{T}_J) \cdots (\mathcal{I} - \mathcal{T}_1)(u - u^{\text{old}}). \tag{4.6}$$ eqn:SSC

  If $J = N$ and each subspace $V_j = \text{span}\{\phi_j\}$ $(j = 1, \ldots, N)$ and $\mathcal{S}_j = \mathcal{A}_j^{-1}$, then the corresponding SSC method is exactly the G-S method; see (2.17).

- For the PSC method, the iterator (or, more often, the preconditioner) satisfies

$$\mathcal{B} = \sum_{j=1}^{J} \mathcal{S}_j \mathcal{Q}_j = \sum_{j=1}^{J} \mathcal{I}_j \mathcal{S}_j \mathcal{Q}_j \qquad \text{and} \qquad \mathcal{B}\mathcal{A} = \sum_{j=1}^{J} \mathcal{S}_j \mathcal{Q}_j \mathcal{A} = \sum_{j=1}^{J} \mathcal{T}_j. \tag{4.7}$$ eqn:PSC

  If $\mathcal{S}_j$'s $(j = 1, \ldots, J)$ are all SPD, then the preconditioner $\mathcal{B}$ is also SPD; see HW 4.3. If each subspace $V_j = \text{span}\{\phi_j\}$ $(j = 1, \ldots, N)$, then the PSC methods with $\mathcal{S}_j = \omega(\cdot, \phi_j)\phi_j$ and $\mathcal{S}_j = \mathcal{A}_j^{-1}$ correspond to the Richardson method and the Jacobi method, respectively.

4.2 Expanded system and block solvers

Back in §2.1, we have discussed a modified block Gauss–Seidel method. In this section, we discuss an expanded system of (4.1) and its block iterative solvers. Moreover, we will show how these block solvers are related to the subspace correction methods for the original linear system (4.1). This relation will become important in the next section for deriving the X-Z identity, which gives the convergence rate of SSC.

Expansion of the original problem

Suppose that the finite dimensional vector space V can be decomposed as the summation of linear vector subspaces (might not be linearly independent), $V_1, V_2, \ldots, V_J$, i.e., $V = \sum_{j=1}^{J} V_j$. We define a new vector space

$$\mathbf{V} := V_1 \times V_2 \times \cdots \times V_J.$$

Define an operator $\mathbf{\Pi} : \mathbf{V} \mapsto V$ such that $\mathbf{\Pi u} = \sum_{j=1}^{J} u_j$, where $\mathbf{u} = (u_1, \ldots, u_J)^T \in \mathbf{V}$ with each component $\mathbf{u}_j = u_j \in V_j$. From the definition, $\mathbf{\Pi}$ is surjective. This operator can be formally interpreted as

$$\mathbf{\Pi} = (\mathcal{I}_1, \ldots, \mathcal{I}_J),$$

where $\mathcal{I}_j : V_j \mapsto V$ is the natural embedding. Hence, we obtain

$$\mathbf{\Pi u} = (\mathcal{I}_1, \ldots, \mathcal{I}_J) \begin{pmatrix} u_1 \\ \vdots \\ u_J \end{pmatrix} = \sum_{j=1}^{J} \mathcal{I}_j u_j = \sum_{j=1}^{J} u_j.$$

So we have

$$\mathbf{\Pi}^T = \begin{pmatrix} \mathcal{I}_1^T \\ \vdots \\ \mathcal{I}_J^T \end{pmatrix} = \begin{pmatrix} \mathcal{Q}_1 \\ \vdots \\ \mathcal{Q}_J \end{pmatrix}.$$

Define $\mathbf{A} : \mathbf{V} \mapsto \mathbf{V}$ such that $\mathbf{A}_{i,j} = \mathcal{A}_{i,j} := \mathcal{I}_i^T \mathcal{A} \mathcal{I}_j : V_j \mapsto V_i$. And we denote $\mathcal{A}_j := \mathcal{A}_{j,j}$. Hence we can write the operator $\mathbf{A}$ in a matrix form

$$\mathbf{A} := \mathbf{\Pi}^T \mathcal{A} \mathbf{\Pi} = \left(\mathbf{A}_{i,j}\right)_{J \times J} = \begin{pmatrix} \mathcal{A}_{1,1} & \cdots & \mathcal{A}_{1,J} \\ \vdots & \ddots & \vdots \\ \mathcal{A}_{J,1} & \cdots & \mathcal{A}_{J,J} \end{pmatrix}.$$

Given any right hand side function $f \in V$, we define

$$\mathbf{f} := \mathbf{\Pi}^T f = \begin{pmatrix} \mathcal{I}_1^T f \\ \vdots \\ \mathcal{I}_J^T f \end{pmatrix} \in \mathbf{V}.$$

In this setting, we can consider the following problem: Find $\mathbf{u} \in \mathbf{V}$, such that

$$\mathbf{A u} = \mathbf{f}. \tag{4.8}$$

eqn:expanded

This system is called the *expanded equation* of the original linear equation (4.1). We will see how the solution of these two problems are related.

If $\mathcal{A}$ is SPD, then $\mathbf{A}$ is a symmetric positive semidefinite (SPSD) operator. Note that $\mathbf{A}$ could be singular due to nontrivial null space, null($\mathbf{\Pi}$). However, its diagonal entries $\mathcal{A}_j$ ($j = 1, 2, \ldots, J$) are non-singular. We can define a semi-norm for $\mathbf{B} : \mathbf{V} \mapsto \mathbf{V}$

$$\|\mathbf{B}\|_{\mathbf{A}} := \sup_{\|\mathbf{v}\|_{\mathbf{A}} \neq 0} \frac{\|\mathbf{B}\mathbf{v}\|_{\mathbf{A}}}{\|\mathbf{v}\|_{\mathbf{A}}}.$$

Block solvers for expanded equation

As before, we denote the lower, upper, and diagonal part of $\mathbf{A}$ as $\mathbf{L}$, $\mathbf{U}$, and $\mathbf{D}$, respectively. We can immediately see that the stationary iterative methods discussed in §1.3 can be easily adapted to solve (4.8). The linear stationary iterative methods for (4.8) can be written in the following abstract form

$$\mathbf{u}^{\text{new}} = \mathbf{u}^{\text{old}} + \mathbf{B}(\mathbf{f} - \mathbf{A}\mathbf{u}^{\text{old}}), \tag{4.9}$$ eqn:iterEx

where the iterator $\mathbf{B} : \mathbf{V} \mapsto \mathbf{V}$. If $\mathbf{B} = \mathbf{D}^{-1}$, then we have the block Jacobi method for (4.8); if $\mathbf{B} = (\mathbf{D} + \mathbf{L})^{-1}$, then we have the block Gauss–Seidel method.

Motivated by (2.13), we can generalize the block Jacobi and G-S methods a little bit. Assume there is a non-singular block diagonal smoother (or relaxation operator) $\mathbf{S} : \mathbf{V} \mapsto \mathbf{V}$, i.e.,

$$\mathbf{S} = \text{diag}(\mathcal{S}_1, \mathcal{S}_2, \ldots, \mathcal{S}_J), \quad \text{with } \mathcal{S}_j : V_j \mapsto V_j, \quad j = 1, 2, \ldots, J.$$

We define modified block Jacobi method by $\mathbf{B} = \mathbf{S}$ and the modified block Gauss–Seidel method by $\mathbf{B} = (\mathbf{S}^{-1} + \mathbf{L})^{-1}$.

m:equivalence **Theorem 4.1** (Solution of expanded and original systems)**.** *The linear stationary iteration* (4.9) *for the equation* (4.8) *reduces to an equivalent stationary iteration* (4.2) *with the iterator* $\mathcal{B} = \mathbf{\Pi}\,\mathbf{B}\,\mathbf{\Pi}^T$ *for the original equation* (4.1)*. Moreover, these two methods have the same convergence rate, i.e.,*

$$\|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}} = \|\mathbf{I} - \mathbf{B}\mathbf{A}\|_{\mathbf{A}}.$$

Proof. The linear stationary iterative method

$$\mathbf{u}^{\text{new}} = \mathbf{u}^{\text{old}} + \mathbf{B}(\mathbf{f} - \mathbf{A}\mathbf{u}^{\text{old}})$$

is equivalent to

$$\begin{aligned} \mathbf{u}_j^{\text{new}} &= \mathbf{u}_j^{\text{old}} + \sum_k \mathbf{B}_{j,k}\Big(\mathcal{I}_k^T f - \sum_i \mathbf{A}_{k,i}\mathbf{u}_i^{\text{old}}\Big) \\ &= \mathbf{u}_j^{\text{old}} + \sum_k \mathbf{B}_{j,k}\mathcal{I}_k^T\Big(f - \sum_i \mathcal{A}\mathcal{I}_i\mathbf{u}_i^{\text{old}}\Big) \;=\; \mathbf{u}_j^{\text{old}} + \sum_k \mathbf{B}_{j,k}\mathcal{I}_k^T\Big(f - \mathcal{A}u^{\text{old}}\Big). \end{aligned}$$

Therefore, we have

$$u^{\text{new}} = \sum_j \mathcal{I}_j \mathbf{u}_j^{\text{new}} = u^{\text{old}} + \sum_{j,k} \mathcal{I}_j \mathbf{B}_{j,k} \mathcal{I}_k^T \Big(f - \mathcal{A} u^{\text{old}} \Big) = u^{\text{old}} + \mathcal{B} \Big(f - \mathcal{A} u^{\text{old}} \Big).$$

This proves the equivalence of (4.9) and (4.2).

A key observation is that

$$(\mathbf{BAv}, \mathbf{v})_{\mathbf{A}} = (\mathbf{ABAv}, \mathbf{v}) = (\mathbf{\Pi}^T \mathcal{A} \mathbf{\Pi}\, \mathbf{B}\, \mathbf{\Pi}^T \mathcal{A} \mathbf{\Pi v}, \mathbf{v}) = (\mathcal{AB}\mathcal{A} \mathbf{\Pi v}, \mathbf{\Pi v}) = (\mathcal{B}\mathcal{A} \mathbf{\Pi v}, \mathbf{\Pi v})_{\mathcal{A}}.$$

The contraction factor can be written

$$\begin{aligned}
\|\mathcal{I} - \mathcal{BA}\|_{\mathcal{A}}^2 &= \sup_{v \neq 0} \frac{\|(\mathcal{I} - \mathcal{BA})v\|_{\mathcal{A}}^2}{\|v\|_{\mathcal{A}}^2} = \sup_{v \neq 0} \frac{(v,v)_{\mathcal{A}} - \big((\mathcal{B}^T + \mathcal{B} - \mathcal{B}^T \mathcal{AB})\mathcal{A}v, v\big)_{\mathcal{A}}}{(v,v)_{\mathcal{A}}} \\
&= \sup_{\mathbf{\Pi v} \neq 0} \frac{(\mathbf{\Pi v}, \mathbf{\Pi v})_{\mathcal{A}} - \big((\mathcal{B}^T + \mathcal{B} - \mathcal{B}^T \mathcal{AB})\mathcal{A}\mathbf{\Pi v}, \mathbf{\Pi v}\big)_{\mathcal{A}}}{(\mathbf{\Pi v}, \mathbf{\Pi v})_{\mathcal{A}}} \\
&= \sup_{\|\mathbf{v}\|_{\mathbf{A}} \neq 0} \frac{(\mathbf{v}, \mathbf{v})_{\mathbf{A}} - \big((\mathbf{B}^T + \mathbf{B} - \mathbf{B}^T \mathbf{AB})\mathbf{Av}, \mathbf{v}\big)_{\mathbf{A}}}{\|\mathbf{v}\|_{\mathbf{A}}^2} \\
&= \|\mathbf{I} - \mathbf{BA}\|_{\mathbf{A}}^2.
\end{aligned}$$

Hence we get the desired result. □

Example 4.1 (Block Jacobi method and PSC)**.** We now apply the block Jacobi method for the expanded system (4.8), i.e.,

$$\mathbf{u}^{\text{new}} = \mathbf{u}^{\text{old}} + \mathbf{D}^{-1}(\mathbf{f} - \mathbf{A}\mathbf{u}^{\text{old}}).$$

We notice that $\mathbf{D}^{-1}\mathbf{A} = \mathbf{D}^{-1}\mathbf{\Pi}^T \mathcal{A} \mathbf{\Pi}$, which is spectrally equivalent[3] to $\mathbf{\Pi} \mathbf{D}^{-1} \mathbf{\Pi}^T \mathcal{A}$. In fact, from Theorem 4.1, we can see that the above iterative method is equivalent to

$$u^{\text{new}} = u^{\text{old}} + \mathbf{\Pi} \mathbf{D}^{-1} \mathbf{\Pi}^T (f - \mathcal{A} u^{\text{old}}) = u^{\text{old}} + \sum_{j=1}^{J} \mathcal{I}_j \mathcal{A}_j^{-1} \mathcal{I}_j^T \left(f - \mathcal{A} u^{\text{old}} \right).$$

We immediately recognize that this is the PSC method or the additive Schwarz method with exact subspace solvers. □

Example 4.2 (Block G-S method and SSC)**.** Similar to the above example, we can get the block G-S method is just the SSC method or the multiplicative Schwarz method for the original problem. We now apply the block G-S method for the expanded system (4.8), i.e.,

$$\mathbf{u}^{\text{new}} = \mathbf{u}^{\text{old}} + (\mathbf{D} + \mathbf{L})^{-1}(\mathbf{f} - \mathbf{A}\mathbf{u}^{\text{old}}).$$

We can rewrite this method as

$$(\mathbf{D} + \mathbf{L})\mathbf{u}^{\text{new}} = (\mathbf{D} + \mathbf{L})\mathbf{u}^{\text{old}} + (\mathbf{f} - \mathbf{A}\mathbf{u}^{\text{old}}).$$

[3] Note that $\sigma(\mathcal{BA})\backslash\{0\} = \sigma(\mathcal{AB})\backslash\{0\}$.

Hence we have

$$\mathbf{D}\mathbf{u}^{\text{new}} = \mathbf{D}\mathbf{u}^{\text{old}} + \mathbf{f} - \mathbf{L}\mathbf{u}^{\text{new}} - (\mathbf{D}+\mathbf{U})\mathbf{u}^{\text{old}};$$

in turn, we get

$$\mathbf{u}^{\text{new}} = \mathbf{u}^{\text{old}} + \mathbf{D}^{-1}\Big(\mathbf{f} - \mathbf{L}\mathbf{u}^{\text{new}} - (\mathbf{D}+\mathbf{U})\mathbf{u}^{\text{old}}\Big).$$

For $j = 1, \ldots, J$, the block G-S method can be written as

$$u_j^{\text{new}} = u_j^{\text{old}} + \mathcal{A}_j^{-1}\Big(\mathcal{I}_j^T f - \sum_{i<j}\mathcal{I}_j^T\mathcal{A}\mathcal{I}_i u_i^{\text{new}} - \sum_{i\geqslant j}\mathcal{I}_j^T\mathcal{A}\mathcal{I}_i u_i^{\text{old}}\Big).$$

We define iteration

$$u^{\frac{j}{J}} := \sum_{i<j} u_i^{\text{new}} + \sum_{i\geqslant j} u_i^{\text{old}} = \sum_{i<j}\mathcal{I}_i u_i^{\text{new}} + \sum_{i\geqslant j}\mathcal{I}_i u_i^{\text{old}}, \quad j = 1, \ldots, J.$$

By this definition, we can see that

$$u^{\frac{j+1}{J}} = u^{\frac{j}{J}} + \mathcal{I}_j u_j^{\text{new}} - \mathcal{I}_j u_j^{\text{old}} = u^{\frac{j}{J}} + \mathcal{I}_j\mathcal{A}_j^{-1}\mathcal{I}_j^T(f - \mathcal{A}u^{\frac{j}{J}}).$$

Here the term $f - \mathcal{A}u^{\frac{j}{J}}$ is sometimes called the *dynamic residual*, which is the residual at an inner iteration of the G-S method. From the above equation, we notice that the block G-S method is just the SSC method with exact subspace solvers $\mathcal{S}_j = \mathcal{A}_j^{-1}$ for the original linear equation (4.1). □

Convergence of block solvers

Motived by the weighted Jacobi and G-S methods, we assume that there is an invertible smoother or local relaxation $\mathbf{S}$ for solving $\mathbf{A}\mathbf{u} = \mathbf{f}$. Similar to the method presented in §2.1, we define a general or modified block G-S method:

$$\mathbf{B} := \big(\mathbf{S}^{-1} + \mathbf{L}\big)^{-1}. \tag{4.10}$$ eqn:blockMGS

We analyze the convergence rate of this method. Let $\mathbf{K} := \mathbf{B}^{-T} + \mathbf{B}^{-1} - \mathbf{A}$ be a symmetric operator and the symmetrization operator as $\overline{\mathbf{B}} = \mathbf{B}^T\mathbf{K}\mathbf{B}$. Then we get

$$\Big(\overline{\mathbf{B}}^{-1}\mathbf{v}, \mathbf{v}\Big) = \Big(\mathbf{B}^{-1}\mathbf{K}^{-1}\mathbf{B}^{-T}\mathbf{v}, \mathbf{v}\Big) = \Big((\mathbf{S}^{-1}+\mathbf{L})\mathbf{K}^{-1}(\mathbf{S}^{-T}+\mathbf{U})\mathbf{v},\, \mathbf{v}\Big), \quad \forall \mathbf{v} \in \mathbf{V} \tag{4.11}$$ eqn:invBbar1

By the definition of $\mathbf{K}$, it is clear that $\mathbf{K}$ is diagonal and

$$\mathbf{K} = (\mathbf{S}^{-T}+\mathbf{U}) + (\mathbf{S}^{-1}+\mathbf{L}) - (\mathbf{D}+\mathbf{L}+\mathbf{U}) = \mathbf{S}^{-T} + \mathbf{S}^{-1} - \mathbf{D} = \mathbf{S}^{-T}\overline{\mathbf{S}}\mathbf{S}^{-1},$$

where $\overline{\mathbf{S}} := \mathbf{S}^T + \mathbf{S} - \mathbf{S}^T\mathbf{D}\mathbf{S}$. Hence, we get

$$\mathbf{K}^{-1} = \big(\mathbf{S}^{-T} + \mathbf{S}^{-1} - \mathbf{D}\big)^{-1} = \mathbf{S}\overline{\mathbf{S}}^{-1}\mathbf{S}^T. \tag{4.12}$$ eqn:invK

We can also obtain $\mathbf{B}^{-1} = \mathbf{K} + \mathbf{A} - \mathbf{B}^{-T}$. Hence we have a representation of $\overline{\mathbf{B}}^{-1}$ by simple manipulations:

$$\overline{\mathbf{B}}^{-1} = (\mathbf{K} + \mathbf{A} - \mathbf{B}^{-T})\mathbf{K}^{-1}(\mathbf{K} + \mathbf{A} - \mathbf{B}^{-1}) = \mathbf{A} + (\mathbf{A} - \mathbf{B}^{-T})\mathbf{K}^{-1}(\mathbf{A} - \mathbf{B}^{-1}).$$

This identity and the definition of $\mathbf{B}$ immediately yield another important identity:

$$\left(\overline{\mathbf{B}}^{-1}\mathbf{v}, \mathbf{v}\right) = (\mathbf{A}\mathbf{v}, \mathbf{v}) + \left(\mathbf{K}^{-1}(\mathbf{D} + \mathbf{U} - \mathbf{S}^{-1})\mathbf{v}, (\mathbf{D} + \mathbf{U} - \mathbf{S}^{-1})\mathbf{v}\right), \quad \forall \mathbf{v} \in \mathbf{V}. \tag{4.13}$$

eqn:invBbar0

Now we apply a modification of Theorem 2.3 (i.e., general convergence rate estimate for SPD problems[4]) and get the following convergence result:

thm:MBGS **Theorem 4.2** (Convergence rate of modified block G-S)**.** *If* $\overline{\mathbf{S}} := \mathbf{S}^T + \mathbf{S} - \mathbf{S}^T\mathbf{D}\mathbf{S}$ *(or* $\mathbf{K} := \mathbf{S}^{-T} + \mathbf{S}^{-1} - \mathbf{D}$*) is SPD, then the modified block G-S method converges and*

$$\|\mathbf{I} - \mathbf{B}\mathbf{A}\|_{\mathbf{A}}^2 = 1 - \frac{1}{1 + c_0}, \quad \textit{with } c_0 := \sup_{\|\mathbf{v}\|_{\mathbf{A}}=1} \left\|\mathbf{K}^{-\frac{1}{2}}\left(\mathbf{D} + \mathbf{U} - \mathbf{S}^{-1}\right)\mathbf{v}\right\|^2.$$

4.3 Convergence analysis of SSC

sec:XZ

In the previous section, we have found that the SSC method for the original equation is equivalent to the block G-S method for the expanded equation using the same subspaces $\{V_j\}_{j=1}^J$. Now we use the same argument discussed in Chapter 2 to analyze the convergence rate of the block G-S method for the expanded system. In this way, we can give a convergence analysis for the successive subspace correction method. The proof here follows the discussion in [24].

A technical lemma

Suppose $V = \sum_{j=1}^J V_j$. It is clear that $\mathbf{\Pi} : \mathbf{V} \mapsto V$ is surjective and $\mathbf{\Pi}\mathbf{u} = \sum_{j=1}^J \mathcal{I}_j \mathbf{u}_j$. We have the following simple but useful lemma:

lemma:equiv **Lemma 4.2.** *If the iterator* $\mathbf{B}$ *in* (4.9) *is SPD, then* $\mathcal{B} = \mathbf{\Pi}\,\mathbf{B}\,\mathbf{\Pi}^T$ *is also SPD and*

$$(\mathcal{B}^{-1}v, v) = \inf_{\substack{\mathbf{v}\in\mathbf{V}\\ \mathbf{\Pi}\mathbf{v}=v}} (\mathbf{B}^{-1}\mathbf{v}, \mathbf{v}), \quad \forall v \in V.$$

Proof. It is clear that $(\mathcal{B}v, v) \geqslant 0$ for any $v \in V$ due to positive definiteness of $\mathbf{B}$. Furthermore, we have

$$0 = (\mathcal{B}v, v) = (\mathbf{B}\mathbf{\Pi}^T v, \mathbf{\Pi}^T v) \quad \Longrightarrow \quad \mathbf{\Pi}^T v = 0 \quad \Longrightarrow \quad v \in \mathrm{null}(\mathbf{\Pi}^T) = \mathrm{range}(\mathbf{\Pi})^{\perp}.$$

[4] In order to apply the convergence rate estimate Theorem 2.3 for stationary iterative methods to a symmetric positive semi-definite problem, we can restrict the domain of operator $\mathbf{A}$ inside the subspace range($\mathbf{A}$). This way the operator $\mathbf{A}$ is still non-singular.

Since $\mathbf{\Pi}$ is surjective, we have $v = 0$. This proves the iterator $\mathcal{B}$ is SPD.

Define $\mathbf{v}_* := \mathbf{B}\mathbf{\Pi}^T\mathcal{B}^{-1}v$. It is easy to see that

$$\mathbf{\Pi}\mathbf{v}_* = \mathbf{\Pi}\mathbf{B}\mathbf{\Pi}^T\mathcal{B}^{-1}v = \mathcal{B}\mathcal{B}^{-1}v = v, \quad \forall v \in V,$$

and

$$(\mathbf{B}^{-1}\mathbf{v}_*, \mathbf{w}) = (\mathbf{\Pi}^T\mathcal{B}^{-1}v, \mathbf{w}) = (\mathcal{B}^{-1}v, \mathbf{\Pi}\mathbf{w}).$$

If $\mathbf{w} \in \text{null}(\mathbf{\Pi})$, then $(\mathbf{B}^{-1}\mathbf{v}_*, \mathbf{w}) = 0$. This ensures that, for any vector $\mathbf{v} \in \mathbf{V}$, there exists a $\mathbf{B}^{-1}$-orthogonal decomposition $\mathbf{v} = \mathbf{v}_* + \mathbf{w}$ with $\mathbf{w} \in \text{null}(\mathbf{\Pi})$. Hence, we get

$$(\mathbf{B}^{-1}\mathbf{v}, \mathbf{v}) = \big(\mathbf{B}^{-1}(\mathbf{v}_* + \mathbf{w}), \mathbf{v}_* + \mathbf{w}\big) = \big(\mathbf{B}^{-1}\mathbf{v}_*, \mathbf{v}_*\big) + \big(\mathbf{B}^{-1}\mathbf{w}, \mathbf{w}\big).$$

Thus

$$\begin{aligned}\inf_{\substack{\mathbf{v}\in\mathbf{V}\\ \mathbf{\Pi}\mathbf{v}=v}} (\mathbf{B}^{-1}\mathbf{v}, \mathbf{v}) &= \big(\mathbf{B}^{-1}\mathbf{v}_*, \mathbf{v}_*\big) + \inf_{\mathbf{w}\in\text{null}(\mathbf{\Pi})} \big(\mathbf{B}^{-1}\mathbf{w}, \mathbf{w}\big) \\ &= \big(\mathbf{B}^{-1}\mathbf{v}_*, \mathbf{v}_*\big) = \big(\mathbf{\Pi}^T\mathcal{B}^{-1}v, \mathbf{B}\mathbf{\Pi}^T\mathcal{B}^{-1}v\big) = \big(\mathcal{B}^{-1}v, v\big).\end{aligned}$$

Hence the result. □

Remark 4.4 (Minimizer for the expanded problem)**.** From the above proof, we can easily see $\mathbf{v}_* = \mathbf{B}\mathbf{\Pi}^T\mathcal{B}^{-1}v$ is the minimizer.

Remark 4.5 (Auxiliary space problem)**.** The above lemma for relation between the expanded problem and the original problem can also be extended to the auxiliary space method: For two vector spaces V and $\tilde{V}$ and a surjective $\Pi : \tilde{V} \mapsto V$, if the iterator $\tilde{\mathcal{B}} : \tilde{V}' \mapsto \tilde{V}$ is SPD, then $\mathcal{B} = \Pi\,\tilde{\mathcal{B}}\,\Pi^T$ is also SPD and

$$(\mathcal{B}^{-1}v, v) = \inf_{\substack{\tilde{v}\in\tilde{V}\\ \Pi\tilde{v}=v}} (\tilde{\mathcal{B}}^{-1}\tilde{v}, \tilde{v}), \quad \forall v \in V.$$

The X-Z identity

We now give the well-known X-Z identity originally proved by Xu and Zikatanov [58] which gives the exact convergence rate of the SSC method.

:xzidentityc0 **Theorem 4.3** (X-Z Identity)**.** *Assume that $\mathcal{B}$ is defined by Algorithm 4.2 and, for $j = 1, \ldots, J$, $\mathbf{w}_j := \mathcal{A}_j\mathit{\Pi}_j\sum_{i\geqslant j}\mathbf{v}_i - \mathcal{S}_j^{-1}\mathbf{v}_j$. If $\mathcal{S}_j^{-T} + \mathcal{S}_j^{-1} - \mathcal{A}_j$ are SPD's for $j = 1, \ldots, J$, then*

$$\|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}}^2 = 1 - \frac{1}{1+c_0} = 1 - \frac{1}{c_1}, \tag{4.14}$$ eq:xzidentity

where

$$c_0 = \sup_{\|v\|_{\mathcal{A}}=1} \inf_{\sum_j \mathbf{v}_j = v} \sum_{j=1}^{J} \|\mathcal{S}_j^T\mathbf{w}_j\|_{\mathcal{S}_j^{-1}}^2 \tag{4.15}$$ eq:xzc0

and

$$c_1 = \sup_{\|v\|_{\mathcal{A}}=1} \inf_{\sum_j \mathbf{v}_j = v} \sum_{j=1}^{J} \left\| \overline{\mathcal{S}}_j \mathcal{S}_j^{-1} \mathbf{v}_j + \mathcal{S}_j^T \mathbf{w}_j \right\|_{\overline{\mathcal{S}}_j^{-1}}^2 . \tag{4.16}$$ eq:xzc1

Proof. By applying Theorem 2.3 and Lemma 4.2, we know

$$\|\mathcal{I} - \mathcal{B}\mathcal{A}\|_{\mathcal{A}}^2 = 1 - \left(\sup_{\|v\|_{\mathcal{A}}=1} (\overline{\mathcal{B}}^{-1} v, v) \right)^{-1} = 1 - \left(\sup_{\|v\|_{\mathcal{A}}=1} \inf_{\mathbf{\Pi}\mathbf{v}=v} (\overline{\mathbf{B}}^{-1}\mathbf{v}, \mathbf{v}) \right)^{-1}.$$

From (4.13) and (4.12), we have, for any $\mathbf{v} \in \mathbf{V}$, that

$$\left(\overline{\mathbf{B}}^{-1}\mathbf{v}, \mathbf{v}\right) = (\mathbf{A}\mathbf{v}, \mathbf{v}) + \left(\mathbf{K}^{-1}(\mathbf{D} + \mathbf{U} - \mathbf{S}^{-1})\mathbf{v},\ (\mathbf{D} + \mathbf{U} - \mathbf{S}^{-1})\mathbf{v}\right).$$

By simple calculation, we get

$$\begin{aligned} (\mathbf{D} + \mathbf{U})\,\mathbf{v} &= \Big(\sum_{j\geqslant 1} \mathcal{Q}_1 \mathcal{A} \mathcal{Q}_j^T \mathbf{v}_j,\ \sum_{j\geqslant 2} \mathcal{Q}_2 \mathcal{A} \mathcal{Q}_j^T \mathbf{v}_j,\ \cdots \Big)^T \\ &= \Big(\sum_{j\geqslant 1} \mathcal{A}_1 \mathit{\Pi}_1 \mathcal{I}_j \mathbf{v}_j,\ \sum_{j\geqslant 2} \mathcal{A}_2 \mathit{\Pi}_2 \mathcal{I}_j \mathbf{v}_j,\ \cdots \Big)^T \\ &= \Big(\mathcal{A}_1 \mathit{\Pi}_1 \sum_{j\geqslant 1} \mathbf{v}_j,\ \mathcal{A}_2 \mathit{\Pi}_2 \sum_{j\geqslant 2} \mathbf{v}_j,\ \cdots \Big)^T. \end{aligned}$$

We then have

$$\left(\mathbf{D} + \mathbf{U} - \mathbf{S}^{-1}\right)\mathbf{v} = (\mathbf{w}_1, \mathbf{w}_2, \ldots, \mathbf{w}_J)^T, \quad \text{with } \mathbf{w}_j := \mathcal{A}_j \mathit{\Pi}_j \sum_{i\geqslant j} \mathbf{v}_i - \mathcal{S}_j^{-1}\mathbf{v}_j.$$

Due to the fact that $\mathbf{K}$ is diagonal, we have

$$\left(\mathbf{K}^{-1}(\mathbf{D}+\mathbf{U}-\mathbf{S}^{-1})\mathbf{v},\ (\mathbf{D}+\mathbf{U}-\mathbf{S}^{-1})\mathbf{v}\right) = \sum_{j=1}^{J} \left(\left(\mathcal{S}_j^{-T} + \mathcal{S}_j^{-1} - \mathcal{A}_j\right)^{-1} \mathbf{w}_j, \mathbf{w}_j \right) = \sum_{j=1}^{J} \left\| \mathcal{S}_j^T \mathbf{w}_j \right\|_{\overline{\mathcal{S}}_j^{-1}}^2,$$

where $\overline{\mathcal{S}}_j := \mathcal{S}_j^T \left(\mathcal{S}_j^{-T} + \mathcal{S}_j^{-1} - \mathcal{A}_j\right)\mathcal{S}_j$ is the symmetrization of $\mathcal{S}_j$. We then obtain, for any $v \in V$, that

$$\sup_{\|v\|_{\mathcal{A}}=1} \inf_{\mathbf{\Pi}\mathbf{v}=v} \left(\overline{\mathbf{B}}^{-1}\mathbf{v}, \mathbf{v}\right) = 1 + \sup_{\|v\|_{\mathcal{A}}=1} \inf_{\mathbf{\Pi}\mathbf{v}=v} \sum_{j=1}^{J} \left\| \mathcal{S}_j^T \mathbf{w}_j \right\|_{\overline{\mathcal{S}}_j^{-1}}^2.$$

This gives the desired estimate for the constant c_0.

On the other hand, from (4.11), we have

$$\begin{aligned} \left(\overline{\mathbf{B}}^{-1}\mathbf{v}, \mathbf{v}\right) &= \left(\mathbf{K}^{-1}(\mathbf{S}^{-T} + \mathbf{U})\mathbf{v},\ (\mathbf{S}^{-T} + \mathbf{U})\mathbf{v}\right) \\ &= \sum_{j=1}^{J} \left\| \left(\mathcal{S}_j^{-1} + \mathcal{S}_j^{-T} - \mathcal{A}_j\right)^{-\frac{1}{2}} \left(\mathcal{S}_j^{-T}\mathbf{v}_j + \sum_{i>j} \mathcal{Q}_j \mathcal{A} \mathcal{I}_i \mathbf{v}_i\right) \right\|^2. \end{aligned} \tag{4.17}$$

We notice that

$$\begin{aligned}\mathcal{S}_j^{-T}\mathbf{v}_j+\sum_{i>j}\mathcal{Q}_j\mathcal{A}\mathcal{I}_i\mathbf{v}_i &= \mathcal{S}_j^{-T}\mathbf{v}_j+\mathcal{A}_j\Pi_j\sum_{i>j}\mathbf{v}_i=\big(\mathcal{S}_j^{-T}+\mathcal{S}_j^{-1}-\mathcal{A}_j\big)\mathbf{v}_j+\mathbf{w}_j\\ &= \mathcal{S}_j^{-T}\overline{\mathcal{S}}_j\mathcal{S}_j^{-1}\mathbf{v}_j+\mathbf{w}_j=\mathcal{S}_j^{-T}\Big(\overline{\mathcal{S}}_j\mathcal{S}_j^{-1}\mathbf{v}_j+\mathcal{S}_j^{T}\mathbf{w}_j\Big).\end{aligned}$$

Plug this into the previous identity, we get

$$\begin{aligned}\Big(\overline{\mathbf{B}}^{-1}\mathbf{v},\mathbf{v}\Big) &= \sum_{j=1}^{J}\Big\|\big(\mathcal{S}_j^{-1}+\mathcal{S}_j^{-T}-\mathcal{A}_j\big)^{-\frac{1}{2}}\big(\mathcal{S}_j^{-T}\mathbf{v}_j+\sum_{i>j}\mathcal{Q}_j\mathcal{A}\mathcal{I}_i\mathbf{v}_i\big)\Big\|^2\\ &= \sum_{j=1}^{J}\Big\|\overline{\mathcal{S}}_j\mathcal{S}_j^{-1}\mathbf{v}_j+\mathcal{S}_j^{T}\mathbf{w}_j\Big\|^2_{\overline{\mathcal{S}}_j^{-1}}.\end{aligned}$$

Hence the estimate for the constant c_1. □

Remark 4.6 (An equivalent form)**.** We have introduced operators $\mathcal{T}_j:=\mathcal{S}_j\mathcal{A}_j: V_j\mapsto V_j$. Hence $\mathcal{T}_{\overline{\mathcal{S}}_j}:=\overline{\mathcal{S}}_j\mathcal{A}_j=\mathcal{T}_j+\mathcal{T}_j^*-\mathcal{T}_j^*\mathcal{T}_j$ and we can rewrite the above estimate (4.16) in a slightly different form. Notice that, in (4.17),

$$\mathcal{S}_j^{-T}\mathbf{v}_j+\sum_{i>j}\mathcal{Q}_j\mathcal{A}\mathcal{I}_i\mathbf{v}_i=\mathcal{A}_j\big(\mathcal{S}_j^{T}\mathcal{A}_j\big)^{-1}\mathbf{v}_j+\mathcal{A}_j\Pi_j\sum_{i>j}\mathbf{v}_i=\mathcal{A}_j\Big[\big(\mathcal{T}_j^*\big)^{-1}\mathbf{v}_j+\Pi_j\sum_{i>j}\mathbf{v}_i\Big]$$

and

$$\big(\mathcal{S}_j^{-1}+\mathcal{S}_j^{-T}-\mathcal{A}_j\big)^{-1}\mathcal{A}_j=\big(\mathcal{T}_j^{-1}+(\mathcal{T}_j^*)^{-1}-\mathcal{I}_j\big)^{-1}=\mathcal{T}_j\mathcal{T}_{\overline{\mathcal{S}}_j}^{-1}\mathcal{T}_j^*.$$

Thus we have

$$c_1=\sup_{\|v\|_{\mathcal{A}}=1}\inf_{\sum_j\mathbf{v}_j=v}\sum_{j=1}^{J}\Big\|\mathcal{T}_{\overline{\mathcal{S}}_j}^{-\frac{1}{2}}\Big(\mathbf{v}_j+\mathcal{T}_j^*\Pi_j\sum_{i>j}\mathbf{v}_i\Big)\Big\|^2_{\mathcal{A}}. \tag{4.18}$$ eq:xzc1T

Example 4.3 (Simple stationary iterative method)**.** One-level linear stationary iterative method

$$u^{\text{new}}=u^{\text{old}}+\overline{\mathcal{S}}(f-\mathcal{A}u^{\text{old}}),$$

can be viewed as a special subspace correction method with only one subspace V. Hence, using (4.18), we immediately have

$$c_1=\sup_{\|v\|_{\mathcal{A}}=1}\|\mathcal{T}_{\overline{\mathcal{S}}}^{-\frac{1}{2}}v\|^2_{\mathcal{A}}=\sup_{\|v\|_{\mathcal{A}}=1}\big((\overline{\mathcal{S}}\mathcal{A})^{-1}v,v\big)_{\mathcal{A}}=\sup_{\|v\|_{\mathcal{A}}=1}\big(\overline{\mathcal{S}}^{-1}v,v\big),$$

which is exactly the convergence rate derived in Theorem 2.3. □

Example 4.4 (Two-grid method)**.** Theorem 3.3 can be viewed as a special case of the X-Z identity in the case of space decomposition with two subspaces, i.e., $V=V_c+V$. Suppose we use $\mathcal{A}_c^{-1}$ and $\overline{\mathcal{S}}$ as subspace solvers, respectively. According to (4.18), we get

$$c_1=\sup_{\|w\|_{\mathcal{A}}=1}\inf_{\substack{w=v_c+v\\ v_c\in V_c,v\in V}}\|v_c+\Pi_c v\|^2_{\mathcal{A}}+\|(\overline{\mathcal{S}}\mathcal{A})^{-\frac{1}{2}}v\|^2_{\mathcal{A}}.$$

We can prove that

$$c_1 = \sup_{\|v\|_{\mathcal{A}}=1} \inf_{v \in V_c^{\perp_{\mathcal{A}}}} \|(\overline{\mathcal{S}}\mathcal{A})^{-\frac{1}{2}} v\|_{\mathcal{A}}^2 = \sup_{\|v\|_{\mathcal{A}}=1} \inf_{v \in V_c^{\perp_{\mathcal{A}}}} \left(\overline{\mathcal{S}}^{-1} v, v\right).$$

Hence, we immediately obtain

$$c_1 = \sup_{\|v\|_{\mathcal{A}}=1} \|(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})v\|_{\overline{\mathcal{S}}^{-1}}^2 = \sup_{\|v\|_{\mathcal{A}}=1} \|\mathcal{T}_{\overline{\mathcal{S}}}^{-\frac{1}{2}}(\mathcal{I} - \mathcal{Q}_{\overline{\mathcal{S}}^{-1}})v\|_{\mathcal{A}}^2,$$

which is consistent with the X-Z identity. □

cor:XZ-exact **Corollary 4.1** (SSC with exact subspace solvers). *If an exact subspace solver $\mathcal{S}_j = \mathcal{A}_j^{-1}$ for each subspace is used, then we have, in (4.14), that*

$$c_0 = \sup_{\|v\|_{\mathcal{A}}=1} \inf_{\sum_j \mathbf{v}_j = v} \sum_{j=1}^{J} \left\| \Pi_j \sum_{i>j} \mathbf{v}_i \right\|_{\mathcal{A}_j}^2 \tag{4.19}$$ eq:c0p

and

$$c_1 = \sup_{\|v\|_{\mathcal{A}}=1} \inf_{\sum_j \mathbf{v}_j = v} \sum_{j=1}^{J} \left\| \Pi_j \sum_{i\geqslant j} \mathbf{v}_i \right\|_{\mathcal{A}_j}^2. \tag{4.20}$$ eq:c1p

4.4 Convergence analysis of PSC

sec:ConvPSC

In this section, we estimate the condition number of the PSC method.

Condition number of PSC

To obtain estimates on the condition number of the preconditioned problems, we first give the following assumptions:

Assump:MSC **Assumption 4.1** (Convergence assumptions for MSC). We assume that

1. For any $v \in V$, there exists a decomposition $v = \sum_{j=1}^{J} v_j$ with $v_j \in V_j$ such that

$$\sum_{j=1}^{J} \left(\mathcal{S}_j^{-1} v_j, v_j\right) \leqslant K_1 (\mathcal{A}v, v); \tag{4.21}$$ assump:MSC1

2. For any $u, v \in V$,

$$\sum_{(i,j)} \left(\mathcal{T}_i u, \mathcal{T}_j v\right)_{\mathcal{A}} \leqslant K_2 \left(\sum_{i=1}^{J} (\mathcal{T}_i u, u)_{\mathcal{A}}\right)^{\frac{1}{2}} \left(\sum_{j=1}^{J} (\mathcal{T}_j v, v)_{\mathcal{A}}\right)^{\frac{1}{2}}. \tag{4.22}$$ assump:MSC2

thm:PSC **Theorem 4.4** (Condition number of PSC). *If Assumption 4.1 holds true, the PSC method (4.7) satisfies*

$$\kappa(\mathcal{B}\mathcal{A}) \leqslant K_1 K_2.$$

Proof. For any $v \in V$, suppose that $v = \sum_{j=1}^{J} v_j$ is a decomposition satisfies the first condition of Assumption 4.1. It is easy to see that

$$\begin{aligned}
(v,v)_{\mathcal{A}} &= \sum_{j=1}^{J}(v_j,v)_{\mathcal{A}} = \sum_{j=1}^{J}(v_j,\Pi_j v)_{\mathcal{A}} = \sum_{j=1}^{J}(v_j,\mathcal{A}_j\Pi_j v) = \sum_{j=1}^{J}\left(\mathcal{S}_j^{-\frac{1}{2}}v_j,\mathcal{S}_j^{\frac{1}{2}}\mathcal{A}_j\Pi_j v\right)\\
&\leqslant \sum_{j=1}^{J}\left(\mathcal{S}_j^{-1}v_j,v_j\right)^{\frac{1}{2}}\left(\mathcal{S}_j\mathcal{A}_j\Pi_j v,\mathcal{A}_j\Pi_j v\right)^{\frac{1}{2}} = \sum_{j=1}^{J}\left(\mathcal{S}_j^{-1}v_j,v_j\right)^{\frac{1}{2}}\left(\mathcal{S}_j\mathcal{A}_j\Pi_j v,v\right)_{\mathcal{A}}^{\frac{1}{2}}\\
&\leqslant \left(\sum_{j=1}^{J}\left(\mathcal{S}_j^{-1}v_j,v_j\right)\right)^{\frac{1}{2}}\left(\sum_{j=1}^{J}\left(\mathcal{T}_j v,v\right)_{\mathcal{A}}\right)^{\frac{1}{2}} \leqslant \sqrt{K_1}\,\|v\|_{\mathcal{A}}\left(\mathcal{B}\mathcal{A}v,v\right)_{\mathcal{A}}^{\frac{1}{2}}.
\end{aligned}$$

Consequently, we have the lower bound

$$\frac{1}{K_1}(v,v)_{\mathcal{A}} \leqslant (\mathcal{B}\mathcal{A}v,v)_{\mathcal{A}}, \quad \forall\, v \in V.$$

From the second assumption, we have

$$\|\mathcal{B}\mathcal{A}v\|_{\mathcal{A}}^2 = \sum_{i,j=1}^{J}\left(\mathcal{T}_i v,\mathcal{T}_j v\right)_{\mathcal{A}} \leqslant K_2(\mathcal{B}\mathcal{A}v,v)_{\mathcal{A}} \leqslant K_2\|\mathcal{B}\mathcal{A}v\|_{\mathcal{A}}\|v\|_{\mathcal{A}}.$$

So we obtain the upper bound

$$(\mathcal{B}\mathcal{A}v,v)_{\mathcal{A}} \leqslant K_2(v,v)_{\mathcal{A}}, \quad \forall\, v \in V.$$

Thus Lemmas 2.5 and 2.6 yield the desired estimate. □

According to Theorem 4.4, if we can find a space decomposition and corresponding smoothers with uniform constants K_1 and K_2, then we are able to construct a uniformly convergent preconditioner using the PSC framework.

Remark 4.7 (Similar estimate for SSC)**.** In fact, with the same assumptions (Assumption 4.1), we can also show that the SSC method also converges with

$$\|\mathcal{I}-\mathcal{B}\mathcal{A}\|_{\mathcal{A}}^2 \leqslant 1-\frac{2-\omega_1}{K_1(1+K_2)^2} \qquad \text{and} \qquad \omega_1 := \max_j \rho(\mathcal{S}_j\mathcal{A}_j) = \max_j \rho(\mathcal{T}_j).$$

Because a sharp result has been given in §4.3, we will just leave the proof to the readers (cf., for example, [55]).

Estimates of K_1 and K_2

Assumption 4.1 is not easy to verify directly. So we now give a few useful estimates for the constants in these conditions. We first give a straight-forward estimate of K_1, which clearly separates the condition on space decomposition part and smoother part. We leave the proof of the following lemma to the readers; see HW 4.5.

lem:K1 **Lemma 4.3** (Estimates of K_1)**.** *Assume that, for any $v \in V$, there is a decomposition $v = \sum_{j=1}^{J} v_j$ with $v_j \in V_j$:*

(i) If the decomposition satisfies that

$$\sum_{j=1}^{J} (v_j, v_j)_{\mathcal{A}} \leqslant C_1 (v, v)_{\mathcal{A}},$$

then we have

$$K_1 \leqslant C_1/\omega_0, \quad \text{where } \omega_0 := \min_{j=1,\dots,J} \{\lambda_{\min}(\mathcal{S}_j \mathcal{A}_j)\};$$

(ii) If $\rho_j := \rho(\mathcal{A}_j)$ and

$$\sum_{j=1}^{J} \rho_j (v_j, v_j) \leqslant \hat{C}_1 (v, v)_{\mathcal{A}},$$

then we have

$$K_1 \leqslant \hat{C}_1/\hat{\omega}_0, \quad \text{where } \hat{\omega}_0 := \min_{j=1,\dots,J} \{\rho_j \lambda_{\min}(\mathcal{S}_j)\}.$$

We introduce a nonnegative symmetric matrix

$$\Sigma = (\sigma_{i,j}) \in \mathbb{R}^{J \times J}, \tag{4.23}$$ eqn:Sigma

where each entry $\sigma_{i,j}$ is the smallest constant such that

$$(\mathcal{T}_i u, \mathcal{T}_j v)_{\mathcal{A}} \leqslant \omega_1 \sigma_{i,j} (\mathcal{T}_i u, u)_{\mathcal{A}}^{\frac{1}{2}} (\mathcal{T}_j v, v)_{\mathcal{A}}^{\frac{1}{2}}, \quad \forall\, u, v \in V. \tag{4.24}$$ eqn:Strengthe

It is clear that $0 \leqslant \sigma_{i,j} \leqslant 1$. Moreover, $\sigma_{i,j} = 0$, if $\Pi_i \Pi_j = 0$.

lem:K2 **Lemma 4.4** (Estimate of K_2)**.** *The constant $K_2 \leqslant \omega_1 \rho(\Sigma)$. Furthermore, if $\sigma_{i,j} \lesssim \gamma^{|i-j|}$ holds for some parameter $0 < \gamma < 1$, then $\rho(\Sigma) \lesssim (1-\gamma)^{-1}$; in this case, the inequality (4.22) is the well-known* strengthened Cauchy-Schwarz inequality*.*

Proof. From the definition of Σ as in (4.23), it is immediately clear that $K_2 \leqslant \omega_1 \rho(\Sigma)$. Furthermore, because the matrix Σ is a real symmetric matrix and $\rho(\Sigma) \leqslant \max_{j=1,\dots,J} \sum_{i=1}^{J} \sigma_{i,j}$, we have

$$\rho(\Sigma) \leqslant \max_{1 \leqslant j \leqslant J} \sum_{i=1}^{J} \sigma_{i,j} \lesssim \sum_{i=1}^{J} \gamma^{i-1} \leqslant \frac{1}{1-\gamma}.$$

Hence the result. □

Auxiliary space method ⋆

Sometimes, we cannot apply subspace correction methods directly due to difficulties in obtaining an appropriate space decomposition. In this case, we can introduce an auxiliary or fictitious space $\tilde{V}$ for assistance. If $\Pi : \tilde{V} \mapsto V$ is surjective and satisfies the following two conditions: Firstly,

$$\|\Pi\tilde{v}\|_{\mathcal{A}} \leqslant \mu_1 \|\tilde{v}\|_{\tilde{\mathcal{A}}}, \qquad \forall\, \tilde{v} \in \tilde{V}.$$

Secondly, for any $v \in V$, there exists $\tilde{v} \in \tilde{V}$ such that $\Pi\tilde{v} = v$ and

$$\mu_0 \|\tilde{v}\|_{\tilde{\mathcal{A}}} \leqslant \|v\|_{\mathcal{A}}, \qquad \forall\, \tilde{v} \in \tilde{V}.$$

Under the above assumptions, if $\tilde{B}$ is a SPD preconditioner for $\tilde{A}$, then $\mathcal{B} = \Pi\tilde{\mathcal{B}}\Pi^T$ is SPD and

$$\kappa(\mathcal{B}\mathcal{A}) \leqslant \left(\frac{\mu_1}{\mu_0}\right)^2 \kappa(\tilde{\mathcal{B}}\tilde{\mathcal{A}}).$$

This suggests that we can construct a subspace correction method on $\tilde{V}$ instead of the original space V. This simple result is sometimes called the *Fictitious Space Lemma*; see [47, 56].

4.5 Homework problems

HW 4.1. Prove the statements in Remark 4.3.

hw:projections **HW 4.2.** Prove Lemma 4.1.

hw:B-spd **HW 4.3.** If $\mathcal{S}_j$ $(j = 1, \ldots, J)$ are all SPD, then the preconditioner $\mathcal{B} = \sum_{j=1}^{J} \mathcal{S}_j \mathcal{Q}_j$ is also SPD.

hw:bGS **HW 4.4.** Show that the block G-S method for the expanded system is just the SSC method for the original problem.

hw:lemK1 **HW 4.5.** Prove the estimates in Lemma 4.3.

Part II

Multilevel Iterative Methods and Their Applications

Chapter 5

Multilevel Subspace Correction Preconditioners

ch:examples

In Chapter 4, we have discussed stationary iterative methods in the framework of method of subspace correction (MSC). In this chapter, we give a few examples of multilevel methods and their convergence analysis based on the framework of subspace corrections.

5.1 Two-grid overlapping DDM ⋆

In this section, we will investigate the two-level overlapping domain decomposition method in Chapter 2 using the MSC framework.

Two-level space decomposition

Based on the previous discussions, it is now easy to understand that the additive and multiplicative Schwarz domain decomposition methods can be considered as PSC and SSC, respectively. For proof-of-concept, we use the Poisson's equation on Ω as an example. In this case, $\mathscr{V} = H_0^1(\Omega)$, $\Omega = \bigcup_{j=1}^J \Omega_j$, and $\mathscr{V}_j := \{v \in \mathscr{V} : \operatorname{supp} v \subset \hat{\Omega}_j\} \subset \mathscr{V}$; see Figure 2.2. We define a finite-dimensional coarse space $V_0 \subset \mathscr{V}$ of meshsize $H = \operatorname{diam}(\Omega_j)$. Apparently, we have a space decomposition

$$\mathscr{V} = V_0 + \mathscr{V}_1 + \cdots + \mathscr{V}_J.$$

The SSC method based on this space decomposition with exact sub-problem solvers for each sub-domain as well as the coarse space gives an abstract multiplicative Schwarz DDM method.

We give a partition of unity $\theta_j \in C^1(\overline{\Omega})$ $(j = 1, \ldots, J)$ such that

(1) $0 \leqslant \theta_j \leqslant 1$ and $\sum_{j=1}^J \theta_j = 1$;

(2) $\operatorname{supp}\theta_j \subset \hat{\Omega}_j$;

(3) $\max|\nabla\theta_j| \leqslant C_\beta/H$, where C_β depends on the relative overlap size β.

This way, for any function $v \in \mathscr{V}$, we have a decomposition

$$v = v_0 + v_1 + \cdots + v_J,$$

where

$$v_0 \in V_0 \qquad \text{and} \qquad v_j := \theta_j(v - v_0) \in \mathscr{V}_j,\ j = 1, \ldots, J.$$

Thus $\sum_{j=1}^J v_j = v - v_0$ and

$$\begin{aligned}\sum_{j=0}^J \Big|\Pi_j \sum_{i=j+1}^J v_i\Big|_1^2 &= \sum_{j=0}^J \Big|\Pi_j \sum_{i=j+1}^J \theta_i(v-v_0)\Big|_1^2 \\ &= \Big|\Pi_0(v-v_0)\Big|_1^2 + \sum_{j=1}^J \Big|\Pi_j \sum_{i=j+1}^J \theta_i(v-v_0)\Big|_1^2.\end{aligned}$$

Convergence analysis of DDM

Since $\Pi_j : \mathscr{V} \mapsto \mathscr{V}_j$ is a $\mathcal{A}$-projection for $j = 1, \ldots, J$, it is easy to see that $|\Pi_j(v-v_0)|_1 \leqslant |v-v_0|_1$. Furthermore,

$$\begin{aligned}\Big|\Pi_j \sum_{i=j+1}^J \theta_i(v-v_0)\Big|_1^2 &= \Big|\Pi_j \sum_{i=j+1}^J \theta_i(v-v_0)\Big|_{1,\hat{\Omega}_j}^2 \leqslant \Big|\sum_{i=j+1}^J \theta_i(v-v_0)\Big|_{1,\hat{\Omega}_j}^2 \\ &\leqslant \Big\|(\sum_{i>j}\theta_i)\nabla(v-v_0)\Big\|_{0,\hat{\Omega}_j}^2 + \Big\|\nabla(\sum_{i>j}\theta_i)(v-v_0)\Big\|_{0,\hat{\Omega}_j}^2 \\ &\leqslant |v-v_0|_{1,\hat{\Omega}_j}^2 + \frac{C_\beta}{H}\|v-v_0\|_{0,\hat{\Omega}_j}^2.\end{aligned}$$

Summing up all the terms, we have

$$\begin{aligned}\sum_{j=0}^J \Big|\Pi_j \sum_{i=j+1}^J v_i\Big|_1^2 &\leqslant |v-v_0|_1^2 + \sum_{j=1}^J |v-v_0|_{1,\hat{\Omega}_j}^2 + C_\beta^2 H^{-2}\sum_{j=1}^J \|v-v_0\|_{0,\hat{\Omega}_j}^2 \\ &\lesssim |v-v_0|_1^2 + C_\beta^2 H^{-2}\|v-v_0\|_0^2,\end{aligned}$$

where the constant in the last inequality depends on the maximal number of overlaps in domain decomposition. Because v_0 could be any function in V_0, we can choose $v_0 = \mathcal{Q}_0 v$ and, in view of Proposition 3.5, obtain

$$\sum_{j=0}^J \Big|\Pi_j \sum_{i=j+1}^J v_i\Big|_1^2 \lesssim |v|_1^2.$$

Using the X-Z identity (Corollary 4.1), we get the following result. We leave the full proof to the readers; see HW 5.1.

prop:ConvDDM **Proposition 5.1** (Uniform convergence of two-level DDM)**.** *The abstract domain decomposition method with coarse space correction converges uniformly.*

:DDM-onelevel **Remark 5.1** (DDM without coarse space)**.** From the above analysis, we immediately see the importance of having the coarse space V_0. With the same proof, one can show that the convergence rate depends on H^{-2} if without the coarse space correction.

5.2 HB preconditioner

sec:HB

In the previous section, we have seen a two-level domain decomposition method in the setting of subspace correction. Now we investigate a multilevel example.

Nested space decomposition

We consider the Poisson's equation on a sequence of nested meshes $\mathcal{M}_l$ $(l = 0, \ldots, L)$ generated from an initial mesh $\mathcal{M}_0$ by uniform regular refinements. Hence meshsize h_l of $\mathcal{M}_l$ is proportional to γ^{2l} with $\gamma \in (0, 1)$. For example, in Figure 1.5, there is a hierarchy of grids with $h_l = (1/2)^{l+1}$ $(l = 0, 1, \ldots, L)$. Clearly,

$$h_0 > h_1 > h_2 > \cdots > h_L =: h.$$

Define continuous piecewise linear finite element spaces on each mesh as

$$V_l := \big\{ v \in \mathscr{V} \,:\, v|_\tau \in \mathcal{P}_1(\tau),\ \forall\, \tau \in \mathcal{M}_l \big\}. \tag{5.1}$$ eqn:SpaceV

This way, we build a nested subspaces

$$V_0 \subset V_1 \subset \cdots \subset V_L =: V \subset \mathscr{V} = H_0^1(\Omega).$$

The set of interior grid points on the l-th level is denoted as $x_{l,i} \in \mathring{G}(\mathcal{M}_l)$ $(i = 1, \ldots, n_l)$. The subspace V_l is assigned with a nodal basis $\{\phi_{l,i}\}_{i=1}^{n_l}$, where $n_l := \big|\mathring{G}(\mathcal{M}_l)\big|$. The space V_l can be further decomposed as the sum of the one-dimensional subspaces spanned with the nodal basis $V_{l,i} := \operatorname{span}\{\phi_{l,i}\}$ $(i = 1, \ldots, n_l)$. In this way, we obtain a natural multilevel space decomposition

$$V = \sum_{l=0}^{L} V_l = \sum_{l=0}^{L} \sum_{i=1}^{n_l} V_{l,i}. \tag{5.2}$$ eqn:MGdecomp

We can also define

$$W_l := \big\{ v \in V_l \,:\, v(x) = 0,\ \forall\, x \in \mathring{G}(\mathcal{M}_{l-1}) \big\} \tag{5.3}$$ eqn:SpaceW

and obtain a decomposition

$$V = W_0 \oplus W_1 \oplus \cdots \oplus W_L. \tag{5.4}$$ eqn:DecompW

Let $\mathcal{J}_l : V \mapsto V_l$ be the interpolation operator and define $\mathcal{J}_{-1} := 0$. It is easy to see that

$$W_l = (\mathcal{J}_l - \mathcal{J}_{l-1})V = (\mathcal{I} - \mathcal{J}_{l-1})V_l, \quad l = 0, \ldots, L.$$

Notice that the decomposition (5.4) is a direct sum and there is no redundancy in this decomposition at all.

For level $l = 0, \ldots, L$, we define a nodal basis

$$\psi_{l,i}(x) = \phi_{l,i}(x), \quad \text{for } x_{l,i} \in \mathring{G}(\mathcal{M}_l)\backslash\mathring{G}(\mathcal{M}_{l-1}),\ i = 1, \ldots, m_l := n_l - n_{l-1}.$$

Apparently, $\sum_{l=0}^{L} m_l = n_L = N$ and this basis

$$\{\psi_{l,i}(x) \,:\, i = 1, \ldots, m_l,\ \ l = 0, \ldots, L\} \tag{5.5}$$ eqn:HB

is called *hierarchical basis*.

Hierarchical basis preconditioner

We now use the Richardson iteration discussed in §3.3 as the subspace solver, i.e.,

$$\mathcal{S}_{l,i}\mathcal{Q}_{l,i}v = h_l^{2-d}\big(\mathcal{Q}_{l,i}v, \psi_{l,i}\big)\,\psi_{l,i} = h_l^{2-d}\big(v, \psi_{l,i}\big)\,\psi_{l,i}.$$

The PSC method based on the space decomposition (5.4) can then be written

$$\mathcal{B}_{\mathrm{HB}} r = \sum_{j=1}^{N} \mathcal{S}_j \mathcal{Q}_j r = \sum_{l=0}^{L} \left(h_l^{2-d} \sum_{i=1}^{m_l} (r, \psi_{l,i})\psi_{l,i} \right). \tag{5.6}$$ eqn:HBprecond

And this is the explicit form of the well-known *hierarchical basis* (HB) *preconditioner* proposed by Yserentant [62].

We shall now analyze this preconditioner in the framework of PSC in §4.4. In order to do that, we need a few important estimates.

:StableInterp **Lemma 5.1** (Stability of interpolation in H^1)**.** *We have*

$$\big\|(\mathcal{J}_l - \mathcal{J}_{l-1})v\big\|_0^2 + h_l^2\big|\mathcal{J}_l v\big|_1^2 \lesssim c_d(l) h_l^2 |v|_1^2, \quad \forall v \in V,$$

where $c_1(l) \equiv 1$, $c_2(l) = L - l$, *and* $c_3(l) = \gamma^{-2(L-l)}$.

Proof. Using Proposition 3.2, we have

$$\|(\mathcal{J}_l - \mathcal{J}_{l-1})v\|_0 = \|\mathcal{J}_l v - \mathcal{J}_{l-1}\mathcal{J}_l v\|_0 \lesssim h_l |\mathcal{J}_l v|_1.$$

Let $\tau \in \mathcal{M}_l$ and $v_\tau := |\tau|^{-1}\int_\tau v\,dx$ be the average of v on τ. Using the standard scaling argument for $|\cdot|_{1,\tau}$, the discrete Sobolev inequality Proposition 3.4, and the Poincaré inequality Proposition 1.2, we can obtain that

$$|\mathcal{J}_l v|_{1,\tau} = |\mathcal{J}_l v - v_\tau|_{1,\tau} \lesssim \|\mathcal{J}_l v - v_\tau\|_{\infty,\tau} \leqslant \|v - v_\tau\|_{\infty,\tau} \lesssim C_d\|v - v_\tau\|_{1,\tau} \lesssim C_d|v|_{1,\tau}.$$

Hence the desired result follows by summing up terms on all elements in $\mathcal{M}_l$. □

Remark 5.2 (Condition number in hierarchical basis)**.** The above lemma suggests that, if $v \in W_l$ for any $0 \leqslant l \leqslant L$, we have

$$c_d^{-1}(l)h_l^{-2}(v,v) \lesssim a[v,v].$$

Compare this with the general Poincaré inequality in Proposition 1.3. Furthermore, from the inverse inequality Proposition 3.3, we always have

$$a[v,v] = |v|_1^2 \lesssim h_l^{-2}\|v\|_0^2 = h_l^{-2}(v,v).$$

Hence the operator $\mathcal{A}_l$ is "well-conditioned" up to a constant $c_d(l)$; compare this property with the standard Lagrangian finite element basis case in Remark 3.4.

Strengthened Cauchy-Schwarz inequality

lem:SCS1 **Lemma 5.2** (Inner product between two levels)**.** *Let $i \leqslant j$; then we have*

$$a[u,v] \lesssim \gamma^{j-i}h_j^{-1}|u|_1\|v\|_0, \quad \forall\, u \in V_i, v \in V_j.$$

Proof. We first restrict our attention to an element $\tau_i \in \mathcal{M}_i$. For $v \in \mathcal{M}_j$, there is a unique function $v_1 \in V$, such that v_1 vanishes on $\partial\tau_i$ and equals to v at all other grid points. Let $v_0 := v - v_1$. Because $u \in W_i$ is a linear function on τ_i, we have $\int_{\tau_i}\nabla u\nabla v_1 = 0$.

Define $T := \bigcup_{\tau_j\in\mathcal{M}_j,\overline{\tau}_j\bigcap\partial\tau_i\neq\emptyset}\tau_j$. Then $|T| \cong \left(\frac{h_i}{h_j}\right)^{d-1}h_j^d = h_i^{d-1}h_j$ and $\operatorname{supp} v_0 \subset \overline{T}$. We have

$$\|\nabla v_0\|_{0,\tau_i}^2 \lesssim \sum_{x\in\mathring{G}(\mathcal{M}_j)\bigcap\partial\tau_i} h_j^d h_j^{-2}\, v_0^2(x) = \sum_{x\in\mathring{G}(\mathcal{M}_j)\bigcap\partial\tau_i} h_j^{d-2}\, v^2(x) \lesssim h_j^{-2}\|v\|_{0,\tau_i}^2.$$

Since ∇u is a constant on τ_i, we have

$$\|\nabla u\|_{0,T} = \frac{|T|^{1/2}}{|\tau_i|^{1/2}}\|\nabla u\|_{0,\tau_i} \lesssim \left(\frac{h_i^{d-1}h_j}{h_i^d}\right)^{1/2}\|\nabla u\|_{0,\tau_i} \lesssim \gamma^{j-i}|u|_{1,\tau_i}.$$

Combining the above two inequalities, we have

$$\int_{\tau_i}\nabla u\cdot\nabla v = \int_{\tau_i}\nabla u\cdot\nabla v_0 \lesssim \gamma^{j-i}h_j^{-1}|u|_{1,\tau_i}\|v\|_{0,\tau_i}, \quad \forall\tau_i \in \mathcal{M}_i.$$

By the Cauchy-Schwarz inequality, we obtain the estimate:

$$\begin{aligned} a[u,v] &= \sum_{\tau_i\in\mathcal{M}_i}\int_{\tau_i}\nabla u\cdot\nabla v \lesssim \gamma^{j-i}h_j^{-1}\sum_{\tau_i\in\mathcal{M}_i}|u|_{1,\tau_i}\|v\|_{0,\tau_i} \\ &\leqslant \gamma^{j-i}h_j^{-1}\Big(\sum_{\tau_i\in\mathcal{M}_i}|u|_{1,\tau_i}\Big)^{1/2}\Big(\sum_{\tau_i\in\mathcal{M}_i}\|v\|_{0,\tau_i}\Big)^{1/2} = \gamma^{j-i}h_j^{-1}|u|_1\|v\|_0. \end{aligned}$$

Hence the result. □

lem:SCS2 **Lemma 5.3** (Strengthened Cauchy-Schwarz inequality for interpolation)**.** *If $u,v\in V$, let $u_i := (\mathcal{J}_i-\mathcal{J}_{i-1})u$, and $v_j := (\mathcal{J}_j-\mathcal{J}_{j-1})v$, then we have*

$$a[u_i,v_j]\lesssim\gamma^{|i-j|}\|u_i\|_{\mathcal{A}}\|v_j\|_{\mathcal{A}}.$$

Proof. If $j\geqslant i$, we have $v_j = v_j-\mathcal{J}_{j-1}v_j$. So $\|v_j\|_0 = \|v_j-\mathcal{J}_{j-1}v_j\|_0\lesssim h_j\|v_j\|_{\mathcal{A}}$ follows from Proposition 3.2. If $i\geqslant j$, we can argue in a similar way. Hence the result follows directly from Lemma 5.2. □

lem:SCS3 **Lemma 5.4** (Estimating K_2)**.** *Assume that $\mathcal{T}_j=\mathcal{S}_j\mathcal{A}_j\Pi_j$ and the subspace smoother $\mathcal{S}_j: V_j\mapsto V_j$ satisfies*

$$\|\mathcal{S}_j\mathcal{A}_jv\|_0^2\lesssim\rho_j^{-1}(\mathcal{A}_jv,v),\quad\forall\, v\in V_j,$$

where $\rho_j := \rho(\mathcal{A}_j)$. Then, if $i<j$, we have

$$(u_i,\mathcal{T}_jv)_{\mathcal{A}}\lesssim\gamma^{j-i}\|u_i\|_{\mathcal{A}}\|v\|_{\mathcal{A}},\quad\forall\, u_i\in V_i, v\in V. \tag{5.7}$$ eqn:SCSIneq1

For $0\leqslant i,j\leqslant L$, we have the strengthened Cauchy-Schwarz inequality

$$(\mathcal{T}_iu,\mathcal{T}_jv)_{\mathcal{A}}\lesssim\gamma^{|j-i|/2}(\mathcal{T}_iu,u)_{\mathcal{A}}^{\frac12}(\mathcal{T}_jv,v)_{\mathcal{A}}^{\frac12},\quad\forall\, u,v\in V. \tag{5.8}$$ eqn:SCSIneq2

Proof. By applying Lemma 5.2, we get

$$(u_i,\mathcal{T}_jv)_{\mathcal{A}} = a[u_i,\mathcal{T}_jv]\lesssim\gamma^{j-i}h_j^{-1}\|u_i\|_{\mathcal{A}}\|\mathcal{T}_jv\|_0.$$

Furthermore, we have

$$\|\mathcal{T}_jv\|_0=\|\mathcal{S}_j\mathcal{A}_j\Pi_jv\|_0\lesssim h_j\|\mathcal{A}_j^{1/2}\Pi_jv\|_0\leqslant h_j\|\Pi_jv\|_{\mathcal{A}}\leqslant h_j\|v\|_{\mathcal{A}}.$$

This proves the first inequality (5.7).

First consider the case when $j\geqslant i$. By the Cauchy-Schwarz inequality and the inequality (5.7), we get

$$(\mathcal{T}_iu,\mathcal{T}_jv)_{\mathcal{A}}\leqslant(\mathcal{T}_j\mathcal{T}_iu,\mathcal{T}_iu)_{\mathcal{A}}^{\frac12}(\mathcal{T}_jv,v)_{\mathcal{A}}^{\frac12}\lesssim\gamma^{(j-i)/2}\|\mathcal{T}_iu\|_{\mathcal{A}}(\mathcal{T}_jv,v)_{\mathcal{A}}^{\frac12}.$$

Also observe that $(\mathcal{T}_iu,\mathcal{T}_iu)_{\mathcal{A}}\lesssim\|\mathcal{T}_iu\|_{\mathcal{A}}(\mathcal{T}_iu,u)_{\mathcal{A}}^{\frac12}$ and the second inequality (5.8) follows immediately. □

Convergence analysis of HB preconditioner ⋆

thm:rateHB **Theorem 5.1** (Convergence of HB preconditioner)**.** *The multilevel PSC preconditioner $\mathcal{B}_{HB}$ defined in (5.6) satisfies*

$$\kappa(\mathcal{B}_{HB}\mathcal{A}) \lesssim C_d(h),$$

where $C_1(h) \equiv 1$, $C_2(h) = |\log h|^2$, and $C_3(h) = h^{-1}$.

Proof. We choose a decomposition $v = \sum_{l=0}^{L} v_l := \sum_{l=0}^{L} (\mathcal{J}_l - \mathcal{J}_{l-1})v$, where $\mathcal{J}_{-1} = 0$. With careful calculations, Proposition 3.3 and Lemma 5.1 ($\mathcal{J}_l = \Pi_l$ in 1D) yield

$$\sum_{l=0}^{L} \|v_l\|_{\mathcal{A}}^2 \lesssim \sum_{l=0}^{L} h_l^{-2} \|v_l\|_0^2 \eqsim \sum_{l=0}^{L} \rho_l \|v_l\|_0^2 \lesssim C_d(h) \|v\|_{\mathcal{A}}^2. \tag{5.9}$$ HB:lowerbd

On the other hand, we know $\hat{\omega}_0 = \min_l \rho_l \lambda_{\min}(\mathcal{S}_l) \lesssim 1$. Therefore $K_1 \lesssim C_d(h)$ due to Lemma 4.3. The strengthened Cauchy-Schwarz inequality (5.8) and Lemma 4.4 give that $K_2 \lesssim 1$. The convergence result then follows directly from the general theory in Theorem 4.4. □

Define an operator $\mathcal{H} : V \mapsto V$ such that

$$(\mathcal{H}v, w) := \sum_{l=0}^{L} \sum_{x_i \in \mathring{G}(\mathcal{M}_l) \backslash \mathring{G}(\mathcal{M}_{l-1})} h_l^{d-2} \Big((\mathcal{J}_l v - \mathcal{J}_{l-1} v)(x_i),\, (\mathcal{J}_l w - \mathcal{J}_{l-1} w)(x_i) \Big).$$

Hence we get

$$(\mathcal{H}v, v) = \sum_{l=0}^{L} \sum_{x_i \in \mathring{G}(\mathcal{M}_l) \backslash \mathring{G}(\mathcal{M}_{l-1})} h_l^{d-2} \Big| (\mathcal{J}_l v - \mathcal{J}_{l-1} v)(x_i) \Big|^2, \quad \forall\, v \in V.$$

This operator is in fact the inverse of the HB preconditioner, i.e., $\mathcal{H} = \mathcal{B}_{\mathrm{HB}}^{-1}$; see [63]. In fact, in the proof of Theorem 5.1, we have shown the following norm equivalence result:

$$\|v\|_{\mathcal{A}}^2 \lesssim (\mathcal{H}v, v) = \sum_{l=0}^{L} h_l^{-2} \|(\mathcal{J}_l - \mathcal{J}_{l-1})v\|_0^2 \lesssim C_d(h) \|v\|_{\mathcal{A}}^2. \tag{5.10}$$ eqn:HBYser

Let Π_l be the $(\cdot,\cdot)_{\mathcal{A}}$-projection from V to V_l. It is easy to check that

$$a\big[(\Pi_i - \Pi_{i-1})v, (\Pi_j - \Pi_{j-1})v\big] = 0, \quad \forall i \neq j.$$

We can easily obtain that

$$\begin{aligned} \|v\|_{\mathcal{A}}^2 &= \Big\| \sum_{l=0}^{L} (\Pi_l - \Pi_{l-1})v \Big\|_{\mathcal{A}}^2 = \sum_{0 \leqslant i,j \leqslant L} a\big[(\Pi_i - \Pi_{i-1})v, (\Pi_j - \Pi_{j-1})v\big] \\ &= \sum_{l=0}^{L} a\big[(\Pi_l - \Pi_{l-1})v, (\Pi_l - \Pi_{l-1})v\big] = \sum_{l=0}^{L} \Big| (\Pi_l - \Pi_{l-1})v \Big|_1^2. \end{aligned}$$

Motivated by the above norm equivalence and (5.10), we can construct a multilevel PSC method

$$\mathcal{B} = \sum_{j=1}^{J} \mathcal{S}_j \Pi_j.$$

However, Π_j is not good for computation in general. In the next section, we explore the idea of telescope expansion using L^2-projection instead of interpolation or Ritz-projection.

5.3 BPX preconditioner

Using the same multilevel space decomposition setting in the previous section, we can give parallel (PSC) versions of multilevel subspace correction method. The most prominent (multilevel) example of PSC methods is the BPX preconditioner [14] based on the multilevel subspace decomposition (5.2):

$$\mathcal{B} = \sum_{j=1}^{J} \mathcal{S}_j \mathcal{Q}_j, \tag{5.11}$$ eqn:Multileve

which is computationally more appealing and converges uniformly; see §3.2 for details. The HB and BPX preconditioners belong to the class of multilevel nodal basis preconditioners.

Telescope expansion of L^2-projections

Using notations in Definition 4.1, we have

$$\begin{cases} \mathcal{A}_l : V_l \mapsto V_l & (\mathcal{A}_l u_l, v_l) = a[u_l, v_l], \quad \forall\, u_l, v_l \in V_l; \\ \mathcal{Q}_l : L^2 \mapsto V_l & (\mathcal{Q}_l u, v_l) = (u, v_l), \quad \forall\, v_l \in V_l; \\ \Pi_l : \mathscr{V} \mapsto V_l & (\Pi_l u, v_l) = a[u, v_l], \quad \forall\, v_l \in V_l. \end{cases} \tag{5.12}$$

We introduce a new notation $i \wedge j := \min(i, j)$. It is trivial to see

$$\mathcal{Q}_i \mathcal{Q}_j = \mathcal{Q}_{i\wedge j}, \quad \Pi_i \Pi_j = \Pi_{i\wedge j}, \tag{5.13}$$ eqn:proj1

and

$$(\mathcal{Q}_i - \mathcal{Q}_{i-1})(\mathcal{Q}_j - \mathcal{Q}_{j-1}) = (\Pi_i - \Pi_{i-1})(\Pi_j - \Pi_{j-1}) = 0, \quad \forall\, i \neq j. \tag{5.14}$$ eqn:proj2

If we define $\mathcal{Q}_{-1} = \Pi_{-1} = 0$, we can give space decomposition

$$v = \sum_{l=0}^{L} (\mathcal{Q}_l - \mathcal{Q}_{l-1}) v = \sum_{l=0}^{L} (\Pi_l - \Pi_{l-1}) v. \tag{5.15}$$ eqn:decomp2

Norm equivalence

lem:telescope **Lemma 5.5** (Norm equivalence for L^2-projection telescope sum)**.** *For all $v \in V$, we have*

$$\left|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\right|_1 \cong h_l^{-1}\left\|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\right\|_0.$$

Proof. Using the inverse inequality, Proposition 3.3, we get

$$\left|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\right|_1 \lesssim h_l^{-1}\left\|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\right\|_0.$$

Proposition 3.5, together with the trivial equality

$$(\mathcal{Q}_l - \mathcal{Q}_{l-1})v = (\mathcal{I} - \mathcal{Q}_{l-1})(\mathcal{Q}_l - \mathcal{Q}_{l-1})v,$$

gives the other direction. □

lem:SCS4 **Lemma 5.6** (Strengthened Cauchy-Schwarz inequality for L^2-projection)**.** *If $u, v \in V$, let $u_i := (\mathcal{Q}_i - \mathcal{Q}_{i-1})u$, and $v_j := (\mathcal{Q}_j - \mathcal{Q}_{j-1})v$, then we have*

$$a[u_i, v_j] \lesssim \gamma^{|i-j|}\|u_i\|_{\mathcal{A}}\|v_j\|_{\mathcal{A}}.$$

Proof. If $j \geqslant i$, Lemma 5.5 shows $\|v_j\|_0 \lesssim h_j\|v_j\|_{\mathcal{A}}$ and the result follows directly from Lemma 5.2. If $i \geqslant j$, we can argue in a similar way. □

em:norm-equiv **Lemma 5.7** (Norm equivalences)**.** *For all $v \in V$, we have*

$$\sum_{l=0}^{L}\left\|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\right\|_1^2 \cong \|v\|_1^2.$$

Proof. Due to that $\Pi_l : \mathscr{V} \mapsto V_l$ is the standard H^1-projection, the finite element approximation theory gives

$$\left\|v - \Pi_l v\right\|_{1-\alpha} \lesssim h_l^{\alpha}\left\|v\right\|_1, \quad \forall\, v \in \mathscr{V}. \tag{5.16}$$ eqn:GalerkinE

Since $\mathcal{Q}_l$ is a L^2-projection, we have $\|\mathcal{Q}_l v\|_0 \leqslant \|v\|_0$, $\forall\, v \in L^2(\Omega)$. Furthermore, using Proposition 3.5, we obtain

$$\|\mathcal{Q}_l v\|_1 \leqslant \|v\|_1, \quad \forall\, v \in \mathscr{V}.$$

By space interpolation, we have, for any $\sigma \in (0, \frac{1}{2})$, that

$$\|\mathcal{Q}_l v\|_\sigma \leqslant \|v\|_\sigma, \quad \forall\, v \in \mathscr{V}.$$

Let $v_i := (\Pi_i - \Pi_{i-1})v$. Note that $\rho_l = \rho(\mathcal{A}_l) \cong h_l^{-2}$. It is easy to show, with help from Proposition 3.3 and (5.16), that

$$\left\|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v_i\right\|_1^2 \lesssim \rho_l^{\alpha}\left\|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v_i\right\|_{1-\alpha}^2 \lesssim \rho_l^{\alpha}\left\|v_i\right\|_{1-\alpha}^2 \lesssim \rho_l^{\alpha}h_i^{2\alpha}\left\|v_i\right\|_1^2.$$

Using this inequality and the Cauchy-Schwarz inequality, we can derive that

$$\sum_l \sum_{i,j} \big(\nabla(\mathcal{Q}_l - \mathcal{Q}_{l-1})v_i, \nabla(\mathcal{Q}_l - \mathcal{Q}_{l-1})v_j\big) = \sum_{i,j} \sum_{l=1}^{i\wedge j} \big(\nabla(\mathcal{Q}_l - \mathcal{Q}_{l-1})v_i, \nabla(\mathcal{Q}_l - \mathcal{Q}_{l-1})v_j\big)$$
$$\lesssim \sum_{i,j} \sum_{l=1}^{i\wedge j} \rho_l^\alpha h_i^\alpha h_j^\alpha \|v_i\|_1 \|v_j\|_1 \lesssim \sum_{i,j} \rho_{i\wedge j}^\alpha h_i^\alpha h_j^\alpha \|v_i\|_1 \|v_j\|_1 \lesssim \sum_{i,j} \gamma^{\alpha|i-j|} \|v_i\|_1 \|v_j\|_1.$$

We have shown that $\sum_{i,j} \gamma^{\alpha|i-j|} \|v_i\|_1 \|v_j\|_1 \lesssim \sum_i \|v_i\|_1^2 = \|v\|_1^2$, which shows

$$\sum_l \big\|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\big\|_1^2 \lesssim \|v\|_1^2.$$

On the other hand, using Lemma 5.6, we obtain

$$\begin{aligned} |v|_1^2 &= \sum_{i,j} \big(\nabla(\mathcal{Q}_i - \mathcal{Q}_{i-1})v, \nabla(\mathcal{Q}_j - \mathcal{Q}_{j-1})v\big) \\ &\lesssim \sum_{i,j} \gamma^{|i-j|} \big\|(\mathcal{Q}_i - \mathcal{Q}_{i-1})v\big\|_1 \big\|(\mathcal{Q}_j - \mathcal{Q}_{j-1})v\big\|_1 \lesssim \sum_i \big\|(\mathcal{Q}_i - \mathcal{Q}_{i-1})v\big\|_1^2. \end{aligned}$$

Hence we get the norm equivalence using Proposition 1.3. □

Remark 5.3 (Fractional norm)**.** We have shown the norm equivalence in H^1-norm. In fact, similar results also hold for $H^\alpha(\Omega)$ with $\frac{1}{2} < \alpha < \frac{3}{2}$.

BPX preconditioner and its convergence

All subspace problems are one-dimensional and, thus, very easy to solve. We can write the subspace solver (exact solver on each one-dimensional subspace) as follows:

$$\mathcal{S}_l^0 v := \sum_{i=1}^{n_l} (\mathcal{A}\phi_{l,i}, \phi_{l,i})^{-1} (v, \phi_{l,i})\, \phi_{l,i} = \sum_{i=1}^{n_l} (\nabla\phi_{l,i}, \nabla\phi_{l,i})^{-1} (v, \phi_{l,i})\, \phi_{l,i}.$$

Since we are now considering the uniform refinement for the linear finite element discretization, we can use an approximation of $\mathcal{S}_l^0$, for example a local relaxation method:

$$\mathcal{S}_l v := \sum_{i=1}^{n_l} h_l^{2-d} (v, \phi_{l,i})\, \phi_{l,i} \quad (\approx \mathcal{S}_l^0 v).$$

This simplification helps to reduce the cost of computation and implementation. Apparently, we have

$$(\mathcal{S}_l v, v) = h_l^{2-d} (\vec{v}, \vec{v}) = h_l^2 (v, v).$$

We have seen that the Richardson method, the damped Jacobi method, and the G-S method all satisfy such a condition; see (3.22).

onRequirement **Remark 5.4** (Behavior of the smoother)**.** Note that this "new" method is just the Richardson method with a weight $\omega = h_l^{2-d}$ on level l.

Using the above space decomposition and subspace solvers $\mathcal{S}_l$, the PSC method yields the well-known *BPX preconditioner*

$$\mathcal{B} = \sum_{l=0}^{L} \mathcal{S}_l \mathcal{Q}_l = \sum_{l=0}^{L} \mathcal{I}_l \mathcal{S}_l \mathcal{Q}_l = \sum_{l=0}^{L} \mathcal{I}_l \mathcal{S}_l \mathcal{I}_l^T \tag{5.17}$$ eqn:BPX

in operator form [14].

thm:BPX **Theorem 5.2** (Uniform convergence of BPX)**.** *The BPX preconditioner* (5.17) *is uniformly convergent, i.e.,* $\kappa(\mathcal{B}\mathcal{A}) \lesssim 1$.

Proof. We take a decomposition $v = \sum_{l=0}^{L} v_l := \sum_{l=0}^{L} (\mathcal{Q}_l - \mathcal{Q}_{l-1})v$, where $\mathcal{Q}_{-1} = 0$. Then we can obtain, from Lemmas 5.7 and 5.5, that

$$(\mathcal{A}v, v) \cong \sum_{l=0}^{L} \left|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\right|_1^2 \cong \sum_{l=0}^{L} h_l^{-2} \|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\|_0^2 = \big(\sum_{l=0}^{L} h_l^{-2}(\mathcal{Q}_l - \mathcal{Q}_{l-1})v, v\big).$$

Define $\tilde{\mathcal{A}} := \sum_{l=0}^{L} h_l^{-2}(\mathcal{Q}_l - \mathcal{Q}_{l-1})$. Apparently, $(\mathcal{A}v, v) \cong (\tilde{\mathcal{A}}v, v)$, $\forall v \in V$. Using (5.13) and (5.14), we can easily verify that

$$\tilde{\mathcal{A}}^{-1} = \sum_{l=0}^{L} h_l^2 (\mathcal{Q}_l - \mathcal{Q}_{l-1}).$$

Hence

$$(\tilde{\mathcal{A}}^{-1}v, v) = \sum_{l=0}^{L} h_l^2 (\mathcal{Q}_l v, v) - \sum_{l=0}^{L} h_l^2 (\mathcal{Q}_{l-1} v, v) = h_L^2 (\mathcal{Q}_L v, v) + \sum_{l=0}^{L-1} (1-\gamma^2) h_l^2 (\mathcal{Q}_l v, v).$$

Namely, $(\tilde{\mathcal{A}}^{-1}v, v) \cong (\mathcal{B}v, v) \Longrightarrow (\mathcal{A}v, v) \cong (\tilde{\mathcal{A}}v, v) \cong (\mathcal{B}^{-1}v, v)$. Then Lemma 2.5 gives the uniform convergence result. □

Remark 5.5 (Multilevel decomposition according to frequencies)**.** From the above analysis, we find that, for any $v \in V$

$$\left|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\right|_1 \cong h_l^{-1} \left\|(\mathcal{Q}_l - \mathcal{Q}_{l-1})v\right\|_0 \quad \Longrightarrow \quad \|\nabla v_l\|_0 \sim \|h_l^{-1} v_l\|_0.$$

This fact draws close comparison with the Fourier expansion. That is to say $v = \sum_{l=0}^{L} v_l$ is a multilevel decomposition to different frequencies. Hence $\tilde{\mathcal{A}}$ can be viewed as a multi-resolution expansion of $\mathcal{A}$ and $\kappa(\tilde{\mathcal{A}}^{-1}\mathcal{A}) \lesssim 1$.

Matrix representation of BPX

Using the matrix representation notations introduced in §3.2 and §3.5, the equation (3.37) in particular, we immediately obtain the matrix representation of the BPX method:

$$\underline{\mathcal{B}}\,\underline{u} = \underline{\mathcal{B}u} = \sum_{l=0}^{L} \underline{\mathcal{I}_l}\,\underline{\mathcal{S}_l}\,\underline{\mathcal{Q}_l}\,\underline{u} = \sum_{l=0}^{L} P_l\,(h_l^{2-d} M_l)\,(M_l^{-1} P_l^T M)\,\underline{u} = \sum_{l=0}^{L} h_l^{2-d} P_l\,P_l^T M\,\underline{u}.$$

In view of (3.16), we get the matrix form of the BPX preconditioner

$$B := \underline{\mathcal{B}} M^{-1} = \sum_{l=0}^{L} h_l^{2-d} P_l\,P_l^T. \tag{5.18}$$ eqn:MatBPX

This is the matrix form of the BPX preconditioner when we implement it.

To improve efficiency, we can use prolongation between two consecutive levels to obtain P_l.

5.4 Homework problems

hw:ConvDDM **Problem 5.1.** Give the complete proof of the uniform convergence of the two-level domain decomposition method (Proposition 5.1). What will happen if we do not include the coarse-level correction (Remark 5.1)?

hw:BPX **Problem 5.2.** Implement the BPX preconditioner for the Poisson's equation on a uniform grid. You can choose your favorite discretization method.

Chapter 6

Multigrid Methods

ch:mg

Multigrid (MG) methods are a group of algorithms for solving differential equations using a hierarchy of discretizations. The method has been proposed initially by Fedorenko [31] for 2D finite difference systems from the Poisson's equation. Its main idea is to accelerate the convergence of a basic iterative method (known as relaxation) by a global correction from time to time, accomplished by solving a coarse problem approximately. The coarse problem is "similar" to the fine grid problem, while cheaper to solve. This recursive process is repeated until a grid is reached where the cost of direct solution is negligible compared to the cost of one relaxation sweep on the finest grid. Later on Brandt [16] noticed that this method was considerably faster than standard relaxation methods and brought it to attention of the western scientific community.

6.1 Geometric multigrid method

sec:GMG

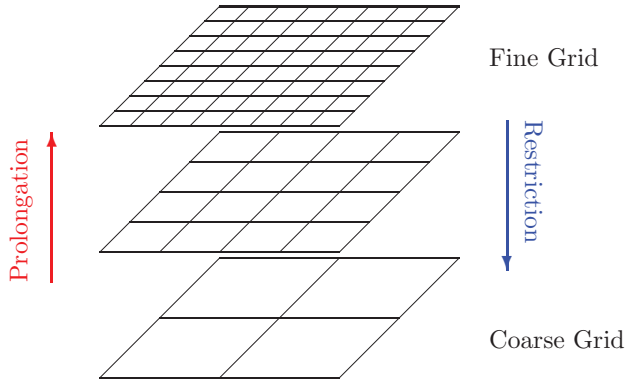
Geometric multigrid (GMG) method is an optimal iterative solver for linear algebraic system (2.1) arising from discretizations of elliptic partial differential equations such as the Poisson's equation. It is based on two important observations we made earlier in Chapter 3:

- A local relaxation method damps out non-smooth (high-frequency) components of the error and the residual becomes a relatively smooth vector after a few relaxation sweeps;
- A smooth (low-frequency) vector can be well approximated on coarse grids.

MG establishes and makes use of hierarchical structures. It is a good example of the idea of *divide and conquer*. MG can be used as a preconditioner for specific problems. This idea has been applied in two-grid methods; see §3.4. However, for large-scale problems in scientific and engineering computation, the coarse grid problem might be still too large to be solved quickly.

This makes introducing multiple grids a natural idea. The key steps in the multigrid method (see Figure 6.1) are as follows:

- **Relaxation/Smoothing**: Reduce high-frequency errors using one or more smoothing steps based on a simple iterative method;
- **Restriction**: Restrict the residual on a finer grid to a coarser grid;
- **Coarse correction**: Solve an approximate problem on a coarse grid;
- **Prolongation**: Represent the correction computed on a coarser grid to a finer grid.



fig_multigrid

Figure 6.1: Pictorial representation of a multigrid method with three grid levels.

V-cycle multigrid method

Now we will explain the multigrid algorithms using the $\mathcal{P}_1$ finite element method for the Poisson's equation on $\Omega \subset \mathbb{R}^d$ as an example. Suppose we have a sequence of meshes $\mathcal{M}_l$ $(l = 0, \ldots, L)$ generated from an initial mesh $\mathcal{M}_0$ by (uniform) regular refinements. Hence meshsize h_l of $\mathcal{M}_l$ is proportional to γ^{2l} with $\gamma \in (0, 1)$. Clearly,

$$h_0 > h_1 > h_2 > \cdots > h_L =: h.$$

It is easy to see that a multigrid method can be viewed a recursive two-grid method. So we only need to introduce how to do the iteration on two consecutive levels. We denote $\mathcal{I}_{l-1,l} : V_{l-1} \mapsto V_l$ $(l = 1, \cdots, L)$ as the natural embedding and $\mathcal{Q}_{l,l-1} = \mathcal{I}_{l-1,l}^T : V_l \mapsto V_{l-1}$ as the $(\cdot, \cdot)$-projection. Define $\mathcal{A}_l$ $(l = 1, \cdots, L)$ as the operator form of $\mathcal{A}$ on the subspace V_l in (5.1). Then a V-cycle multigrid method is given as follows:

alg:V-cycle **Algorithm 6.1** (One iteration of MG V-cycle)**.** Assume that $\mathcal{B}_{l-1} : V_{l-1} \mapsto V_{l-1}$ is defined and the coarsest level solver $\mathcal{B}_0 = \mathcal{A}_0^{-1}$ is exact. We shall recursively define $\mathcal{B}_l : V_l \mapsto V_l$ which is an iterator for the equation $\mathcal{A}_l v_l = r_l$. Let v_l be the initial guess on each level, i.e., $v_L = u^{(0)}$ and $v_l = 0$ for $0 < l < L$. Do the following steps:

(1) **Pre-smoothing**: For $k = 1, 2, \ldots, m$, compute

$$v_l \leftarrow v_l + \mathcal{S}_l\big(r_l - \mathcal{A}_l v_l\big);$$

(2) **Coarse-grid correction**: Find an approximate solution $e_{l-1} \in V_{l-1}$ of the residual equation on level $l-1$, i.e., $\mathcal{A}_{l-1} e_{l-1} = \mathcal{Q}_{l,l-1}\big(r_l - \mathcal{A}_l v_l\big)$, by an iterative method:

$$e_{l-1} \leftarrow \mathcal{B}_{l-1}\mathcal{Q}_{l,l-1}(r_l - \mathcal{A}_l v_l), \quad v_l \leftarrow v_l + \mathcal{I}_{l-1,l} e_{l-1};$$

(3) **Post-smoothing**: For $k = 1, 2, \ldots, m$, compute

$$v_l \leftarrow v_l + \mathcal{S}_l^T\big(r_l - \mathcal{A}_l v_l\big).$$

From this algorithm, we can see this MG V-cycle method is just a generalization of Algorithm 3.2 (the abstract two-grid method). Clearly, this geometric multigrid method (with one G-S iteration as pre-smoothing and one backward G-S iteration as post-smoothing) is just a special successive subspace correction (SSC) method based on the following multilevel space decomposition

$$V = \sum_{j=1}^{J} \tilde{V}_j = \sum_{l=L:-1:1} \sum_{i=1:n_l} V_{l,i} \; + \; V_0 \; + \sum_{l=1:L} \sum_{i=n_l:-1:1} V_{l,i},$$

which is a modification of (5.2). Furthermore, on each one-dimensional subspace $\tilde{V}_j$, the subspace problem is solved exactly.

According to Lemma 3.6, the error transfer operator of MG V-cycle on the l-th level can be written as

$$\mathcal{E}_l := \mathcal{I} - \mathcal{B}_l\mathcal{A}_l = \big(\mathcal{I} - \mathcal{S}_l^T\mathcal{A}_l\big)\big(\mathcal{I} - \mathcal{B}_{l-1}\mathcal{A}_{l-1}\Pi_{l-1}\big)\big(\mathcal{I} - \mathcal{S}_l\mathcal{A}_l\big),$$

where Π_{l-1} is the Ritz-projection from V to V_{l-1}. By applying this operator recursively, we obtain the error transfer operator for the MG V-cycle:

$$\mathcal{E}_L = \mathcal{I} - \mathcal{B}_L\mathcal{A}_L\Pi_L = \big(\mathcal{I} - \mathcal{S}_L^T\mathcal{A}_L\big)\cdots\big(\mathcal{I} - \mathcal{S}_1^T\mathcal{A}_1\big)\big(\mathcal{I} - \Pi_0\big)\big(\mathcal{I} - \mathcal{S}_1\mathcal{A}_1\big)\cdots\big(\mathcal{I} - \mathcal{S}_L\mathcal{A}_L\big).$$

Matrix representation of GMG

Similar to the matrix representation of two-grid method discussed in §3.4, we can write the matrix representation of multigrid method. By definition, we have

$$(\mathcal{A}_l u_l, v_l) = (\mathcal{A} u_l, v_l), \quad \forall\, u_l, v_l \in V_l.$$

Hence,

$$(\mathcal{A}_l \mathcal{Q}_l u, \mathcal{Q}_l v) = (\mathcal{I}_l^T \mathcal{A} \mathcal{I}_l \mathcal{Q}_l u, \mathcal{Q}_l v) = (\mathcal{A} \mathcal{I}_l \mathcal{Q}_l u, \mathcal{I}_l \mathcal{Q}_l v), \quad \forall\, u, v \in V.$$

It is easy to see that

$$\mathcal{Q}_l^T \mathcal{A}_l \mathcal{Q}_l = \mathcal{Q}_l^T \mathcal{I}_l^T \mathcal{A} \mathcal{I}_l \mathcal{Q}_l \quad \Longrightarrow \quad \underline{\mathcal{A}_l} = \underline{\mathcal{I}_l^T \mathcal{A} \mathcal{I}_l} = \underline{\mathcal{I}_l^T}\, \underline{\mathcal{A}}\, \underline{\mathcal{I}_l}.$$

This and (3.37), in turn, give the inter-grid transformations:

$$\hat{\mathcal{A}}_l = M_l\, \underline{\mathcal{A}_l} = M_l\, \underline{\mathcal{I}_l^T}\, \underline{\mathcal{A}}\, \underline{\mathcal{I}_l} = M_l\, \underline{\mathcal{Q}_l}\, M^{-1} \hat{\mathcal{A}}\, \underline{\mathcal{I}_l} = \underline{\mathcal{I}_l}^T \hat{\mathcal{A}}\, \underline{\mathcal{I}_l}, \quad 0 \leqslant l < L.$$

Hence we get the dual matrix form of the coarse level operator

$$\hat{\mathcal{A}}_l = P_l^T \hat{\mathcal{A}}\, P_l, \quad 0 \leqslant l < L.$$

Convergence analysis of GMG method ⋆

Now we show the MG /-cycle (i.e., half of the MG V-cycle) method converges uniformly using the X-Z identity. For simplicity, we show the proof in 1D here. For general analysis in $\mathbb{R}^d$, see HW 6.1. The multidimensional cases and other MG methods can be analyzed in a similar way, but much more technically involved.

Assume the subspace problems are solved exactly, i.e., $\mathcal{S}_{l,i} = \mathcal{A}_{l,i}^{-1}$, for $i = 1, \ldots, n_l$ and $l = 0, \ldots, L$. We denote the operators $\mathcal{J}_l : V \mapsto V_l$, i.e., for any function $v \in V$,

$$\big(\mathcal{J}_l v\big)(x) = \sum_{i=1}^{n_l} v(x_i^l)\, \phi_i^l(x), \quad \forall\, l = 0, \ldots, L.$$

Let $\mathcal{J}_{-1} v := 0$, $v_0 := \mathcal{J}_0 v$, and $v_l := (\mathcal{J}_l - \mathcal{J}_{l-1}) v$, $l = 1, \ldots, L$. Using the interpolants in multilevel spaces, we can write

$$v = \mathcal{J}_L v = \sum_{l=0}^{L} \big(\mathcal{J}_l - \mathcal{J}_{l-1}\big) v = \sum_{l=0}^{L} v_l. \tag{6.1}$$ eqn:HBdecomp

We also have

$$v = \sum_{l=0}^{L} v_l = \sum_{l=0}^{L} \sum_{i=1}^{n_l} v(x_i^l)\, \phi_i^l(x) =: \sum_{l=0}^{L} \sum_{i=1}^{n_l} v_{l,i}.$$

It is easy to check that

$$(\mathcal{I}-\mathcal{J}_k)v=\sum_{l=k+1}^{L}v_l=\sum_{l=k+1}^{L}\sum_{i=1}^{n_l}v_{l,i}$$

To estimate the convergence rate, in view of Corollary 4.1, we only need to estimate the quantity:

$$c_1:=\sup_{|v|_1=1}\inf_{\sum_{l,i}v_{l,i}=v}\sum_{l=0}^{L}\sum_{i=1}^{n_l}\Big|\Pi_{l,i}\sum_{(k,j)\geqslant(l,i)}v_{k,j}\Big|_1^2.$$

We now define and estimate

$$c_1(v):=\sum_{l=0}^{L}\sum_{i=1}^{n_l}\Big|\Pi_{l,i}\Big(\sum_{j=i}^{n_l}v_{l,j}+\sum_{k=l+1}^{L}\sum_{j=1}^{n_l}v_{k,j}\Big)\Big|_1^2.$$

We use the same notations introduced in Chapter 4 for projections, $\Pi_{l,i}: V\mapsto V_{l,i}$ is the $(\cdot,\cdot)_{\mathcal{A}}$-projection. For one-dimensional problems, it is easy to see that $\Pi_l=\mathcal{J}_l$; see HW 6.2. This leads to the following identity

$$\Pi_{l,i}(\mathcal{I}-\mathcal{J}_l)=0,\quad \forall\, 1\leqslant i\leqslant n_l,\ 0\leqslant l\leqslant L.$$

Furthermore, we also have $\Pi_{l,i}(\sum_{j\geqslant i}v_{l,j})=\Pi_{l,i}\big(v_{l,i}+v_{l,i+1}\big)$. Using these properties, we have

$$\begin{aligned}c_1(v)&=\sum_{l=0}^{L}\sum_{i=1}^{n_l}\big|\Pi_{l,i}\big(v_{l,i}+v_{l,i+1}\big)+\Pi_{l,i}(\mathcal{I}-\mathcal{J}_l)v\big|_1^2=\sum_{l=0}^{L}\sum_{i=1}^{n_l}\big|\Pi_{l,i}\big(v_{l,i}+v_{l,i+1}\big)\big|_1^2\\ &\lesssim\sum_{l=0}^{L}\sum_{i=1}^{n_l}|v_{l,i}|_1^2=\sum_{l=0}^{L}h_l^{-2}\|(\mathcal{J}_l-\mathcal{J}_{l-1})v\|_0^2\lesssim\sum_{l=0}^{L}|v_l|_1^2=|v|_1^2.\end{aligned}$$

The last equality is easy to check; see HW 6.3. This estimate shows the convergence rate of MG is uniformly bounded.

rem:HB **Remark 6.1** (Relation with the HB preconditioner)**.** Note that several places in the above analysis depend on the one-dimensional ($d=1$) assumption. In fact, the decomposition (6.1) used in this proof is the hierarchical basis (HB) decomposition in §5.2 and convergence rate of the HB method is not optimal in multidimensional cases ($d>1$). So the proof must be changed in higher dimensions.

Problems with anisotropic coefficients ⋆

For GMG, smoothness of error is in the usual geometric sense. But it is not always easy to do so. A representative example is the second-order elliptic problem

$$-\epsilon u_{xx}-u_{yy}=f(x,y),\quad \forall\,(x,y)\in\Omega, \tag{6.2}$$ eqn:aniso

where $\epsilon > 0$ is usually small. Other examples include problems with high-contrast coefficients, problems on anisotropic meshes, etc.

If we apply the standard finite different discretization in §1.2 on uniform $n \times n$ tensor-product grid for this problem, or equivalently the $\mathcal{P}_1$ finite element discretization on uniform triangular grid from regular refinements, then the coefficient matrix for (6.2) is

$$A_\epsilon = I \otimes A_{1,\epsilon} + C \otimes I, \quad \text{with } A_{1,\epsilon} = \text{tridiag}(-\epsilon, 2+2\epsilon, -\epsilon),\ C = \text{tridiag}(-1, 0, -1).$$

The eigenvalues of A are given

$$\lambda_{i,j}(A_\epsilon) = 2(1+\epsilon) - 2\epsilon\cos\frac{i\pi}{n+1} - 2\cos\frac{j\pi}{n+1} = 4\epsilon\sin^2\frac{i\pi}{2(n+1)} + 4\sin^2\frac{j\pi}{2(n+1)},$$

with eigenvectors

$$\vec{\xi}_{i,j} = \left(\sin\frac{ki\pi}{n+1}\sin\frac{lj\pi}{n+1}\right)_{k,l=1,\ldots,n}.$$

If $\epsilon \ll 1$, then $\lambda_{1,1} < \lambda_{2,1} < \cdots < \lambda_{n,1} < \lambda_{1,2} < \lambda_{2,2} < \cdots$. We notice that, unlike the Poisson's equation, these eigenvalues are ordered in a different pattern. The geometric low frequencies can be highly oscillatory in x-direction. It is natural to expect such a behavior from the PDE itself as the x-direction is much less diffusive than the y-direction. We call the x-direction (with small coefficient) the weak direction and the y-direction the strong direction.

We can also view this problem from a different perspective. Using the analysis in §3.3, we obtain that the error of the G-S method satisfies

$$(2+2\epsilon)e_{i,j}^{\text{new}} = \epsilon e_{i-1,j}^{\text{new}} + \epsilon e_{i+1,j}^{\text{old}} + e_{i,j-1}^{\text{new}} + e_{i,j+1}^{\text{old}}, \quad i,j = 1,\ldots,n.$$

According to the local Fourier analysis, we can obtain that

$$\lambda(\theta) := \frac{\alpha_\theta^{\text{new}}}{\alpha_\theta^{\text{old}}} = \frac{\epsilon e^{\sqrt{-1}\theta_1} + e^{\sqrt{-1}\theta_2}}{2 + 2\epsilon - \epsilon e^{-\sqrt{-1}\theta_1} - e^{-\sqrt{-1}\theta_2}}.$$

In this case, the smoothing factor of the G-S method is

$$\bar{\rho}_{\text{GS}} = \lambda\left(\frac{\pi}{2}, \arctan\left(\frac{\epsilon(1-\bar{\rho}_{\text{GS}}^2)}{2(\epsilon+1)\bar{\rho}_{\text{GS}}^2}\right)\right) = \frac{\sqrt{5\epsilon^2 - 2\epsilon + 1} + 2}{5\epsilon + 3} \longrightarrow 1, \quad \text{as } \epsilon \to 0.$$

This means the standard G-S method does not have smoothing effect on the anisotropic problem when ϵ is small.

On the other hand, if we apply the line G-S smoother, things will be a lot different. Suppose we apply the line smoother in natural ordering (from left to right), namely,

$$(2+2\epsilon)u_{i,j}^{\text{new}} = \epsilon u_{i-1,j}^{\text{new}} + \epsilon u_{i+1,j}^{\text{old}} + u_{i,j-1}^{\text{new}} + u_{i,j+1}^{\text{new}}, \quad j = 1,\ldots,n, \quad i = 1,\ldots,n.$$

Then the error satisfies

$$(2+2\epsilon)e_{i,j}^{\text{new}} = \epsilon e_{i-1,j}^{\text{new}} + \epsilon e_{i+1,j}^{\text{old}} + e_{i,j-1}^{\text{new}} + e_{i,j+1}^{\text{new}}, \quad j = 1,\ldots,n, \quad i = 1,\ldots,n.$$

And we get

$$\lambda(\theta) := \frac{\alpha_\theta^{\text{new}}}{\alpha_\theta^{\text{old}}} = \frac{\epsilon e^{\sqrt{-1}\theta_1}}{2 + 2\epsilon - \epsilon e^{-\sqrt{-1}\theta_1} - 2e^{-\sqrt{-1}\theta_2}}.$$

The maximal smoothing factor is then

$$\bar{\rho}_{\text{LGS}} = \max\left\{\frac{\epsilon}{2+\epsilon}, \frac{\sqrt{5}}{5}\right\}.$$

If $0 < \epsilon \leqslant 1$, we always have $\bar{\rho}_{\text{LGS}} = \sqrt{5}/5 < 1$ independent of ϵ.

In multigrid methods setting, one could handle such an equation using special techniques like: (1) an line smoother (group all y-variables corresponding to the same x-coordinate together), or (2) semi-coarsening (only coarse in y-direction), or (3) operator-dependent interpolations. Later in this chapter, we will turn our attention to the third approach, which leads to algebraic multigrid methods for solving such difficult problems.

General procedure of multigrid methods

From our discussions above, we observe that a typical MG algorithm contains two phases—the "setup" phase and the "solve" phase. The setup phase initializes a hierarchical structure, including coarse spaces, prolongations and restrictions, coarse problem solvers for multilevel iterations. Notice that the setup phase only needs to be called once before iterations; sometimes, the same setup phase can be used at different time levels for time-dependent problems. For geometric multigrid (GMG) methods, the setup phase is trivial using the hierarchical grid structure. However, GMG methods are difficult to apply for equations on general domains with unstructured grids. Algebraic multigrid (AMG) methods can be viewed as a generalization of geometric multigrid methods; see [60] and references therein for more details.

We now explain how to perform multigrid setup phase in a relatively general setting. Once the setup is done, an appropriate nested iteration scheme should be chosen for the solve phase; see §6.2. It is immediately clear that we only need to discuss how to setup hierarchical information in two consecutive grids/levels for multigrid methods. We can summarize a general multigrid setup procedure as the following steps:

Step 1. **Selecting a smoother:** Choose a smoother S for $Au = f$.

Step 2. **Coarsening:** Identify a coarser space $V_c \subset V$, which contains smooth vectors.

Step 3. **Constructing a prolongation:** Construct a prolongation P in two steps:

3a. Decide, for each fine variable, which coarse variables will be used for interpolation;

3b. Determine the weights for prolongation P.

Step 4. **Recursion or nested iteration:** Define $A_c = P^T AP$ and $f_c = P^T f$; then apply the algorithm for the coarse problem $A_c u_c = f_c$.

For GMG methods discussed by far, Steps 2–4 are determined by the information of nested grids and the users can only find an appropriate smoother S. For example, in §1.4, we have presented a 1D GMG method in a purely algebraic fashion. We have observed that:

(1) GMG coarsening explores the topology of the graph representing the stiffness matrices on different levels are explicitly clear from the geometric refinement procedure;

(2) Prolongation and restriction for GMG usually depend only on the topological structure of the graph without knowing the grid coordinates;

(3) For GMG, smoothness of error is in the geometric sense and, in the algebraic setting, smooth error can be geometrically non-smooth.

6.2 Nested iterations

sec:cycles

The solve phase approximates corresponding problems by calling a two-grid algorithm recursively. There are different approaches for the solve phase; for example, we have seen the V-cycle method in §6.1. In this section, we discuss a few popular methods for the solve phase.

V-cycle and its generalizations

The multigrid V-cycle iterator $\mathcal{B}$, Algorithm 6.1, is a two-grid method with an inexact coarse-level solver defined recursively, i.e., the coarse-level iterator $\mathcal{B}_c$ is just $\mathcal{B}$ on the coarse-grid. On the coarse level, we start from the initial guess $u_c^{\text{old}} = 0$ and then iterate

$$u_c^{\text{new}} = u_c^{\text{old}} + \mathcal{B}_c\big(f_c - \mathcal{A}_c u_c^{\text{old}}\big), \quad \text{where } \mathcal{B}_c \text{ is the two-grid method for } \mathcal{A}_c.$$

In the the V-cycle, we only apply the above iteration once on the coarse-level. Apparently, this procedure can be generalized. For example, we can iterate multiple steps:

$$u_c^{(0)} = 0, \quad u_c^{(k)} = u_c^{(k-1)} + \mathcal{B}_c\big(f_c - \mathcal{A}_c u_c^{(k-1)}\big), \quad k = 1, \ldots, \nu.$$

This gives an estimate

$$u_c^{(\nu)} = \mathcal{B}_c f_c + (\mathcal{I} - \mathcal{B}_c\mathcal{A}_c)u_c^{(\nu-1)} = \mathcal{B}_c f_c + \mathcal{E}_c u_c^{(\nu-1)} = \cdots = \big(\mathcal{I} + \mathcal{E}_c + \cdots + \mathcal{E}_c^{\nu-1}\big)\mathcal{B}_c f_c,$$

where $\mathcal{E}_c := \mathcal{I} - \mathcal{B}_c\mathcal{A}_c$. We can define a new iterator $\mathcal{B}_{c,\nu}$ such that

$$\mathcal{B}_{c,\nu} f_c := \big(\mathcal{I} - \mathcal{E}_c^{\nu}\big)\big(\mathcal{I} - \mathcal{E}_c\big)^{-1}\mathcal{B}_c f_c = \big(\mathcal{I} - \mathcal{E}_c^{\nu}\big)\mathcal{A}_c^{-1} f_c. \tag{6.3}$$ eqn:nu-cycle

Motivated by (6.3), we can introduce a polynomial $q_\nu(t) := (1-t)^\nu \in \mathcal{P}_\nu$ and let

$$\mathcal{B}_{c,\nu} := \Big(\mathcal{I} - q_\nu\big(\mathcal{B}_c\mathcal{A}_c\big)\Big)\mathcal{A}_c^{-1}.$$

Then $\nu = 1$ yields the V-cycle. And the first non-trivial example is the well-known W-cycle ($\nu = 2$), which is a simple extension of the V-cycle algorithm; see Figure 6.2. By iterate twice as in (6.4), we can obtain $\mathcal{B}_{c,2}$ (the W-cycle); see HW 6.4.

alg:W-cycle **Algorithm 6.2** (One iteration of MG W-cycle)**.** Assume that $\mathcal{B}_{l-1} : V_{l-1} \mapsto V_{l-1}$ is defined and the coarsest level solver $\mathcal{B}_0 = \mathcal{A}_0^{-1}$ is exact. We shall recursively define $\mathcal{B}_l : V_l \mapsto V_l$ which is an iterator for the equation $\mathcal{A}_l v = r_l$. Let $v = v^{(0)}$ be the initial guess.

(1) **Pre-smoothing**: For $k = 1, 2, \ldots, m$, compute

$$v \leftarrow v + \mathcal{S}_l\big(r_l - \mathcal{A}_l v\big);$$

(2) **Coarse-grid correction**: Find an approximate solution $e_{l-1} \in V_{l-1}$ of the residual equation on level $l-1$, i.e., $\mathcal{A}_{l-1}e_{l-1} = \mathcal{Q}_{l,l-1}\big(r_l - \mathcal{A}_l v\big)$ using the iteration: Let $e_{l-1} = 0$ initially. For $k = 1, 2$,

$$e_{l-1} \leftarrow e_{l-1} + \mathcal{B}_{l-1}\Big(\mathcal{Q}_{l,l-1}(r_l - \mathcal{A}_l v) - \mathcal{A}_{l-1}e_{l-1}\Big); \tag{6.4}$$ eqn:W-cycle

Update the solution with

$$v \leftarrow v + \mathcal{I}_{l-1,l}e_{l-1};$$

(3) **Post-smoothing**: For $k = 1, 2, \ldots, m$,

$$v \leftarrow v + \mathcal{S}_l^T\big(r_l - \mathcal{A}_l v\big).$$

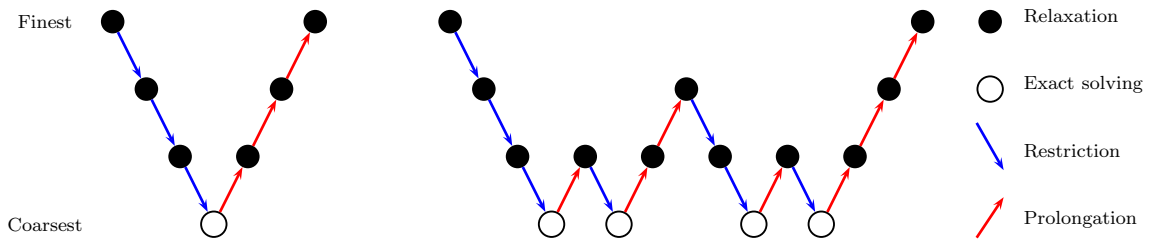


fig:VW-cycle Figure 6.2: Multigrid V-cycle (left) and W-cycle (right).

In V-cycle and W-cycle, the iterators on all the coarser levels are the same. We can also use variable polynomial orders ν_l on level l $(0 < l < L)$. From the above discussion, we also have a lot of freedom in the choice of $q_\nu(t)$. For example, we can use a polynomial $q_\nu(t)$ such that $q_\nu(0) = 1$ and $0 \leqslant q_\nu(t) < 1$ on the spectrum of $\mathcal{B}_c\mathcal{A}_c$. This type of methods are referred to as the AMLI-cycle (Algebraic Multi-Level Iteration cycle[1]); see [1] and references therein for details.

Example 6.1 (A simple AMLI-cycle)**.** A simple AMLI-cycle method is to give $l_0 \geqslant 1$, $\mu_1 \geqslant \mu_2 \geqslant 1$, and use the following polynomial orders

$$\nu_l := \begin{cases} \mu_1, & \text{if } \ l = kl_0; \\ \mu_2, & \text{otherwise.} \end{cases}$$

It is clear that, if $l_0 = \mu_1 = \mu_2 = 1$, then this method is just the V-cycle. □

Now we turn our attention to the work estimate of nested cycles. For simplicity, we consider the AMLI-cycle with $\mu_2 \equiv 1$ only. Denote the work needed by $\mathcal{B}_l$ is W_l. Assume the each smoothing sweep costs $O(N_l)$ operations and $N_l \sim h_l^{-d} \sim \gamma^{-2ld}$. Then it requires $2m$ $O(N_l)$ operations for the pre- and post-smoothing on level l. The prolongation and restriction also requires $O(N_l)$ operations. Hence, for the AMLI-cycle, we have

$$\begin{aligned} W_{(k+1)l_0} &= O\big(N_{kl_0+1} + \cdots + N_{kl_0+l_0}\big) + \mu_1 O(N_{kl_0}) + \mu_1 W_{kl_0} \\ &= O\big(N_{(k+1)l_0}\big) + \mu_1 O(N_{kl_0}) + \mu_1 W_{kl_0} \\ &= O\big(N_{(k+1)l_0}\big) + 2\mu_1 O(N_{kl_0}) + \mu_1^2 W_{(k-1)l_0} \\ &= \cdots\cdots \\ &= \mu_1^k W_{l_0} + O\Big(\sum_{j=2}^{k+1} \mu_1^{k+1-j} N_{jl_0}\Big) \\ &= O\Big(\sum_{j=1}^{k+1} \mu_1^{k+1-j} N_{jl_0}\Big) \\ &= O\big(N_{(k+1)l_0}\big) \sum_{j=0}^{k} \big(\mu_1 \gamma^{2dl_0}\big)^j. \end{aligned}$$

Hence, this AMLI method costs $O(N)$ operations in each cycle, if $\mu_1\gamma^{2dl_0} < 1$. Here $N = N_L$ is the number of unknowns on the finest grid. This analysis also yields the complexity of the standard V-cycle method quickly.

Indeed, we can choose some optimal polynomial $q_\nu(t)$ like the Chebyshev polynomials. This reminds us about the Krylov subspace methods discussed in §2.2. Inspired by this similarity, we

[1] Here "algebraic" stands for the fact that certain inner polynomial iterations are used in the definition of the multilevel cycle.

can apply a preconditioned Krylov methods (like Flexible CG or GCR methods) on some of the coarse levels to improve convergence. This type of methods are called Krylov-cycle (K-cycle) methods or Nonlinear AMLI methods [49].

The previous multigrid methods converge uniformly with respect to the meshsize h and requires $O(N)$ operations in each cycle. This means the computation cost is $O(N)$ to reach a fixed tolerance. On the other hand, when we solve a discrete partial differential equation, we usually need to use smaller tolerances for finer meshsizes. This means, in order to reach the discretization accuracy, the V-cycle multigrid method requires $O(N \log N)$ operations.

Full multigrid method

One way to improve the multigrid method is to provide good initial guesses using coarse approximation (cheap in computation). This idea leads to the so-called full multigrid (FMG) cycle; see Figure 6.3. From this figure, we can see the full multigrid method is a sequence of V-cycles on

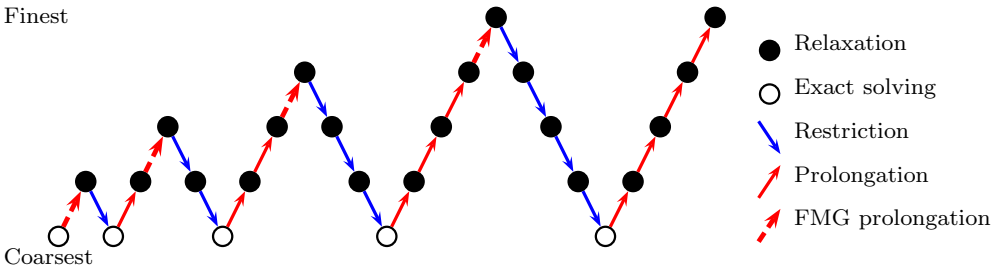


fig:F-cycle

Figure 6.3: Full multigrid cycle.

different levels. We can write the concrete algorithm as follows:

Listing 6.1: Full multigrid method

```
1  ũ_0 ← A_0^{-1} f_0;
2  for l = 1, ..., L
3      u_l^{(0)} ← I_{l-1,l} ũ_{l-1};
4      u_l^{(k)} ← V-cycle(l, f_l, u_l^{(k-1)}),   k = 1, ..., ν;
5      ũ_l ← u_l^{(ν)};
6  end
```

thm:FMG

Theorem 6.1 (Full multigrid convergence)**.** *Assume that the l-th level iteration is a contraction with a contraction factor $0 < \delta < 1$ independent of level l. If ν is large enough, then we have*

$$\|\!|u_l - \tilde{u}_l|\!\| \lesssim h_l |u|_{2,\Omega},$$

where u_l is the exact solution of finite element problem on level l and $\tilde{u}_l$ is the approximation solution on l-th level.

Proof. Let $e_l := u_l - \tilde{u}_l$. Apparently, on the coarsest level, we have $e_0 = 0$. On the l-th level $(0 < l \leqslant L)$, we have

$$\begin{aligned} \|\!| e_l |\!\| &\leqslant \delta^\nu \|\!| u_l - \tilde{u}_{l-1} |\!\| \leqslant \delta^\nu \Big(\|\!| u_l - u |\!\| + \|\!| u_{l-1} - u |\!\| + \|\!| u_{l-1} - \tilde{u}_{l-1} |\!\| \Big) \\ &\leqslant \delta^\nu \Big(C h_l |u|_{2,\Omega} + \|\!| e_{l-1} |\!\| \Big). \end{aligned}$$

By iteration, if $\delta^\nu < \gamma^2$, we obtain that

$$\|\!| e_l |\!\| \leqslant C \Big(\delta^\nu h_l + \delta^{2\nu} h_{l-1} + \cdots + \delta^{l\nu} h_1 \Big) |u|_{2,\Omega} \leqslant \frac{C\delta^\nu h_l}{1 - \gamma^{-2}\delta^\nu} |u|_{2,\Omega} \lesssim h_l |u|_{2,\Omega}.$$

Hence the result. □

The theorem indicates that, if we do enough V-cycles on each level, we can obtain an approximate solution within the accuracy of discretization error. This result shows that FMG can reach discretization error tolerance using $O(N)$ operations.

6.3 From geometric to algebraic multigrid

sec:AMG

Consider the system of equations arising form the Poisson's equation on unstructured meshes or the second-order elliptic equation with anisotropic coefficients

$$Au = f, \quad \text{where } A \in \mathbb{R}^{N \times N} \text{ and } u, f \in \mathbb{R}^N.$$

Problems with anisotropic coefficients on regular meshes, or problems with isotropic coefficients but on anisotropic meshes, will cause troubles for geometric multigrid methods. While geometric multigrid essentially relies on the availability of robust smoothers, AMG takes a different approach. AMG assumes a simple relaxation process to be given (typically point-wise relaxations) and then attempts to construct a suitable operator-dependent interpolation using the algebraic information of the coefficient matrix A. In §6.1, we have already found that the G-S method is not a good smoother in the usual geometric sense for (6.2) if the coefficient ϵ is small. However, we can choose coarse spaces adapted to such a smoother by mimicking the idea of semi-coarsening.

Sparse matrices and graphs ⋆

A sparse matrix can be represented as a graph. As the sparse matrices that we consider are mainly symmetric in the following we only discuss undirected graphs here. We first introduce a few elementary concepts from the graph theory. An *undirected graph* (or simply a *graph*) G is a pair (V, E), where V is a finite set of points called *vertices* and E is a finite set of *edges*. As set

of vertices we always consider subsets of $\{1, \ldots, N\}$. An edge in E is an unordered pair $(j,\, k)$ with $j,\, k \in \mathsf{V}$. A graph $\mathsf{G}_0 = (\mathsf{V}_0, \mathsf{E}_0)$ is called a *subgraph* of $\mathsf{G} = (\mathsf{V}, \mathsf{E})$, if $\mathsf{V}_0 \subset \mathsf{V}$ and $\mathsf{E}_0 \subset \mathsf{E}$.

If $(j, k) \in \mathsf{E}$ is an edge in an undirected graph $\mathsf{G} = (\mathsf{V},\, \mathsf{E})$, vertices j and k are said to be *adjacent*. The set of neighboring vertices of i is the set of all vertices that are adjacent to i; and it is denoted as $\mathsf{N}_i \subseteq \mathsf{V}$. A *path* from a vertex i to another vertex j is a sequence of edges $\big\{(i,\, j_1),\, (j_1, j_2),\, \ldots,\, (j_{l-2}, j_{l-1}),\, (j_{l-1}, j)\big\} \subseteq \mathsf{E}$ and the number of edges l is called the length of this path. A vertex j is *connected* to a vertex k if there is a path from j to k. The distance between j and k is defined as the length of the shortest path between these two vertices. Apparently, the distance between two vertices is equal to 1 if they are adjacent and is set to ∞ if they are not connected.

An undirected graph $\mathsf{G} = (\mathsf{V},\, \mathsf{E})$ is *connected* if any pair of vertices are connected by a path, otherwise G is said to be *disconnected*. An *independent set* of V is a set of vertices of G, no two of which are adjacent. A *maximal independent set* (MIS) is an independent set such that adding any other vertex to the set will introduce at least one adjacent pair.

Let $A \in \mathbb{R}^{N \times N}$ be a sparse matrix. The *adjacency graph* of A, denoted by $\mathsf{G}(A)$, is a graph $\mathsf{G} = (\mathsf{V},\, \mathsf{E})$ with $\mathsf{V} := \{1, 2, \ldots, N\}$ and

$$\mathsf{E} := \big\{(j,\, k) \,:\, a_{j,k} \neq 0\big\}.$$

As a general rule, sparse matrices do not provide any geometric information for the underlying graph and only the combinatorial/topological properties of $\mathsf{G}(A)$ or its subgraphs; see Figure 6.4. We note that two different discretizations on different meshes could lead to same sparse coefficient matrix A and, in turn, same graph $\mathsf{G}(A)$.

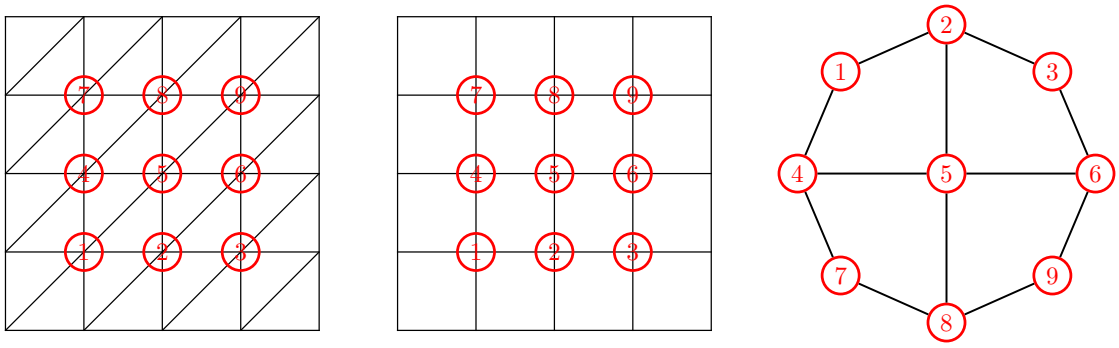


Figure 6.4: Finite element grid (left), difference grid (middle), and graph of their corresponding stiffness matrices (right).

fig:graph

Let A be the coefficient matrix corresponding to the finite element discretization of the second-order elliptic equation with Neumann boundary condition. Apparently A has zero row

sum. Hence we can write

$$(Au, v) = \sum_{\substack{(i,j)\in \mathsf{E} \\ i<j}} -a_{i,j}(u_i - u_j)(v_i - v_j). \tag{6.5}$$ eqn:Auv-Neuma

On the other hand, we can easily derive the corresponding equality for the Dirichlet boundary condition or the mixed boundary condition:

$$(Au, v) = \sum_{\substack{(i,j)\in \mathsf{E} \\ i<j}} -a_{i,j}(u_i - u_j)(v_i - v_j) \qquad \text{and} \qquad u_j = v_j = 0, \ \forall\, x_j \in \Gamma_D. \tag{6.6}$$ eqn:Auv-Diric

M-matrix and Delaunay triangulation ⋆

We first introduce the concept of M-matrix. We call A an M-matrix if it is irreducible (i.e., the graph $\mathsf{G}(A)$ is connected) and

$$a_{i,i} > 0, \quad a_{i,j} \leqslant 0 \ \ (i \neq j), \quad a_{j,j} \geqslant \sum_{i\neq j} |a_{i,j}|, \quad a_{j,j} > \sum_{i\neq j} |a_{i,j}| \ \text{ for at least one } j.$$

Apparently, the stiffness matrix in (1.28) is an M-matrix. It is in general not the case for the stiffness matrices from finite element discretizations. In fact, whether a stiffness matrix is an M-matrix depends on the underlying mesh $\mathcal{M}$.

First we introduce a few notations using Figure 6.5. In any give simplicial element τ in $\mathbb{R}^3$; similar definitions can be introduced in $\mathbb{R}^d$ for $d \geqslant 2$. An edge (i, j) has two vertices x_i and x_j and denote this edge as E. Let $\kappa_E(\tau) := F_i \bigcap F_j$ and $\theta_E(\tau)$ be the angle between faces F_i and F_j. Define a quantity

$$\omega_E(\tau) := \frac{1}{d(d-1)} |\kappa_E(\tau)| \cot \theta_E(\tau). \tag{6.7}$$ eqn:OET

We then have the following result; see [57] for details.

prop:M-matrix **Proposition 6.1** (Condition for M-matrix)**.** *The stiffness matrix for the Poisson's equation is an M-matrix if and only if, for any edge E, $\sum_{\tau\supset E} \omega_E(\tau) \geqslant 0$ with $\omega_E(\tau)$ defined in (6.7).*

Remark 6.2 (Delaunay triangulation and M-matrix)**.** In $\mathbb{R}^2$, the above proposition simply means the sum of the angle opposite to any edge is less than or equal to π, which means the underlying triangulation must be Delaunay. Hence the stiffness matrix for the Poisson's equation is an M-matrix if the triangulation is Delaunay. And the condition is almost sharp[2].

For a given mesh $\mathcal{M}_h$, the stiffness matrix of $\mathcal{P}_1$-finite element method for the Poisson's equation is not necessarily an M-matrix. However, it can be estimated by an M-matrix. More

[2]The opposite direction is true with a few possible exceptions near the boundary

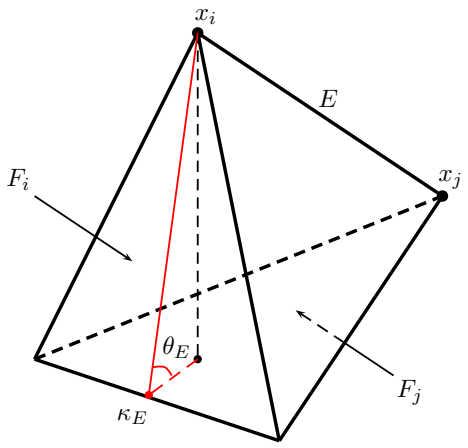


fig:simplex

Figure 6.5: Definition of θ_E and κ_E in a simplex in 3D.

specifically, if we keep all the vertices of $\mathcal{M}_h$ and swap internal edges, we can obtain a Delaunay triangulation $\mathcal{M}_h^D$. We have

$$\left(A_{\mathcal{M}_h^D} v, v\right) \leqslant \left(A_{\mathcal{M}_h} v, v\right), \quad \forall\, v \in \mathbb{R}^N;$$

moreover, the equality in the above inequality holds if and only if $\mathcal{M}_h$ is Delaunay. We refer the interested readers to [50] for details. Let $\phi_{\mathcal{M}_h} \in V_h$ is a piecewise linear function and $\phi_{\mathcal{M}_h}(x) = \sum_{i=1}^N v_i \phi_{i,\mathcal{M}_h}(x)$. Then we have

$$\left|\phi_{\mathcal{M}_h^D}\right|_1^2 \leqslant \left|\phi_{\mathcal{M}_h}\right|_1^2, \quad \forall\, v \in \mathbb{R}^N.$$

This means the Delaunay triangulation results in lower roughness of finite element functions among all possible triangulations on a fixed set of vertices.

Algebraic smooth error

Based on the two-level convergence theory discussed in §3.4 (Theorem 3.5, in particular), we have the following result: For a given smoother S, the best coarse space V_c^{opt} of dimension N_c is given by

$$V_c^{\text{opt}} = \text{span}\{\phi_k\}_{k=1}^{N_c}, \tag{6.8}$$

eqn:optVH

where $\{\phi_k\}_{k=1}^{N_c}$ are the eigenfunctions corresponding to the smallest eigenvalues $\lambda_k(\overline{S}A)$. However, it is difficult to find small eigenvalues of $\overline{S}A$ in practice.

Hence it is important to give a practical characterization of the smooth error. A good interpretation of smooth error in algebraic sense could lead to an efficient AMG method. In view of (3.24), we know that the standard point-wise relaxation methods, like the Richardson, weighted Jacobi, and Gauss–Seidel methods, satisfy that

$$\rho_A^{-1}\left(v, v\right)_A \lesssim \left(SAv, v\right)_A \leqslant \omega_1 \left(v, v\right)_A.$$

And Theorem 3.5 indicates the lower end of the spectrum of $\overline{S}A$ are low frequencies. This motives the following definition of the algebraic smooth vector:

Definition 6.1 (Algebraic smoothness)**.** Let $\varepsilon \in (0,1)$ be a small parameter. If $v \in V$ satisfies

$$\big(\overline{S}Av, v\big)_A \leqslant \varepsilon^2 \big(v, v\big)_A,$$

then v is called algebraically ε-smooth.

Since $\overline{S}$ is SPD, the algebraically smooth vectors satisfy

$$\|v\|_A^2 = \big(\overline{S}Av, \overline{S}^{-1}v\big) \leqslant \big(\overline{S}Av, Av\big)^{1/2} \big(\overline{S}^{-1}v, v\big)^{1/2} \leqslant \varepsilon \big(\overline{S}^{-1}v, v\big)^{1/2} \|v\|_A.$$

Then we can derive an estimate

$$\|v\|_A \leqslant \varepsilon \|v\|_{\overline{S}^{-1}}, \tag{6.9}$$ eqn:algsmooth

which can be also viewed as an alternative characterisation of algebraically smooth vectors. By adding and subtraction and (2.10), we have

$$\Big(\big(I - \overline{S}A\big)v, v\Big)_A \geqslant (1-\varepsilon^2)\big(v, v\big)_A \quad \Longrightarrow \quad \frac{\|(I-SA)v\|_A^2}{\|v\|_A^2} \geqslant 1 - \varepsilon^2.$$

Apparently, the contraction factor for this vector v is close 1 if ε is small. Basically, this means the algebraically smooth error components are those which the smoother cannot damp efficiently; that is to say, an error not eliminated by the smoother is a smooth error; see Remark 1.8 for geometric smooth error.

Remark 6.3 (Local adaptation of AMG)**.** In AMG methods, it is not important whether S smooths the error in any geometric sense or not. On the contrary, the key point is that the error after smoothing sweeps can be characterized algebraically to a degree which makes it possible to construct coarse levels and define interpolations which are locally adapted to the properties of the given smoother.

Remark 6.4 (Smooth error and the Classical AMG)**.** A simple characterization of smooth error is used in methods like the Classical AMG. If e corresponds to the low-end of eigenvalues, then we have $Ae \ll 1$. According to (6.5), this is equivalent to say that

$$(Ae, e) = \sum_{i<j} -a_{i,j}(e_i - e_j)^2 \ll 1. \tag{6.10}$$ eqn:CAMG-smoo

This inequality provides an important motivation for the Classical AMG: Smooth error varies slowly in the direction of relatively large (negative) coefficients of the matrix.

Construction of coarse spaces

From Theorem 3.3, the convergence rate of the two-grid method depends on effectiveness of the smoother S and approximability of the coarse space range(P). Now we discuss a few guidelines on how to construct coarse spaces and prolongation matrices.

In §6.1, we discussed a general procedure of multigrid setup phase. The coarsening algorithms are methods for determining the coarse level variables. Such algorithms are usually based on selecting or combining vertices in the adjacency graph corresponding to the (filtered) coefficient matrix A. We shall discuss concrete examples of coarsening algorithms in the following sections. Now we present two assumptions on the coarse space V_c and prolongation P.

A natural choice of the coarse DOFs is to use a subset of fine-grid DOFs. Under proper re-ordering (coarse variables first and then fine variables) $R = (I, 0) \in \mathbb{R}^{N_c \times N}$. According to Theorem 3.4, we can use the diagonal matrix $D \in \mathbb{R}^{N \times N}$ of A (i.e., the Jacobi method) to analyze the smoother S defined by the point-wise Gauss–Seidel method. Motivated by Lemma 3.4, we can further simplify it and just choose $D := \|A\|_0 I$, for example. This result motivates that we should construct a coarse space, such that

$$\|v - \mathcal{Q}_D v\|_D^2 = \inf_{v_c \in V_c} \|v - v_c\|_D^2 \leqslant C \|v\|_A^2, \quad \forall\, v \in V,$$

where the constant C should be small and uniform with respect to interested parameters like h. This condition is a sufficient condition for the convergence of the two-grid method.

asmp2 **Assumption 6.1** (Weak approximability)**.** $\|(I - PR)v\|_D \leqslant \alpha \|v\|_A, \ \ \forall\, v \in V.$

The matrix PR can be viewed as an approximation of $\mathcal{Q}_D = P(P^T D P)^{-1} P^T D =: PR_*$. Heuristically, the numerical error becomes quite smooth after a few relaxation steps and we can expect the coarse-grid space can approximate v rather accurately if the coarse-grid space is appropriately chosen. Motivated by Theorem 3.3, we have Assumption 6.1, which is equivalent to that V_c reproduces local constant. This condition is equivalent to that V_c has locally supported basis and Assumption 6.2. If v is smooth, i.e., $\|\nabla v\|$ is small, then v can be approximated well in the coarse space V_c.

In view of Remark 3.16, we assume that the prolongation operator preserves the constant (Assumption 6.2). In fact, from the weak approximation property (Assumption 6.1) and let $D := \|A\| I$, we have

$$\|A\|^{1/2} \|v - PR\, v\| \leqslant \alpha \|v\|_A.$$

If v is in the near-null space of A, i.e., $\|v\|_A \approx 0$, then $PR\, v \approx v$. Hence we get the following simplified assumption:

asmp1 **Assumption 6.2** (Constant preserving)**.** $P \mathbf{1}_{N_c} = \mathbf{1}_N$.

:OpComplexity

Remark 6.5 (Operator complexity)**.** When constructing the prolongation P, we must control the sparsity of the coarse level matrices. For efficient overall performance, convergence speed is only one aspect. An equally important aspect is the complexity (sparsity) of the coarser level matrices produced by AMG. We now define a measurement of sparsity, i.e., the operator complexity

$$C_A := \frac{\sum_{l=0}^{L} \mathrm{nnz}(A_l)}{\mathrm{nnz}(A)},$$

where $\mathrm{nnz}(\cdot)$ is the number of nonzeros of a matrix. Apparently, $C_A \geqslant 1$ is always true and $C_A = 1$ corresponds to the one-level methods. During constructing the interpolation operator, we would like to make C_A as close to 1 as possible while keeping good convergence performance.

6.4 Classical algebraic multigrid methods

The original AMG [19] idea (the classical AMG) was developed under the assumption that such a problem with A being an M-matrix was solved. The multilevel hierarchy is constructed based on the coefficient matrix only. Later, the AMG algorithm was further generalized using many heuristics that served to extend its applicability to more general problems. For simplicity, we suppose $A = (a_{i,j}) \in \mathbb{R}^{N\times N}$ be a SPD M-matrix and $\mathsf{G} = (\mathsf{V}, \mathsf{E})$ be its corresponding graph.

Strength of connections

In coarsening, we need to find coarse level variables. This is usually done based on heuristics from properties of M-matrices in AMG methods. Let $\theta_{\mathrm{str}} \in (0, 1)$ be a given real number, usually called *relative strength* parameter. If a pair of indices (i, j) satisfies that

$$-a_{i,j} \geqslant \theta_{\mathrm{str}} \big| \min_k a_{i,k} \big|,$$

then this pair is called *strongly negatively-coupled*. Note that, by this definition, (i, j) and (j, i) are two different pairs by this definition. Of course, there are many different ways to define strongly coupled pairs. For example, we can call i and j strongly negatively- coupled,

$$\text{if } \ a_{i,j} < 0 \ \text{ and } \ |a_{i,j}| > \theta_{\mathrm{str}} \sqrt{a_{i,i} a_{j,j}}.$$

We can easily generalize the concept by considering the positive coupling by removing "negatively".

Denote further

$$\mathsf{S}_i := \big\{ j \in \mathsf{N}_i : \ j \ \text{ strongly coupled to } \ i \big\} \qquad \text{and} \qquad \mathsf{S}_i^T := \big\{ j \in \mathsf{V} : \ i \in \mathsf{S}_j \big\}.$$

So S_i is the set of indices which affects i and S_i^T is the ones which are affected by i. After finding the strongly coupled variables, we can filter the coefficient matrix and obtain A_{S} by removing non-strongly coupled connections.

The above definition strongly coupled variables applies to the direct connections. Sometimes we also need to consider long-range connections; for example, in aggressive coarsening. A variable i is said *strongly connected* to another variable j along a path of length l if there exists a sequence of edges $\{(i, j_1), (j_1, j_2), \ldots, (j_{l-2}, j_{l-1}), (j_{l-1}, j)\} \subseteq \mathsf{E}$ such that $j_{k+1} \in \mathsf{S}_{j_k}$. If there exist at least one path of length less than or equal to l such that i strongly connects to j, then we say that i is l-strongly connects to j and denoted by $j \in \mathsf{S}_i^l$.

We note that, based on the nonzero pattern of the original matrix A^l or a filtered matrix A_{S}^l, one can tell whether there are paths between i and j of length l or not. For example, if we consider five-point stencil finite difference scheme on the mesh given in Figure 6.6 (left). Consider the vertex at the center, the point 13. Then $\mathsf{S}_{13} = \{12, 8, 14, 18\}$ and $\mathsf{S}_{13}^2 = \{12, 8, 14, 18, 11, 3, 15, 23, 7, 17, 9, 19\}$. And we give the weights of A and A^2 in Figure 6.6.

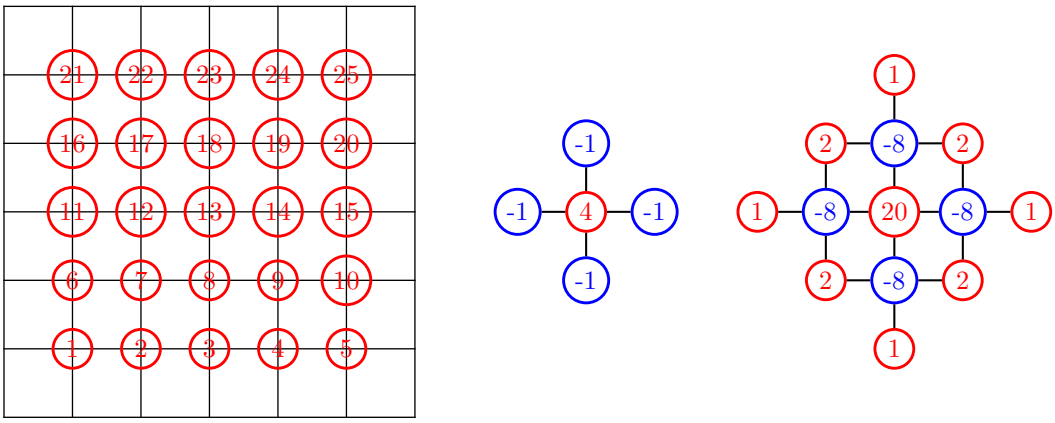


Figure 6.6: Finite difference grid (left), strong connections and weights in A of vertex 13 (middle), and 2-strong connections and weights in A^2 of 13 (right).

ngconnections

C/F splitting

The classical Ruge-Stüben method is to split the set of vertices V to a sum of two non-intersecting sets, the fine variables F and the coarse variables C, such that all the indices in F will be affected by some index in C, while C is expected to contain as few entries as possible. Then F will be chosen as the set of indices of finer grid nodes, and C will be chosen as the set of indices of coarse grid nodes. The indices of nodes are assigned to be coarse or fine successively. Denote by U the

set of indices of nodes that have not been assigned yet, and we summarize the algorithm in the following subroutine:

Listing 6.2: Classical C/F splitting method

```
1 U ← V, C ← ∅, F ← ∅;
2 while U ≠ ∅
3     λ_i ← 2|S_i^T ⋂ F| + |S_i^T ⋂ U|, i ∈ U;
4     k ← arg max{λ_i, i ∈ U};
5     C ← C ⋃ {k}, U ← U\{k};
6     F ← F ⋃ S_k^T, U ← U\S_k^T;
7 end
```

Note that λ_i is a *measure of importance*—It is a measurement about how many points are affected by i. If λ_i is big, we would like to include this point in C; in this way, we can make C contains less points to get bigger coarsening ratio. We weight more on $|\mathsf{S}_i^T \bigcap \mathsf{F}|$ than $|\mathsf{S}_i^T \bigcap \mathsf{U}|$ due to the first part is already determined to be on the fine-grid. In the early stage of coarsening procedure, F does not contain many points, the above algorithm selects a coarse point with as many as neighbors that strongly coupled to it. In the later stage, vertices that strongly coupled to many F-variables are preferred to be selected. We need to add them as F-variables and interpolate indirectly through the F-variables that they are strongly coupled to.

There are a few special cases which require careful treatment during the C/F splitting procedure. We now summarize them in the following remarks:

Remark 6.6 (Isolated points)**.** Before we start the above algorithm, we need to filter out those isolated points (like the Dirichlet boundary points) and define them as F-variables. Similarly, if a point has very strong diagonal dominance, we should also move them to F. These are the trivial cases.

Remark 6.7 (Termination of C/F splitting)**.** If successfully terminated, the set C is an independent set of vertices of the underlying graph G. All F-variables have at least one strongly negatively coupled C-variable, except the trivial ones in the previous remark. However, there might be some U-variables left (with measure $\lambda_i = 0$). They are not strongly negatively coupled to any C-variables or themselves. Furthermore, there are no F-variables are strongly negatively coupled to these points.

Remark 6.8 (Aggressive coarsening)**.** In practice, the standard C/F splitting scheme given above usually results in high operator complexity (refer to Remark 6.5), which leads high computational and storage demands; see Table 6.1. In such cases, we can apply the so-called aggressive coarsening by considering strong connections of length l. Oftentimes $l = 2$ is used. However, A_{S}^2 is expensive to compute and we can apply the regular C/F splitting twice. At

the first pass, find C-variables among all variables using A_{S}; at the second pass, apply the C/F splitting on C-variables using A_{S}^2 but on C only.

Coarsening method	Standard	Aggressive
Operator complexity	2.889	1.606
Setup time (sec)	1.536	1.036
Number of iterations	6	38
Solve time (sec)	0.791	3.293

Table 6.1: Solving 2D five-point stencil finite difference of the Poisson's equation with 1 million DOF using different coarsening methods (stopping criteria for PCG is the relative residual smaller than 10^{-6}).

ab:coarsening

Example 6.2 (Anisotropic elliptic PDE)**.** To illustrate the effect of the above C/F splitting algorithm, we consider an anisotropic diffusion example in §6.1. The computational domain is a unit square. Let us consider the anisotropic diffusion equation $-\epsilon u_{xx} - u_{yy} = 0$ ($\epsilon > 0$). Roughly speaking, we have $\epsilon\|u_{xx}\| \approx \|u_{yy}\|$. This means the solution is smooth in y-direction (low-frequencies); but rough in x direction (high-frequencies). We consider the five-point stencil. The difference equation at the node (x_i, y_j) is

$$-\epsilon\frac{2u_{i,j} - u_{i+1,j} - u_{i-1,j}}{h_x^2} - \frac{2u_{i,j} - u_{i,j-1} - u_{i,j+1}}{h_y^2} = 0.$$

If $\frac{\epsilon}{h_x^2} \ll \frac{1}{h_y^2}$, then $u_{i,j}$ depends on $u_{i,j+1}$ and $u_{i,j-1}$ only. Thus if we process the C/F procedure,

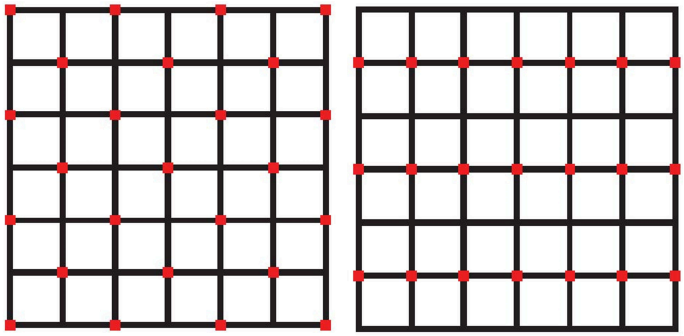


Figure 6.7: C/F splitting for the 2D elliptic problem with $\epsilon = 1$ (left) and $\epsilon \ll 1$ (right), where the red points are C-variables and the black points are F-variables.

ig:coarsening

the coarsening will take place indeed in one direction only (semi-coarsening); see Figure 6.7. □

Construction of prolongation

After obtaining a C/F splitting, upon a reordering or indices, we can always assume the indices of the nodes in C is from 1 to N_c, and those in F are from N_c+1 to N. We can write the stiffness

matrix in the following block structure

$$\begin{pmatrix} A_{\mathsf{C},\mathsf{C}} & A_{\mathsf{C},\mathsf{F}} \\ A_{\mathsf{F},\mathsf{C}} & A_{\mathsf{F},\mathsf{F}} \end{pmatrix} \begin{pmatrix} u_{\mathsf{C}} \\ u_{\mathsf{F}} \end{pmatrix} = \begin{pmatrix} f_{\mathsf{C}} \\ f_{\mathsf{F}} \end{pmatrix}$$

Let $e^H \in \mathbb{R}^{N_c}$ correspond to the variable on coarse grid. We now consider how to prolongate it to $e^h \in \mathbb{R}^N$ corresponding to a variable on fine grid. We use geometric multigrid method for linear finite element method on uniform grids for the 1D Poisson's equation as an example. Let $\{\phi_i^h\}_{i=1}^N$ be the basis of the fine space V and $\{\phi_j^H\}_{j=1}^{N_c}$ be the basis of the coarse space V_c. From the geometrical multigrid point of view, it is natural to expect

$$a(\phi_j^H, \phi_i^h) = 0, \quad j \in \mathsf{C},\ i \in \mathsf{F}. \tag{6.11}$$ eqn:CFconditi

In fact, the main idea is that the fine-grid (high-frequency) part should be captured by fine grid, i.e., $a(u^h - \Pi_H u^h, \phi^h) = a(u^h, \phi^h)$, if ϕ^h is a basis function corresponding to the difference between fine and coarse grid functions.

It is trivial to see that $(Pe^H)_j = e_j^H$, if $j \in \mathsf{C}$. Define

$$P := \begin{pmatrix} I \\ Q \end{pmatrix},$$

where $I \in \mathbb{R}^{N_c \times N_c}$ is the identity matrix and $Q \in \mathbb{R}^{(N-N_c)\times N_c}$. In matrix form, (6.11) can be written as

$$\begin{pmatrix} 0 & 0 \\ 0 & I_{\mathsf{F},\mathsf{F}} \end{pmatrix} \begin{pmatrix} A_{\mathsf{C},\mathsf{C}} & A_{\mathsf{C},\mathsf{F}} \\ A_{\mathsf{F},\mathsf{C}} & A_{\mathsf{F},\mathsf{F}} \end{pmatrix} \begin{pmatrix} I_{\mathsf{C},\mathsf{C}} \\ Q \end{pmatrix} = 0.$$

That is to say, $A_{\mathsf{F},\mathsf{C}} + A_{\mathsf{F},\mathsf{F}}Q = 0$ or $Q = -A_{\mathsf{F},\mathsf{F}}^{-1}A_{\mathsf{F},\mathsf{C}}$. It is easy to check that this prolongation $P = (I,\ Q)^T$ satisfies Assumption 6.2 if the row-sum of A is zero. However this prolongation is too expensive to compute in practice and there are many different ways to approximate Q by a simpler matrix W.

1) Direct interpolation scheme

For the error $e^h \in \mathbb{R}^N$, we have

$$A_{\mathsf{F},\mathsf{F}}e_{\mathsf{F}}^h + A_{\mathsf{F},\mathsf{C}}e_{\mathsf{C}}^h \ll \mathbf{1} \quad \Longrightarrow \quad \sum_{j=1}^{N} a_{i,j}e_j^h \approx 0, \quad i \in \mathsf{F}.$$

Motived by the above observation, we can assume

$$a_{i,i}e_i^h + \sum_{j \in \mathsf{N}_i} a_{i,j}e_j^h = 0, \quad i \in \mathsf{F}. \tag{6.12}$$ eqn:GInterp

This would be an interpolation scheme itself if all points in N_i are C-variables. Of course, it is not always the case. Alternatively, we can throw out the ones that are not strongly negatively coupled and obtain

$$a_{i,i}e_i^h + \sum_{j\in\mathsf{S}_i} a_{i,j}e_j^h = 0, \quad i \in \mathsf{F}. \tag{6.13}$$ eqn:SInterp

We approximate the above equation

$$a_{i,i}e_i^h + \alpha_i \sum_{j\in\mathsf{N}_i\bigcap\mathsf{C}} a_{i,j}e_j^h = 0, \quad \alpha_i = \frac{\sum_{k\in\mathsf{N}_i} a_{i,k}}{\sum_{k\in\mathsf{N}_i\bigcap\mathsf{C}} a_{i,k}}, \quad i \in \mathsf{F}.$$

If row-sum of the i-th row is zero, then $\alpha_i = -\frac{a_{i,i}}{\sum_{k\in\mathsf{P}_i} a_{i,k}}$ and we get an interpolation method

$$e_i^h = \sum_{j\in\mathsf{N}_i\bigcap\mathsf{C}} w_{i,j}e_j^H \qquad \text{and} \qquad w_{i,j} = \frac{a_{i,j}}{\sum_{k\in\mathsf{N}_i\bigcap\mathsf{C}} a_{i,k}}. \tag{6.14}$$ eqn:direct-in

In this case, the matrix form is just $W = \big(\mathrm{diag}(A_{\mathsf{F},\mathsf{C}}\mathbf{1})\big)^{-1}A_{\mathsf{F},\mathsf{C}}$. It is straightforward to show that Assumption 6.2 holds in this case.

We can also make W more sparse by shrinking the support slightly. Define an interpolation set $\mathsf{P}_i := \mathsf{S}_i\bigcap\mathsf{C}$ for $i \in \mathsf{F}$. After further sparsifying the interpolation (by keeping the strongly negatively coupled C-variables only), we get

$$a_{i,i}e_i^h + \alpha_i \sum_{j\in\mathsf{P}_i} a_{i,j}e_j^h = 0, \quad \alpha_i = \frac{\sum_{k\in\mathsf{N}_i} a_{i,k}}{\sum_{k\in\mathsf{P}_i} a_{i,k}}, \quad \forall\, i \in \mathsf{F}.$$

If row-sum of the i-th row is zero, then this gives the well-known direct interpolation

$$e_i^h = \sum_{j\in\mathsf{P}_i} w_{i,j}e_j^H \qquad \text{and} \qquad w_{i,j} = \frac{a_{i,j}}{\sum_{k\in\mathsf{P}_i} a_{i,k}}. \tag{6.15}$$ eqn:direct-in

2) Standard interpolation scheme

We first eliminate all e_j^h for $j \in \mathsf{S}_i\bigcap\mathsf{F}$ by

$$e_j^h := -\sum_{k\in\mathsf{N}_j} a_{j,k}e_k^h/a_{j,j}.$$

This results in

$$\hat{a}_{i,i}e_i^h + \sum_{j\in\hat{\mathsf{N}}_i} \hat{a}_{i,j}e_j^h = 0, \quad \forall\, i \in \mathsf{F},$$

with $\hat{\mathsf{N}}_i = \{j \neq i : \hat{a}_{i,j} \neq 0\}$ and define a new interpolation set $\hat{\mathsf{P}}_i = \big(\bigcup_{j\in\mathsf{S}_i\bigcap\mathsf{F}} \mathsf{S}_j\big)\bigcup\mathsf{S}_i\bigcap\mathsf{C}$. Then we apply the above direct interpolation for this new equation and arrive at the so-called standard interpolation scheme.

3) Jacobi interpolation scheme

We can rewrite the equation (6.13) as

$$a_{i,i}e_i^h + \sum_{j\in \mathsf{P}_i} a_{i,j}e_j^H + \sum_{j\in \mathsf{S}_i\setminus \mathsf{P}_i} a_{i,j}e_j^h = 0, \quad i \in \mathsf{F}.$$

Therefore, in order to obtain an interpolation matrix Q, we just need to approximately solve the above equations for e_i^h ($i \in \mathsf{F}$). For example, we can just apply one Jacobi iteration using $e_j^h \approx \frac{\sum_{k\in \mathsf{P}_i} a_{i,k}e_k^h}{\sum_{k\in \mathsf{P}_i} a_{i,k}}$ as the initial guess of, $j \geqslant N_c + 1,\ j \neq i$. Then the prolongation can be defined as

$$\begin{cases} e_i^h = e_i^H, & i \in \mathsf{C} \\ a_{i,i}e_i^h + \sum_{j\in \mathsf{P}_i} a_{i,j}e_j^h + \sum_{j\in \mathsf{S}_i\setminus \mathsf{P}_i} a_{i,j}\dfrac{\sum_{k\in \mathsf{P}_i} a_{i,k}e_k^h}{\sum_{k\in \mathsf{P}_i} a_{i,k}} = 0, & i \in \mathsf{F}. \end{cases} \tag{6.16}$$

intprol

This is the so-called Jacobi interpolation method.

Remark 6.9 (Some simple alternatives)**.** The biggest advantage of the above approach is that it is simple and local: For the i-th entry, we only need the information on the i-th row of the matrix. We can improve this prolongation matrix P using some straightforward modifications. A few more steps of Jacobi iteration might improve the accuracy. Moreover, the initial guess for the same entry is not the same at the relaxations for different entries. For example, an alternative initial guess could be

$$e_j^h \approx \frac{\sum_{k\in \mathsf{P}_j} a_{j,k}e_k^h}{\sum_{k\in \mathsf{P}_j} a_{j,k}}, \quad j \in \mathsf{F}.$$

Remark 6.10 (Initial guess of W)**.** If the initial guess $W^{(0)}$ preserves constants, then we get

$$Q - W^{(k)} = \left(I - D_{\mathsf{F},\mathsf{F}}^{-1}A_{\mathsf{F},\mathsf{F}}\right)^k \left(Q - W^{(0)}\right).$$

Since both Q and $W^{(0)}$ preserves constants, all improved weights $W^{(k)}$ also preserve constants.

6.5 Aggregation-based algebraic multigrid methods

In this section, we consider the aggregation-base AMG methods whose easy-to-implement feature has drawn quite some attention recently. The idea is to sub-divide the set of vertices into non-intersecting sets (aggregates), i.e., $\mathsf{V} = \bigcup_{j=1,\ldots,N_c} \mathsf{C}_j$. Each aggregate C_j corresponds to a coarser variable.

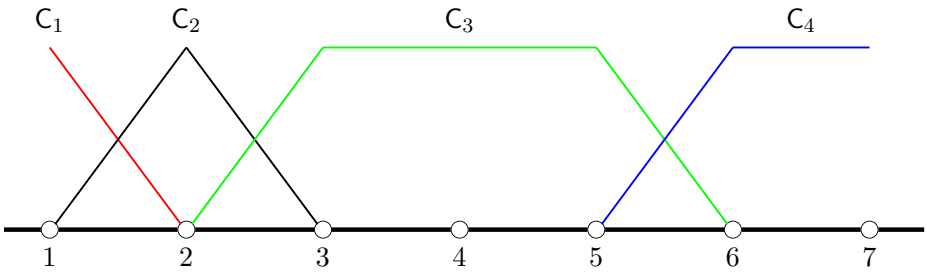
Unsmoothed aggregation AMG

There are several different sophisticated ways to form aggregates. In principle, any combinatorial graph partitioning algorithms can be applied to form aggregation. We now give a simple greedy algorithm to form such an aggregation.

Listing 6.3: A greedy aggregation method

```
1  N_c ← 0, U ← V;
2  for i ∈ U
3     if N_i ⊆ U
4          N_c ← N_c + 1;
5          C_{N_c} ← {i} ⋃ N_i, U ← U\C_{N_c};
6     end
7  end
```

It is possible to have some "left-over" vertices which do not belong to any aggregate after the above procedure. We can, for example, add them to their neighboring aggregates with least points.



ig:aggregates

Figure 6.8: Aggregates and prolongation corresponding to (6.17).

Whence an aggregation is given, it is easy to define the prolongation matrix, for $1 \leqslant i \leqslant N$ and $1 \leqslant j \leqslant N_c$, by

$$\big(P\big)_{i,j} = \begin{cases} 1, & \text{if } i \in \mathsf{C}_j; \\ 0, & \text{if } i \notin \mathsf{C}_j. \end{cases}$$

With this interpolation, it is straight-forward to see that $P\mathbf{1}_{N_c} = \mathbf{1}_N$. We now give an example

to explain P in one dimension. Let

$$P = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 1 \end{pmatrix} \in \mathbb{R}^{N \times N_c}. \tag{6.17}$$ eqP

Figure 6.8 shows the aggregation defined by prolongation P in (6.17).

Of course, there are different ways to form aggregates and we now give another approach. The algorithm to construct coarse grid and prolongation is

Listing 6.4: Another aggregation method

```
1  U ← V;
2  for i ∈ U
3     S_i ← {j ∈ U : j is strongly coupled with i};
4     construct a column of prolongation P based on S_i;
5     U ← U\({i} ⋃ S_i);
6  end
```

Smoothed aggregation AMG

The unsmoothed aggregation methods are very simple but usually converge slowly. There are two ways to improve their convergence behavior. One way is to employ a more complicated nested iteration, like the K-cycle multigrid method discussed in §6.2. And the other way is to enlarge the aggregates and smooth out the basis functions. The latter approach gives the smoothed aggregation AMG methods.

Assume that all variables are partitioned into non-overlapping subsets $\{\mathsf{C}_i\}_{i=1}^{N_c}$. We further assume that each C_i has at least one interior point, i.e., there exists an index $k_i \in \mathsf{C}_i$ such that $(A)_{k_i,j} = 0$ for any $j \notin \mathsf{C}_i$. Suppose that $\mathbf{1}$ is in the null space of A, namely, $A\mathbf{1} = 0$. Define

$$\mathbf{1}_i(x_j) := \begin{cases} \mathbf{1}(x_j), & \text{if } \ j \in \mathsf{C}_i; \\ 0, & \text{otherwise.} \end{cases}$$

Apparently, $\sum_i \mathbf{1}_i = \mathbf{1}$ and $(A\mathbf{1}_i)_{k_i} = 0$. We now smooth out these piecewise basis functions by, for example, one step of weighted Jacobi iteration

$$\psi_i = (I - \omega D^{-1} A)\, \mathbf{1}_i.$$

Hence we have the partition of unity

$$\sum_i \psi_i = (I - \omega D^{-1}A) \sum_i \mathbf{1}_i = (I - \omega D^{-1}A)\,\mathbf{1} = \mathbf{1}.$$

Thus we can obtain

$$\mathbf{1}(x_{k_i}) = \sum_j \psi_j(x_{k_i}) = \sum_j (I - \omega D^{-1}A)\,\mathbf{1}_i(x_{k_i}) = \mathbf{1}_i(x_{k_i}) - \omega D^{-1}A\,\mathbf{1}_i(x_{k_i}),$$

which implies that $D^{-1}A\,\mathbf{1}_i(x_{k_i}) = 0$ and $\psi_i(x_{k_i}) = 1$. We can define the prolongation

$$P_{\text{SA}} := \big(\psi_1, \psi_2, \ldots, \psi_{N_c}\big).$$

Define $\mathbf{1}_c := (1, \ldots, 1)^T \in \mathbb{R}^{N_c}$. Hence we have $P_{\text{SA}}\mathbf{1}_c = \mathbf{1}$. Furthermore, the coarse level matrix $A_c = P_{\text{SA}}^T A P_{\text{SA}}$ satisfies that

$$A_c\mathbf{1}_c = P_{\text{SA}}^T A \sum_i \psi_i = P_{\text{SA}}^T A\mathbf{1} = 0.$$

By applying this definition recursively, we can finish the AMG setup for the smoothed aggregation method.

Listing 6.5: Smoothed aggregation method

```
1 U ← V;
2 for i ∈ U
3     S_i ← {j ∈ U : j is strongly coupled with i};
4     construct a column of prolongation P based on S_i;
5     U ← U\({i} ⋃ S_i);
6 end
7 Smooth the basis functions using the weighted Jacobi method P_SA = (I − ωD^{-1}A)P;
```

We have mentioned in the previous subsection that there are different ways to form aggregates. After forming aggregates one can apply UA or SA to give prolongation. Now we do preliminary tests on aggregation methods for solving the 2D Poisson's equation using the five-point stencil; see Table 6.2. The AMG methods are applied as preconditioners of PCG. Note that, for the SA method, we use the standard V-cycle multigrid in the solve phase; on the other hand, for the UA methods, we use the K-cycle multigrid.

6.6 Homework problems

prob:GMGconv **HW 6.1.** Show the geometric multigrid V-cycle (Algorithm 6.1) is uniformly convergent in $\mathbb{R}^d$.

b:interp-proj **HW 6.2.** If $\mathcal{A} = -\Delta$, show that the interpolant $\mathcal{J}_l : V \mapsto V_l$ is equal to the $(\cdot,\cdot)_{\mathcal{A}}$-projection $\mathit{\Pi}_l : V \mapsto V_l$.

Aggregation method	SA [53]	UA [53]	Pairwise UA [48]
Number of levels	5	5	7
Operator complexity	1.364	1.264	1.332
Setup time (sec)	0.557	0.171	0.277
Number of iterations	16	21	12
Solve time (sec)	1.223	1.696	1.336

Table 6.2: Solving 2D five-point stencil finite difference of the Poisson's equation with 1 million DOF using different aggregation methods (stopping criteria for PCG is the relative residual smaller than 10^{-6}).

b:aggregation

prob:decompH1 **HW 6.3.** Let $\Omega = (0,1)$ and $v \in V_h$ be a $\mathcal{P}_1$ Lagrange finite element function. Show that $|v|_1^2 = \sum_{l=l}^{L} |v_l|_1^2$.

prob:W-cycle **HW 6.4.** Let $q(t) = (1-t)^2$. Show that $\mathcal{B}_{c,2} = (\mathcal{I} - q(\mathcal{B}_c\mathcal{A}_c))\mathcal{A}_c^{-1}$ can be obtained by (6.4).

prob:FMG **HW 6.5.** Show the work estimate of the full multigrid method is $O(N)$.

Chapter 7

Fluid Problems

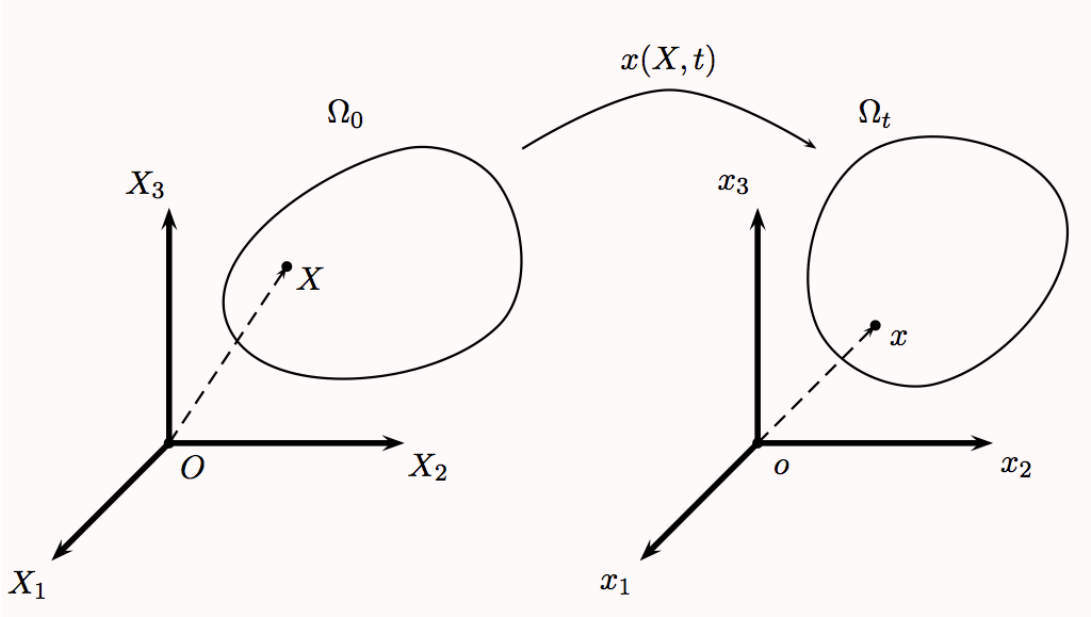
Computational fluid dynamics (CFD) is a branch of fluid mechanics that uses numerical analysis and algorithms to solve and analyze fluid problems. Computers are used to perform the calculations required to simulate liquids or/and gases with surfaces defined by boundary conditions. The fundamental basis of most CFD problems are the Navier–Stokes (NS) equations, which define single-phase fluid flows. These equations can be simplified by removing terms describing viscous actions to yield the Euler equations. These equations can be simplified by dropping the nonlinear convection term to yield the Stokes equation. In this chapter, we discuss multilevel iterative methods suitable for problems arising from CFD.

7.1 The Navier–Stokes equations ⋆

The Navier–Stokes equations describe the motion of viscous fluid substances. These balance equations arise from applying the Newton's second law to fluid motion, together with the assumption that the stress in the fluid is the sum of a diffusing viscous term (proportional to the gradient of velocity) and a pressure term.

Flow map

Let Ω_0 be an open bounded set in $\mathbb{R}^d$ $(d = 2, 3)$. As a convention, we denote the location of a particle in Ω_0 by $X = (X_1, \ldots, X_d)$. This is the configuration at time $t = 0$, which is also called the initial configuration. To describe movement of particles, we denote the current configuration as Ω_t at any time $t \geqslant 0$. The position of a particle at time t is denoted by $x = (x_1, \ldots, x_d)$; see Figure 7.1. The Lagrangian specification of the flow field is a way of looking at particle motion where the observer follows an individual particle as it moves through space and time; see the right figure in Figure 7.1. The Eulerian specification of the flow field is a way of looking at



configuration

Figure 7.1: From initial configuration Ω_0 to current configuration Ω_t.

particle motion that focuses on specific locations in the space through which the fluid flows as time passes; see the left figure in Figure 7.1.

For a vector-valued function $\mathbf{f} : \Omega_t \mapsto \mathbb{R}^d$, the divergence operator can then be written as $\nabla \cdot \mathbf{f} := \sum_{i=1}^d \partial_i \mathbf{f}_i$. The gradient tensor $\nabla \mathbf{f}$ with $(\nabla \mathbf{f})_{i,j} = \partial_j \mathbf{f}_i$. Let $\mathbf{a} \in \mathbb{R}^d$ be a constant vector field and $(\mathbf{a} \cdot \nabla)\mathbf{f} = (\sum_{i=1}^d \mathbf{a}_i \partial_i)\mathbf{f}$. We define an inner product of two gradient matrices $\nabla \mathbf{f} : \nabla \mathbf{g} = \sum_{i=1}^d \nabla \mathbf{f}_i \cdot \nabla \mathbf{g}_i$. Let $\mathbf{u}(\cdot, t) : \Omega_t \mapsto \mathbb{R}^d$ be the velocity field at a fixed time t. The gradient of $\mathbf{u}$ is denoted by $\nabla \mathbf{u} = (\partial_j \mathbf{u}_i)_{i,j}$. Furthermore, $\nabla \mathbf{u}$ is often divided into the symmetric part and the anti-symmetric part. The symmetric gradient is denoted as $\boldsymbol{\varepsilon}(\mathbf{u}) := \frac{1}{2}(\nabla \mathbf{u} + \nabla \mathbf{u}^T)$ and it is the so-called *strain rate*.

We are ready to introduce an important concept to describe trajectory of particles, namely, the flow map $x(X, t)$, which is the trajectory of a particle X along time. We define that

$$\dot{x} = \frac{dx(X,t)}{dt} = \mathbf{u}(x,t) \qquad \text{and} \qquad x(X,0) = X. \tag{7.1}$$

eqn:flowmap

This simple one-dimensional ordinary differential equation (ODE) is called the characteristic equation. Hence $x(\cdot, t)$ is a mapping from the initial configuration Ω_0 to the current configuration Ω_t, or deformation. The deformation gradient and its determinant are then defined as

$$F := \frac{\partial x}{\partial X} \qquad \text{and} \qquad J := |F| = \det(F), \tag{7.2}$$

eqn:defgrad

respectively. F is also called the Jacobian matrix.

For any function $f(\cdot, t) : \Omega_t \mapsto \mathbb{R}$, we can easily derive that

$$\dot{f} := \frac{df(x,t)}{dt} = \nabla f \cdot \frac{dx}{dt} + f_t = f_t + \mathbf{u} \cdot \nabla f, \tag{7.3}$$

eqn:fdot

which is usually called the material derivative of f. Apparently, F and J are functions of t. Using the well-known Jacobi's formula in matrix calculus, we can show that

$$\dot{J} = J \operatorname{tr}(F^{-1}\dot{F}). \tag{7.4}$$ eqn:JacobiFor

Hence we can immediately obtain

$$\dot{J} = J \operatorname{tr}(\frac{\partial X}{\partial x}\frac{\partial \dot{x}}{\partial X}) = J \operatorname{tr}(\nabla \mathbf{u}) = J(\nabla \cdot \mathbf{u}). \tag{7.5}$$ eqn:Jdot

This way, we get an ODE for J, i.e.

$$\dot{J} = (\nabla \cdot \mathbf{u})J \quad \text{and} \quad J(0) = 1. \tag{7.6}$$ eqn:Jode

In fact, we can also obtain the variation of the determinant of F,

$$\delta|F| = |F| \operatorname{tr}(F^{-1}\delta F).$$

We can also derive similar results for the deformation gradient F itself:

$$\dot{F} = \frac{d}{dt}(\frac{\partial x}{\partial X}) = \frac{\partial \dot{x}}{\partial X} = \frac{\partial \mathbf{u}}{\partial X} = \nabla \mathbf{u} F. \tag{7.7}$$ eqn:Fdot

We can easily immediately see that

$$F_t + \mathbf{u} \cdot \nabla F = \nabla \mathbf{u}\, F \qquad \text{and} \qquad F(0) = I. \tag{7.8}$$ eqn:F

Volume and mass conservation

A very useful trick for doing calculus in continuum mechanics is the pull-back (from Ω_t to Ω_0) and push-forward (from Ω_0 to Ω_t) argument. We first give an example:

$$\begin{aligned}
\frac{d}{dt}\int_{\Omega_t} f(x,t)\,dx &= \frac{d}{dt}\int_{\Omega_0} f(x(X,t),t)J\,dX \\
&= \int_{\Omega_0} \frac{d}{dt} f(x(X,t),t)J\,dX + \int_{\Omega_0} f(x(X,t),t)\dot{J}\,dX \\
&= \int_{\Omega_0} (f_t + \mathbf{u}\cdot\nabla f + f\nabla\cdot\mathbf{u})J\,dX \\
&= \int_{\Omega_t} \dot{f} + f\nabla\cdot\mathbf{u}\,dx \;= \int_{\Omega_t} f_t + \nabla\cdot(f\mathbf{u})\,dx.
\end{aligned} \tag{7.9}$$

This identity is often called the *transport formula*.

lem:transport **Lemma 7.1** (Transport formula)**.** *For a function $f : \Omega_t \mapsto \mathbb{R}$ and $\mathbf{u}(x,t) := \frac{dx(X,t)}{dt}$, we have*

$$\frac{d}{dt}\int_{\Omega_t} f(x,t)\,dx = \int_{\Omega_t} f_t + \nabla\cdot(f\mathbf{u})\,dx = \int_{\Omega_0} (f_t + \mathbf{u}\cdot\nabla f + f\nabla\cdot\mathbf{u})J\,dX.$$

For a domain $\Omega \subset \mathbb{R}^d$, we denote its volume (or area) as $|\Omega|$. We then find that

$$|\Omega_t| = \int_{\Omega_t} 1 \, dx = \int_{\Omega_0} J \, dX = J|\Omega_0|$$

For *incompressible* fluids, we have that the volume preserving property

$$|\Omega_t| \equiv |\Omega_0| \qquad \text{or} \qquad J(t) \equiv 1.$$

From the equation (7.6), we can derive that $\nabla \cdot \mathbf{u} = 0$. This is the so-called *divergence-free* condition.

Denote the density of the material occupying Ω_t by $\rho(x,t)$. According to the equation (7.9), for any region $\omega_t \subset \Omega_t$, we have that

$$\frac{d}{dt} \int_{\omega_t} \rho(x,t) \, dx = \int_{\omega_t} \rho_t + \nabla \cdot (\rho \mathbf{u}) \, dx$$

Since this identity holds for any ω, we immediately see that

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0 \qquad \text{and} \qquad \dot{\rho} + \rho \nabla \cdot \mathbf{u} = 0, \tag{7.10}$$

eqn:masscon

which is called the equation of mass conservation or the continuity equation.

It is clear that integrating the density over any domain ω_t gives the mass. Due to mass conservation, we have that

$$\int_{\omega_0} \rho_0(X) \, dX = \int_{\omega_t} \rho(x,t) \, dx = \int_{\omega_0} \rho(x(X,t),t) J \, dX.$$

Hence, we have the relation

$$\rho(x(X,t),t) = \frac{\rho_0(X)}{J}. \tag{7.11}$$

eqn:rhoratio

If the incompressible condition $\nabla \cdot \mathbf{u} = 0$ holds, we obtain that $\rho(x(X,t),t) = \rho_0(X)$.

If $\rho \equiv \rho_0$ is a constant, then (7.10) gives the divergence-free condition immediately. On the other hand, if we assume incompressibility, we can get a simplified equation:

$$\rho_t + (\mathbf{u} \cdot \nabla)\rho = 0 \qquad \text{or} \qquad \dot{\rho} = 0. \tag{7.12}$$

eqn:masscon1

Together with $\rho(X,0) = \rho_0$ being a constant, we can get $\rho \equiv \rho_0$ for all time $t \in [0,T]$.

Balance of momentum

Now we consider the incompressible Newtonian fluids. Due to the Newton's Second Law, we have the balance of momentum

$$\frac{d}{dt} \int_{\Omega_t} \rho \mathbf{u} \, dx = \text{Force}(\Omega_t). \tag{7.13}$$

eqn:Newton2

The left-hand side of the above equation is the rate of change for the momentum. Using the transport formula (Lemma 7.1), we derive that

$$\frac{d}{dt}\int_{\Omega_t}\rho\mathbf{u}\,dx = \int_{\Omega_0}(\rho_t + \mathbf{u}\cdot\nabla\rho)\mathbf{u} + \rho(\mathbf{u}_t + \mathbf{u}\cdot\nabla\mathbf{u})\,dX.$$

Due to the mass conservation and incompressibility (7.12), we then have

$$\frac{d}{dt}\int_{\Omega_t}\rho\mathbf{u}\,dx = \int_{\Omega_t}\rho(\mathbf{u}_t + \mathbf{u}\cdot\nabla\mathbf{u})\,dx. \tag{7.14}$$ eqn:momentumL

On the other hand, the right-hand side of the Newton's Second Law is the total force acting on Ω_t. We have, from the divergence theorem, that

$$\text{Force}(\Omega_t) := \int_{\Omega_t}\mathbf{f}\,dx + \int_{\partial\Omega_t}\mathbf{T}\cdot\mathbf{n}\,dS = \int_{\Omega_t}\mathbf{f} + \nabla\cdot\mathbf{T}\,dx,$$

where $\mathbf{f}$ is the total external body force (like gravity), $\mathbf{T}$ is the traction tensor on the boundary of Ω_t, and $\mathbf{n}$ is the outer normal direction on the boundary $\partial\Omega_t$. The exact form of $\mathbf{T}$ depends on the underlying constitutive laws. For Newtonian fluids, the traction can be defined as

$$\mathbf{T} := -pI + 2\mu\boldsymbol{\varepsilon}(\mathbf{u}), \tag{7.15}$$ eqn:T

where p is the pressure and μ is the viscosity.

For incompressible fluids, we have $\nabla\cdot\mathbf{u} = 0$. In turn, we can obtain (see HW 7.1) that

$$\big(\nabla\cdot(2\boldsymbol{\varepsilon}(\mathbf{u}))\big)_j = \sum_{i=1}^{d}\partial_i(\mathbf{u}_{i,j} + \mathbf{u}_{j,i}) = \sum_{i=1}^{d}\partial_j\mathbf{u}_{i,i} + \sum_{i=1}^{d}\partial_i\mathbf{u}_{j,i} = \Delta\mathbf{u}_j,$$

which means

$$2\,\nabla\cdot\boldsymbol{\varepsilon}(\mathbf{u}) = \Delta\mathbf{u}. \tag{7.16}$$ eqn:div_Lapla

This way we can get the momentum equation (balance of force) for incompressible Newtonian fluids:

$$\rho(\mathbf{u}_t + \mathbf{u}\cdot\nabla\mathbf{u}) = -\nabla p + \mu\Delta\mathbf{u}. \tag{7.17}$$ eqn:momentum

If the density ρ is a constant, we further simplify the above equation (by modifying the definition of p and μ) to give

$$\mathbf{u}_t + \mathbf{u}\cdot\nabla\mathbf{u} = -\nabla p + \mu\Delta\mathbf{u}. \tag{7.18}$$ eqn:momentum1

ssc:model

Mathematical models

To summarize, we have derived the mathematical model for incompressible Newtonian fluids, i.e., the *Navier–Stokes* (NS) equations:

$$\begin{cases} \rho(\mathbf{u}_t + \mathbf{u}\cdot\nabla\mathbf{u}) - \mu\Delta\mathbf{u} + \nabla p = \mathbf{f}, & \Omega_t & \text{balance of momentum;} \\ \rho_t + \nabla\cdot(\rho\mathbf{u}) = 0, & \Omega_t & \text{conservation of mass;} \\ \nabla\cdot\mathbf{u} = 0, & \Omega_t & \text{incompressibility;} \\ \mathbf{u} = 0, & \partial\Omega_t & \text{no-slip boundary;} \\ \mathbf{u}|_{t=0} = \mathbf{u}_0, & \Omega_t & \text{initial condition.} \end{cases} \tag{7.19}$$

eqn:NS

If we assume the density ρ is a constant, then we can write (7.19) as follows:

$$\begin{cases} \mathbf{u}_t + \mathbf{u}\cdot\nabla\mathbf{u} - \mu\Delta\mathbf{u} + \nabla p = \mathbf{f}, & \Omega_t & \text{momentum equation;} \\ \nabla\cdot\mathbf{u} = 0, & \Omega_t & \text{continuity equation;} \\ \mathbf{u} = 0, & \partial\Omega_t & \text{no-slip boundary;} \\ \mathbf{u}|_{t=0} = \mathbf{u}_0, & \Omega_t & \text{initial condition.} \end{cases} \tag{7.20}$$

eqn:NS1

Now we have the mathematical model for incompressible viscous Newtonian fluids. If we consider ideal fluids (viscosity $\mu = 0$) and assume that there is no external body force ($\mathbf{f} = 0$), then we get the incompressible *Euler* equations:

$$\begin{cases} \rho(\mathbf{u}_t + \mathbf{u}\cdot\nabla\mathbf{u}) + \nabla p = 0, & \Omega_t & \text{balance of momentum;} \\ \rho_t + \nabla\cdot(\rho\mathbf{u}) = 0, & \Omega_t & \text{conservation of mass;} \\ \nabla\cdot\mathbf{u} = 0, & \Omega_t & \text{incompressibility;} \\ \mathbf{u}\cdot\mathbf{n} = 0, & \partial\Omega_t & \text{no-flow boundary;} \\ \mathbf{u}|_{t=0} = \mathbf{u}_0, & \Omega_t & \text{initial condition.} \end{cases} \tag{7.21}$$

eqn:Euler

If the density ρ is a constant, then we have the following simplified form:

$$\begin{cases} \mathbf{u}_t + \mathbf{u}\cdot\nabla\mathbf{u} + \nabla p = 0, & \Omega_t & \text{momentum equation;} \\ \nabla\cdot\mathbf{u} = 0, & \Omega_t & \text{continuity equation;} \\ \mathbf{u}\cdot\mathbf{n} = 0, & \partial\Omega_t & \text{no-flow boundary;} \\ \mathbf{u}|_{t=0} = \mathbf{u}_0, & \Omega_t & \text{initial condition.} \end{cases} \tag{7.22}$$

eqn:Euler1

For numerical simulation of the Navier–Stokes and Euler equations, there are several technical difficulties. First of all, the incompressibility condition is a constraint on the velocity field and appropriate finite element spaces need to be selected to discretize this mixed problem. Secondly, these equations have a nonlinear convection term; when the viscosity coefficient μ is small (corresponding to high Reynolds number), the convection is essentially dominant.

7.2 The Stokes equations

For simplicity, we now focus on a linearized problem of the Navier–Stokes equation, namely the Stokes equation.

The time-dependent Stokes equation

On an open bounded set $\Omega \subset \mathbb{R}^d$, we consider

$$\begin{cases} \mathbf{u}_t - \mu\Delta\mathbf{u} + \nabla p &= \mathbf{f}, & \Omega; \\ \nabla\cdot\mathbf{u} &= 0, & \Omega; \\ \mathbf{u} &= 0, & \partial\Omega; \\ \mathbf{u}|_{t=0} &= \mathbf{u}_0, & \Omega. \end{cases} \tag{7.23}$$ eqn:Stokes

This set of equations is usually referred to as the time-dependent Stokes equations. After time discretization, we need to solve the Stokes-like equations

$$\begin{cases} (I - \epsilon^2\Delta)\mathbf{u} + \nabla p &= \mathbf{f}, & \Omega; \\ \nabla\cdot\mathbf{u} &= 0, & \Omega; \\ \mathbf{u} &= 0, & \partial\Omega. \end{cases} \tag{7.24}$$ eqn:Stokes1ti

We can further simplify the discussion and only consider the following steady-state Stokes equations, i.e.,

$$\begin{cases} -\Delta\mathbf{u} + \nabla p &= \mathbf{f}, & \Omega; \\ \nabla\cdot\mathbf{u} &= 0, & \Omega; \\ \mathbf{u} &= 0, & \partial\Omega. \end{cases} \tag{7.25}$$ eqn:Stokes1

Let $\mathscr{V} := \left[H_0^1(\Omega)\right]^d$ and $\mathscr{Q} := L_0^2(\Omega) = \left\{q \in L^2(\Omega) : \int_\Omega q = 0\right\}$. The weak form of the Stokes equation (7.25) can be written as: Find $\mathbf{u} \in \mathscr{V}$ and $p \in \mathscr{Q}$, such that

$$\begin{cases} 2\int_\Omega \boldsymbol{\varepsilon}(\mathbf{u}) : \boldsymbol{\varepsilon}(\mathbf{v})\, dx + (p, \nabla\cdot\mathbf{v}) &= (\mathbf{f}, \mathbf{v}), & \forall\, \mathbf{v} \in \mathscr{V}; \\ (\nabla\cdot\mathbf{u}, q) &= 0, & \forall\, q \in \mathscr{Q}. \end{cases} \tag{7.26}$$ eqn:Stokes1w

The derivation is straightforward and hence leave to the readers; see HW 7.2.

Remark 7.1 (Constrained energy minimization)**.** We can view the Stokes equations as a constrained energy minimization problem

$$\min_{\mathbf{v}\in\mathscr{Z}} \int_\Omega \boldsymbol{\varepsilon}(\mathbf{v}) : \boldsymbol{\varepsilon}(\mathbf{v})\, dx - \int_\Omega \mathbf{f}\cdot\mathbf{v}\, dx,$$

where $\mathscr{Z} := \{\mathbf{v} \in \mathscr{V} : \nabla\cdot\mathbf{v} = 0\}$ is the subspace of divergence-free functions. The equation (7.26) is the first-order optimality condition of this constrained minimization problem and p is the Lagrange multiplier.

The Brezzi theory

Let $\mathscr{V}'$ and $\mathscr{Q}'$ be the dual spaces of $\mathscr{V}$ and $\mathscr{Q}$, respectively. Generally speaking, we can put the Stokes problem in an abstract framework and consider the following *saddle-point* problem: For any given $(f, g) \in \mathscr{V}' \times \mathscr{Q}'$, find a pair $(u, p) \in \mathscr{V} \times \mathscr{Q}$, such that the following system holds

$$\begin{cases} a[u, v] + b[v, p] &= \langle f, v \rangle, \quad \forall v \in \mathscr{V}; \\ b[u, q] &= \langle g, q \rangle, \quad \forall q \in \mathscr{Q}. \end{cases} \tag{7.27}$$ eqn:saddle

Here $a[\cdot, \cdot] : \mathscr{V} \times \mathscr{V} \mapsto \mathbb{R}$ and $b[\cdot, \cdot] : \mathscr{V} \times \mathscr{Q} \mapsto \mathbb{R}$ are *continuous* bilinear forms, i.e.,

$$\begin{aligned} a[u, v] &\leqslant C_a \|u\|_{\mathscr{V}} \|v\|_{\mathscr{V}}, \quad \forall u, v \in \mathscr{V}, \\ b[u, p] &\leqslant C_b \|u\|_{\mathscr{V}} \|p\|_{\mathscr{Q}}, \quad \forall u \in \mathscr{V}, p \in \mathscr{Q}. \end{aligned}$$

We can identify a linear operator $\mathcal{A} : \mathscr{V} \mapsto \mathscr{V}'$ such that

$$\langle \mathcal{A}u, v \rangle = a[u, v], \quad \forall u \in \mathscr{V},\ v \in \mathscr{V}$$

and another linear operator $\mathcal{B} : \mathscr{V} \mapsto \mathscr{Q}'$ (or its adjoint $\mathcal{B}^T : \mathscr{Q} \mapsto \mathscr{V}'$) such that

$$\langle \mathcal{B}u, p \rangle = \left\langle u, \mathcal{B}^T p \right\rangle = b[u, p], \quad \forall u \in \mathscr{V},\ p \in \mathscr{Q}.$$

Hence (7.27) can be written in the following operator form

$$\begin{cases} \mathcal{A}u + \mathcal{B}^T p &= f, \\ \mathcal{B}u &= g. \end{cases}$$

We now analyze under what condition(s) the weak formulation (7.27) is well-posed. We define the kernel space of $\mathcal{B}$ as

$$\mathscr{Z} := \mathrm{null}(\mathcal{B}) = \left\{ v \in \mathscr{V} \,:\, b[v, q] = 0,\ \forall q \in \mathscr{Q} \right\} \subset \mathscr{V}.$$

Because $b[\cdot, \cdot]$ is continuous, $\mathscr{Z}$ is closed. Hence we can give an orthogonal decomposition

$$\mathscr{V} = \mathscr{Z} \oplus \mathscr{Z}^{\perp} = \mathrm{null}(\mathcal{B}) \oplus \mathrm{null}(\mathcal{B})^{\perp}.$$

For any $u \in \mathscr{V}$, we have $u = u_0 + u_{\perp}$, with $u_0 \in \mathrm{null}(\mathcal{B})$ and $u_{\perp} \in \mathrm{null}(\mathcal{B})^{\perp}$.

In order to solve $\mathcal{B}u = g$, we only need to solve $\mathcal{B}u_{\perp} = g$. Using the inf-sup theory discussed in §1.1, we can see that, if $\mathcal{B}$ is surjective, namely,

$$\inf_{q \in \mathscr{Q}} \sup_{v \in \mathscr{V}} \frac{b[v, q]}{\|v\|_{\mathscr{V}} \|q\|_{\mathscr{Q}}} = \beta > 0, \tag{7.28}$$ eqn:infsupB

then $u_\perp$ exists. Furthermore, it is easy to see that $u_\perp$ is also unique[1]. Hence we have $\mathcal{B}: \mathscr{Z}^\perp \mapsto \mathscr{Q}'$ and $\mathcal{B}^T: \mathscr{Q} \mapsto (\mathscr{Z}^\perp)'$ are isomorphisms.

Now we only need to show the existence and uniqueness of the following problem: Find $u_0 \in \mathscr{Z}$, such that

$$a[u_0, v] = \langle f, v\rangle - a[u_\perp, v], \quad \forall\, v \in \mathscr{Z}.$$

According to the Nečas Theorem 1.4, we know that the existence and uniqueness of u_0 is equivalent to the following inf-sup conditions

$$\inf_{u\in\mathscr{Z}} \sup_{v\in\mathscr{Z}} \frac{a[u,v]}{\|u\|_{\mathscr{V}}\|v\|_{\mathscr{V}}} = \inf_{v\in\mathscr{Z}} \sup_{u\in\mathscr{Z}} \frac{a[u,v]}{\|u\|_{\mathscr{V}}\|v\|_{\mathscr{V}}} = \alpha > 0. \tag{7.29}$$ eqn:infsupA

With the conditions (7.29) and (7.28), we obtain a unique solution $u = u_0 + u_\perp$.

We can find the solution for the pressure variable by solving

$$\mathcal{B}^T p = f - \mathcal{A}u. \tag{7.30}$$ eqn:Bp

For any $v \in \mathscr{Z} = \text{null}(\mathcal{B})$, it is easy to see that

$$\langle f - \mathcal{A}u, v\rangle = \langle \mathcal{B}^T p, v\rangle = \langle p, \mathcal{B}v\rangle = 0.$$

Hence, $f - \mathcal{A}u \in (\mathscr{Z}^\perp)' = \{w \in \mathscr{V}' : \langle w, v\rangle = 0,\ \forall v \in \mathscr{Z}\}$. Because $\mathcal{B}^T : \mathscr{Q} \mapsto (\mathscr{Z}^\perp)'$ is an isomorphism, there is a unique solution to (7.30).

Hence we obtain the following well-posedness result [21, Theorem 1.1]:

thm:Brezzi **Theorem 7.1** (Brezzi Theorem)**.** *For continuous bilinear forms $a[\cdot,\cdot]$ and $b[\cdot,\cdot]$, the saddle-point problem (7.27) is well-posed if and only if (7.29) and (7.28) hold. Furthermore, the solution (u,p) satisfies the stability condition*

$$\|u\|_{\mathscr{V}} + \|p\|_{\mathscr{Q}} \lesssim \|f\|_{\mathscr{V}'} + \|g\|_{\mathscr{Q}'}.$$

mixed-inf-sup **Remark 7.2** (Inf-sup condition of the mixed formulation)**.** Let $\mathscr{X} := \mathscr{V} \times \mathscr{Q}$. We define a new bilinear form $\tilde{a}: \mathscr{X} \times \mathscr{X} \mapsto \mathbb{R}$

$$\tilde{a}[(u,p),(v,q)] := a[u,v] + b[v,p] + b[u,q].$$

Then the saddle-point problem (7.27) is equivalent to finding $(u,p) \in \mathscr{X}$ such that

$$\tilde{a}[(u,p),(v,q)] = \langle f, v\rangle + \langle g, q\rangle, \quad \forall (v,q) \in \mathscr{X}. \tag{7.31}$$ eqn:saddle2

If both $a[\cdot,\cdot]$ and $b[\cdot,\cdot]$ are continuous, then $\tilde{a}[\cdot,\cdot]$ is also continuous. If $a[\cdot,\cdot]$ and $b[\cdot,\cdot]$ satisfy the standard Brezzi conditions (7.29) and (7.28), respectively, then $\tilde{a}[\cdot,\cdot]$ satisfies the inf-sup condition as well.

[1] Suppose there is another solution $\tilde{u}_\perp$, then $\mathcal{B}(u_\perp - \tilde{u}_\perp) = 0$. In turn, we have $u_\perp - \tilde{u}_\perp$ is in $\text{null}(\mathcal{B})$. Due to $u_\perp - \tilde{u}_\perp \in \text{null}(\mathcal{B})^\perp$, we find $u_\perp - \tilde{u}_\perp = 0$.

Well-posedness of the Stokes problem

In view of the general theory developed in the previous subsection, we can define

$$a[\mathbf{u},\mathbf{v}] := 2\int_\Omega \boldsymbol{\varepsilon}(\mathbf{u}) : \boldsymbol{\varepsilon}(\mathbf{v})\, dx \qquad \mathcal{A} := -\Delta \tag{7.32}$$

$$b[\mathbf{v}, q] := -\int_\Omega \nabla\cdot\mathbf{v}\, q\, dx \qquad \mathcal{B} := -\nabla\cdot,\ \mathcal{B}^T := \nabla \tag{7.33}$$

In this case, the inf-sup condition (7.29) is trivial since the coercive condition holds, i.e.,

$$\int_\Omega \boldsymbol{\varepsilon}(\mathbf{u}) : \boldsymbol{\varepsilon}(\mathbf{u}) \geqslant \alpha\|\mathbf{u}\|_1^2, \quad \forall \mathbf{u} \in [H_0^1(\Omega)]^d.$$

Hence we only need to check the inf-sup condition for $b[\cdot,\cdot]$.

Lemma 7.2 (Inf-sup condition for divergence operator)**.** *For any $q \in \mathscr{Q} = L_0^2(\Omega)$, there exists $\mathbf{v} \in \mathscr{V} = [H_0^1(\Omega)]^d$ such that*

$$\nabla\cdot\mathbf{v} = q \qquad \textit{and} \qquad \|\mathbf{v}\|_1 \lesssim \|q\|_0.$$

So the inf-sup condition (7.28) holds.

Proof. This non-trivial result goes back to Nečas and a proof can be found in [33, Ⅲ.3.1]. □

Remark 7.3 (Existence of solution)**.** It has been shown in the above lemma that range$(\mathcal{B}) = L^2(\Omega)/\mathbb{R} \cong \mathscr{Q}$. Or equivalently, we have null$(\mathcal{B}^T) \bigcap \mathscr{Q} = \{0\}$.

Using the previous lemma and the Brezzi theorem, we can easily get the following result:

kes_wellposed **Theorem 7.2** (Well-posedness of the Stokes equations)**.** *There exists a unique solution $(\mathbf{u}, p) \in [H_0^1(\Omega)]^d \times L_0^2(\Omega)$ to the weak form of the Stokes equation (7.26) and*

$$\|\mathbf{u}\|_1 + \|p\|_0 \lesssim \|\mathbf{f}\|_{-1}.$$

Penalty method for the Stokes problem ⋆

In general, there are two approaches to approximate the Stokes problem. The first one is to approximate (7.26) directly. An alternative method is to formulate the original problem using a penalty method as

$$\text{Find}\ \ \mathbf{u} \in \mathscr{V}: \quad 2\int_\Omega \boldsymbol{\varepsilon}(\mathbf{u}) : \boldsymbol{\varepsilon}(\mathbf{v})\, dx + \gamma(\nabla\cdot\mathbf{u},\, \nabla\cdot\mathbf{v}) = (\mathbf{f}, \mathbf{v}), \quad \forall\, \mathbf{v} \in \mathscr{V}. \tag{7.34}$$

eqn:Stokes2w

The above equation can also be seen in the linear elasticity problems and it is known for causing the *locking* phenomena[2] for many finite element methods when γ is big. This is usually caused by overly constraint on the velocity space. To cure such a problem, penalty methods introduce selective or reduced integration procedures. It has been shown that penalty methods are sometimes equivalent to mixed methods [42].

[2] The computed velocity is vanishing or unnaturally small for big λ.

7.3 Mixed finite element methods

In this section, we consider conforming mixed finite element methods for the Stokes equations. Let $V_h \subset \mathscr{V} = [H_0^1(\Omega)]^d$ and $Q_h \subset \mathscr{Q} = L_0^2(\Omega)$ be finite dimensional spaces. Find $\mathbf{u}_h \in V_h$ and $p_h \in Q_h$, such that

$$\left\{ \begin{array}{lcll} 2\displaystyle\int_\Omega \boldsymbol{\varepsilon}(\mathbf{u}_h) : \boldsymbol{\varepsilon}(\mathbf{v}_h)\, dx - (p_h,\, \nabla\cdot\mathbf{v}_h) & = & (\mathbf{f}, \mathbf{v}_h), & \forall\, \mathbf{v}_h \in V_h, \\ (\nabla\cdot\mathbf{u}_h,\, q_h) & = & 0, & \forall\, q_h \in Q_h. \end{array} \right. \tag{7.35}$$ eqn:Stokes1d

The existence of the discrete solution $(\mathbf{u}_h, p_h)$ is straightforward due to the conformity of the approximation spaces.

Well-posedness and convergence

Let $Z_h = \mathrm{null}(\mathcal{B}_h)$ be the kernel of the discrete divergence operator. In fact, the coercivity of $a[\cdot,\cdot]$ yields that

$$\inf_{\mathbf{u}_h\in Z_h} \sup_{\mathbf{v}_h\in Z_h} \frac{a[\mathbf{u}_h, \mathbf{v}_h]}{\|\mathbf{u}_h\|_1 \|\mathbf{v}_h\|_1} = \alpha_h > 0. \tag{7.36}$$ eqn:inf_sup_A

If $Z_h \subset \mathscr{Z}$ and the coercivity condition holds, we have the following optimal approximation property by the Céa's lemma (Lemma 3.1):

$$\|\mathbf{u} - \mathbf{u}_h\|_{\mathscr{V}} \leqslant \frac{C_a}{\alpha_h} \inf_{\mathbf{v}_h\in Z_h} \|\mathbf{u} - \mathbf{v}_h\|_{\mathscr{V}}.$$

However, it is not easy to make the finite element kernel space $Z_h \subset \mathscr{Z}$. A sufficient condition for this inclusion property is $\mathcal{B}(V_h) \subset Q_h$, which suggests Q_h should be large enough for a fixed space V_h. In fact, we have

$$\mathcal{B}_h\mathbf{u}_h = 0, \quad \text{in } Q_h' \quad \Longleftrightarrow \quad (\mathcal{B}\mathbf{u}_h, q_h) = 0, \quad \forall q_h \in Q_h.$$

Furthermore, we also have

$$\mathcal{B}\mathbf{u}_h = 0, \quad \text{in } \mathscr{Q}' \quad \Longleftrightarrow \quad (\mathcal{B}\mathbf{u}_h, q) = 0, \quad \forall q \in \mathscr{Q}.$$

If $\mathbf{u}_h \in Z_h$ and $q \in \mathscr{Q}$, then $(\mathcal{B}\mathbf{u}_h, q) = (\mathcal{B}\mathbf{u}_h, q_0 + q_\perp) = (\mathcal{B}\mathbf{u}_h, q_0) + (\mathcal{B}\mathbf{u}_h, q_\perp) = 0$, where $q = q_0 + q_\perp$ with $q_0 \in Q_h$. Notice that $(\mathcal{B}\mathbf{u}_h, q_\perp) = 0$ because the inclusion condition $\mathcal{B}(V_h) \subset Q_h$.

If $Z_h \not\subset \mathscr{Z}$, then there is a variational crime and we have following estimate:

$$\|\mathbf{u} - \mathbf{u}_h\|_{\mathscr{V}} \leqslant \Big(1 + \frac{C_a}{\alpha_h}\Big) \inf_{\mathbf{v}\in Z_h} \|\mathbf{u} - \mathbf{v}\|_{\mathscr{V}} + \frac{1}{\alpha_h} \sup_{\mathbf{w}\in Z_h\setminus\{0\}} \frac{\big|a[\mathbf{u} - \mathbf{u}_h, \mathbf{w}]\big|}{\|\mathbf{w}\|_{\mathscr{V}}}.$$

For $\mathbf{w} \in Z_h$, we have

$$a[\mathbf{u} - \mathbf{u}_h, \mathbf{w}] = a[\mathbf{u}, \mathbf{w}] - (\mathbf{f}, \mathbf{v}) = -b[\mathbf{w}, p] = -b[\mathbf{w}, p - q],$$

for any $q \in Q_h$. Because $b[\cdot,\cdot]$ is continues, we find that

$$\big|a[\mathbf{u}-\mathbf{u}_h,\mathbf{w}]\big| \leqslant C_b\|\mathbf{w}\|_{\mathscr{V}}\|p-q\|_{\mathscr{Q}}.$$

We can then conclude with the following best approximation result:

Lemma 7.3 (Quasi-optimality for velocity)**.** *Let $V_h \subset \mathscr{V}$ and $Q_h \subset \mathscr{Q}$. If the bilinear form $a[\cdot,\cdot]$ is coercive, then we have*

$$\|\mathbf{u}-\mathbf{u}_h\|_{\mathscr{V}} \leqslant \Big(1+\frac{C_a}{\alpha_h}\Big)\inf_{\mathbf{v}\in Z_h}\|\mathbf{u}-\mathbf{v}\|_{\mathscr{V}} + \frac{C_b}{\alpha_h}\inf_{q\in Q_h}\|p-q\|_{\mathscr{Q}}.$$

We have the identity

$$(\mathcal{B}_h\mathbf{u}_h, q_h) = b[\mathbf{u}_h,q_h] = (\mathcal{B}\mathbf{u}_h,q_h), \quad \forall q_h \in Q_h.$$

In the other words, $\mathcal{B}_h\mathbf{u}_h$ is the L^2-projection of $\mathcal{B}\mathbf{u}_h$ onto Q_h. If $\mathrm{null}(\mathcal{B}_h^T)$ is not trivial, then $\mathrm{range}(\mathcal{B}_h)$ is strictly included in Q_h. This could lead to ill-posed problems. For a fixed Q_h, the velocity approximation space V_h should be rich enough in order to guarantee the discrete inf-sup condition:

$$\inf_{q_h\in Q_h}\sup_{\mathbf{v}_h\in V_h}\frac{b(\mathbf{v}_h,q_h)}{\|\mathbf{v}_h\|_1\|q_h\|_0} = \beta_h > 0. \tag{7.37}$$ eqn:inf_sup_B

The condition $\mathrm{null}(\mathcal{B}_h^T) = \{0\}$ is necessary for the inf-sup condition above. If $\mathrm{null}(\mathcal{B}_h^T)$ is non-trivial, then the numerical solution p_h is not unique, namely, $p_h + s_h$ is also a solution when $s_h \in \mathrm{null}(\mathcal{B}_h^T)$. In this case, we usually find the computed pressure is oscillatory and, hence, $\mathrm{null}(\mathcal{B}_h^T)$ is often referred to as the space of spurious pressure modes.

Theorem 7.3 (Quasi-optimality)**.** *Let $V_h \subset \mathscr{V}$ and $Q_h \subset \mathscr{Q}$. If the bilinear form $a[\cdot,\cdot]$ is coercive and the inf-sup condition (7.37) holds with $\beta_h \geqslant \beta_0 > 0$, then we have*

$$\|\mathbf{u}-\mathbf{u}_h\|_{\mathscr{V}} + \|p-p_h\|_{\mathscr{Q}} \lesssim \inf_{\mathbf{v}\in Z_h}\|\mathbf{u}-\mathbf{v}\|_{\mathscr{V}} + \inf_{q\in Q_h}\|p-q\|_{\mathscr{Q}}.$$

Some stable finite element pairs ⋆

From the above discussions, we conclude that: To balance computational efforts and convergence rates for the velocity in $[H_0^1(\Omega)]^d$ and the pressure in $L_0^2(\Omega)$, it is better to use $(k+1)$-th degree of polynomials for V_h and k-th degree of polynomials for Q_h.

Remark 7.4 (Constraint ratio)**.** An empirical approach has been used to check the balance between velocity and pressure approximation spaces. The so-called constraint ratio is defined as

$$C_r := \mathrm{dim}Q_h/\mathrm{dim}V_h.$$

Apparently, if $C_r > 1$ then number of constraints exceeds the number of variables, which will usually cause locking. On the other hand, if C_r is too small, then divergence free condition is not approximated accurately enough.

The easiest and seemingly natural choice for the mixed finite element spaces is the pair of the lowest order polynomials $P_h^{1,0}$–P_h^0. Unfortunately, this pair does not satisfy the discrete inf-sup condition and we have to either enlarge velocity field finite element space or restrict the pressure space. There are many possible stable pairs; see the survey paper [9] and references therein for more details. Here we just name a few:

- $\left[P_h^{k,0}\right]^d$–$P_h^{k-1,0}$ for $k \geqslant 2$, Taylor–Hood
- $\left[Q_h^{k,0}\right]^d$–$Q_h^{k-1,0}$ for $k \geqslant 2$, Taylor–Hood
- $\left[P_h^{1,0} \oplus \mathsf{B}_\tau^3\right]^2$–$P_h^0$, where B_τ^3 are cubic bubble functions, MINI
- $\left[P_{h/2}^{1,0}\right]^2$–$P_h^0$
- $\left[P_h^{2,0}\right]^d$–P_h^0, important theoretically, but degree not matching
- $\left[P_h^{2,0} \oplus \mathsf{B}_\tau^3\right]^2$–$P^{1,-1}$, Crouzeix–Raviart
- $\left[P_h^{2,0} \oplus \mathsf{B}_\tau^4\right]^3$–$P^{1,-1}$, Crouzeix–Raviart
- $\left[P_h^{1,\mathrm{NC}}\right]^d$–$P_h^0$, non-conforming Crouzeix–Raviart
- $\left[P_h^{k,0}\right]^2$–$P_h^{k-1,-1}$ for $k \geqslant 4$, Scott–Vogelius
- $\left[Q_h^{k,0}\right]^d$–$P_h^{k-1,-1}$ for $k \geqslant 2$

Constructing stable finite difference schemes for the Stokes equation lacks of theoretical guidance like the Babuška–Brezzi condition discussed above. However we can expect that the standard five-point stencil does not work for the Stokes equation. This is because the five-point stencil can be viewed as $Q_h^{1,0} - Q_h^{1,0}$ finite element with a specific quadrature rule. If we change the pressure discretization to the center of cells, then it yields $Q_h^{1,0} - Q_h^{0,-1}$. And, apparently, both finite element pairs are not stable. The main idea of the Marker-and-Cell (MAC) scheme is to place the degrees of freedom for velocity and pressure at different locations. More specifically, the pressure p is defined at the cell centers, the velocity component $\mathbf{u}_1$ is defined at the middle points of vertical edges, and the velocity component $\mathbf{u}_2$ defined at the middle points of horizontal edges; see Figure 7.2. This method is same as the RT_0 finite element on rectangular grids.

Mixed methods for the Poisson's equation ⋆

:MixedPoisson

Mixed finite element methods have been applied to our model problem, the Poisson's equation, as well. In this section, we use this model problem to further explain how to construct preconditioners arising from the saddle-point problems.

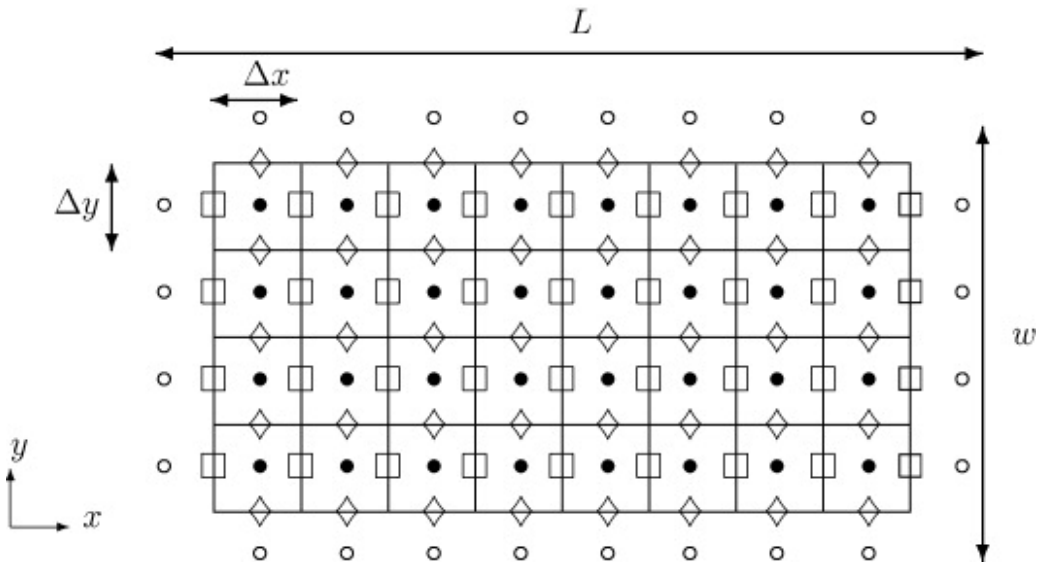


fig:mac

Figure 7.2: A sample discretization using the MAC scheme

① Sometimes the mixed formulation of the Poisson's equation is used for numerical treatment: Find $(\mathbf{u}, p) \in H_0(\mathrm{div}, \Omega) \times L_0^2(\Omega)$ such that

$$\begin{cases} (\mathbf{u}, \mathbf{v}) + (p, \nabla \cdot \mathbf{v}) &= \langle \mathbf{f}, \mathbf{v} \rangle, \quad \forall\, \mathbf{v} \in H_0(\mathrm{div}, \Omega); \\ (\nabla \cdot \mathbf{u}, q) &= \langle g, q \rangle, \quad \forall\, q \in L_0^2(\Omega). \end{cases} \tag{7.38}$$

eqn:PoissonMi

Here $H(\mathrm{div}, \Omega)$ consists of all functions in $[L^2(\Omega)]^d$ with divergence in $L^2(\Omega)$ and $H_0(\mathrm{div}, \Omega)$ contains the $H(\mathrm{div}, \Omega)$-functions with vanishing normal components on the boundary $\partial\Omega$. Define an inner product

$$(\mathbf{u}, \mathbf{v})_{H_0(\mathrm{div},\Omega)} := (\mathbf{u}, \mathbf{v}) + (\nabla \cdot \mathbf{u}, \nabla \cdot \mathbf{v}). \tag{7.39}$$

eqn:Hdiv-inne

This problem corresponds to the mixed formulation of the Poisson's equation with the Neumann boundary condition.

If $\mathbf{u} \in \mathscr{Z}$ is divergence free, then $\|\mathbf{u}\|_{H_0(\mathrm{div},\Omega)} = \|\mathbf{u}\|_{0,\Omega}$. Hence we can easily verify the Brezzi conditions hold for this problem. As a consequence, the operator

$$\tilde{\mathcal{A}}_0 = \begin{pmatrix} \mathcal{I} & -\mathrm{grad} \\ \mathrm{div} & 0 \end{pmatrix} : H_0(\mathrm{div}, \Omega) \times L_0^2(\Omega) \mapsto H_0(\mathrm{div}, \Omega)' \times L_0^2(\Omega)$$

is an isomorphism. The canonical preconditioner is a block diagonal isomorphism

$$\tilde{\mathcal{D}}_0^{(1)} = \begin{pmatrix} (\mathcal{I} - \mathrm{grad}\,\mathrm{div})^{-1} & 0 \\ 0 & \mathcal{I} \end{pmatrix} : H_0(\mathrm{div}, \Omega)' \times L_0^2(\Omega) \mapsto H_0(\mathrm{div}, \Omega) \times L_0^2(\Omega).$$

② There is an alternative mixed formulation for the Poisson's equation: Find $(\mathbf{u}, p) \in [L^2(\Omega)]^d \times \left(H^1(\Omega) \bigcap L_0^2(\Omega)\right)$ such that

$$\left\{ \begin{array}{rcll} (\mathbf{u}, \mathbf{v}) - (\nabla p, \mathbf{v}) & = & \langle \mathbf{f}, \mathbf{v} \rangle, & \forall\, \mathbf{v} \in [L^2(\Omega)]^d; \\ -(\mathbf{u}, \nabla q) & = & \langle g, q \rangle, & \forall\, q \in H^1(\Omega) \bigcap L_0^2(\Omega). \end{array} \right. \tag{7.40}$$

eqn:PoissonMi

The Brezzi conditions can be verified using the Poincaré's inequality. Hence $\tilde{\mathcal{A}}$ is also well-defined on $[L^2(\Omega)]^d \times \left(H^1(\Omega) \bigcap L_0^2(\Omega)\right)$. And in this case, the canonical preconditioner is

$$\tilde{\mathcal{D}}_0^{(2)} = \begin{pmatrix} \mathcal{I} & 0 \\ 0 & (-\Delta)^{-1} \end{pmatrix} : [L^2(\Omega)]^d \times \left(H^1(\Omega) \bigcap L_0^2(\Omega)\right)' \mapsto [L^2(\Omega)]^d \times \left(H^1(\Omega) \bigcap L_0^2(\Omega)\right).$$

Apparently, this preconditioner is significantly different than the one given in the previous subsection. As a result, different choices of approximation space and its norm can yield very different solution methods.

7.4 Canonical preconditioners

In this section, we discuss how to construct canonical preconditioners for the saddle-point problems, like the Stokes equation and the time-dependent Stokes equation. The basic idea follows the discussion in §2.2.

Preconditioning the Stokes problem

We notice that the corresponding operator of the Stokes system

$$\tilde{\mathcal{A}} := \begin{pmatrix} -\Delta & -\operatorname{grad} \\ \operatorname{div} & 0 \end{pmatrix}$$

is an isomorphism mapping from $[H_0^1(\Omega)]^d \times L_0^2(\Omega)$ onto $[H^{-1}(\Omega)]^d \times L_0^2(\Omega)$. A natural preconditioner would be the classical block diagonal preconditioner

$$\tilde{\mathcal{D}} = \begin{pmatrix} (-\Delta)^{-1} & 0 \\ 0 & \mathcal{I} \end{pmatrix}.$$

This observation immediately motivates the classical block diagonal preconditioner [13].

Similar to the continuous case, we can construct natural preconditioners based on the mapping properties. Let $\{X_h\}$ be a family of finite element spaces and it is conforming in the sense that $X_h \subset \mathscr{X} := [H_0^1(\Omega)]^d \times L_0^2(\Omega)$. Consider the discrete Stokes problem: Find $(\mathbf{u}_h, p_h) \in X_h$ such that

$$\tilde{a}[(\mathbf{u}_h, p_h), (\mathbf{v}_h, q_h)] = \langle f, \mathbf{v}_h \rangle, \quad \forall (\mathbf{v}_h, q_h) \in X_h.$$

The corresponding linear map $\tilde{\mathcal{A}}_h : X_h \mapsto X_h'$ is given by

$$\langle \tilde{\mathcal{A}}_h x, y\rangle = \tilde{a}[x, y], \quad \forall x, y \in X_h.$$

Note that, in this case, $\tilde{a}$ is not positive definite and the system $\tilde{\mathcal{A}}_h$ can be singular.

According to Remark 7.2, the stable discretizations can be characterized by a discrete inf-sup condition: There exists a constant α_0, independent of h, such that

$$\inf_{x\in X_h} \sup_{y\in X_h} \frac{\tilde{a}[x, y]}{\|x\|_{\mathscr{X}}\|y\|_{\mathscr{X}}} \geqslant \alpha_0 > 0. \tag{7.41}$$ eqn:dis-mixed

This condition does not follow from the corresponding continuous inf-sup condition. Similar to the continuous case, we can define a preconditioner $\tilde{\mathcal{D}}_h : X_h' \mapsto X_h$ by

$$(\tilde{\mathcal{D}}_h f, y)_{\mathscr{X}} = \langle f, y\rangle, \quad \forall y \in X_h.$$

That is to say

$$\tilde{\mathcal{D}}_h := \begin{pmatrix} (-\Delta_h)^{-1} & 0 \\ 0 & \mathcal{I}_h^{-1} \end{pmatrix}. \tag{7.42}$$ eqn:Stokes-pr

Apparently, if $\tilde{\mathcal{A}}_h$ is symmetric, $\tilde{\mathcal{D}}_h\tilde{\mathcal{A}}_h$ is symmetric with respect to $(\cdot,\cdot)_{\mathscr{X}}$ and

$$\|\tilde{\mathcal{D}}_h\tilde{\mathcal{A}}_h\|_{\mathscr{L}(X_h;X_h)} \leqslant C_a, \qquad \|(\tilde{\mathcal{D}}_h\tilde{\mathcal{A}}_h)^{-1}\|_{\mathscr{L}(X_h;X_h)} \leqslant \alpha_0^{-1}.$$

Hence the condition number $\kappa(\tilde{\mathcal{D}}_h\tilde{\mathcal{A}}_h)$ is uniformly bounded.

Preconditioning the time-dependent Stokes problem ⋆

We now consider preconditioner for the time-dependent Stokes problem (7.24) where the coefficient operator is defined as

$$\tilde{\mathcal{A}}_\epsilon := \begin{pmatrix} \mathcal{I} - \epsilon^2\Delta & -\operatorname{grad} \\ \operatorname{div} & 0 \end{pmatrix}$$

For this problem, we shall construct a preconditioner which is uniformly convergent with respect to both h and ϵ.

We first consider how to construct a preconditioner for $\mathcal{A}_\epsilon = \mathcal{I} - \epsilon^2\Delta$ corresponding to the reaction-diffusion equation. In order to study the reaction-diffusion equation $\mathcal{A}_\epsilon\mathbf{u} = \mathbf{f}$ in Ω and $\mathbf{u}|_{\partial\Omega} = 0$, we can easily see the natural norm is

$$\|\mathbf{u}\|_{L^2\bigcap \epsilon H_0^1} := \left(\|\mathbf{u}\|_0^2 + \epsilon^2\|\nabla\mathbf{u}\|_0^2\right)^{\frac{1}{2}} = \left(\|\mathbf{u}\|_0^2 + \epsilon^2\sum_{i=1}^{d}\|\nabla\mathbf{u}_i\|_0^2\right)^{\frac{1}{2}}.$$

The question is what would be the appropriate norm for $\mathbf{f}$.

Using the classical theory of intersections and sums of Hilbert spaces [8], we can introduce the norms for $\mathscr{X}_1 \bigcap \mathscr{X}_2$ and $\mathscr{X}_1 + \mathscr{X}_2$ as

$$\|u\|_{\mathscr{X}_1 \bigcap \mathscr{X}_2} := \left(\|u\|^2_{\mathscr{X}_1} + \|u\|^2_{\mathscr{X}_2}\right)^{\frac{1}{2}}$$

and

$$\|u\|_{\mathscr{X}_1+\mathscr{X}_2} := \inf_{\substack{u=u_1+u_2 \\ u_1\in\mathscr{X}_1, u_2\in\mathscr{X}_2}} \left(\|u_1\|^2_{\mathscr{X}_1} + \|u_2\|^2_{\mathscr{X}_2}\right)^{\frac{1}{2}}.$$

If $\mathscr{X}_1 \bigcap \mathscr{X}_2$ is dense in both $\mathscr{X}_1$ and $\mathscr{X}_2$, then

$$(\mathscr{X}_1 \bigcap \mathscr{X}_2)' = \mathscr{X}_1' + \mathscr{X}_2' \qquad \text{and} \qquad (\mathscr{X}_1 + \mathscr{X}_2)' = \mathscr{X}_1' \bigcap \mathscr{X}_2'.$$

If $\mathcal{F} \in \mathscr{L}(\mathscr{X}_1; \mathscr{Y}_1) \bigcap \mathscr{L}(\mathscr{X}_2; \mathscr{Y}_2)$, then

$$\mathcal{F} \in \mathscr{L}(\mathscr{X}_1 \bigcap \mathscr{X}_2; \mathscr{Y}_1 \bigcap \mathscr{Y}_2) \bigcap \mathscr{L}(\mathscr{X}_1 + \mathscr{X}_2; \mathscr{Y}_1 + \mathscr{Y}_2).$$

For our purpose, we assume that $\mathscr{X}_1$ and $\mathscr{X}_2$ are real separable Hilbert spaces and $\mathscr{X}_2 \subset \mathscr{X}_1$. Hence it is natural to assume $\|u\|_{\mathscr{X}_1} \leqslant \|u\|_{\mathscr{X}_2}$. For $\epsilon > 0$, we consider the norm for spaces $\mathscr{X}_1 \bigcap \epsilon\mathscr{X}_2$ by

$$\|u\|_{\mathscr{X}_1 \bigcap \epsilon\mathscr{X}_2} := \left(\|u\|^2_{\mathscr{X}_1} + \epsilon^2\|u\|^2_{\mathscr{X}_2}\right)^{\frac{1}{2}}, \qquad \|f\|_{\mathscr{X}_1'+\epsilon^{-1}\mathscr{X}_2'} := \inf_{\substack{f=f_1+f_2 \\ f_1\in\mathscr{X}_1', f_2\in\mathscr{X}_2'}} \left(\|f_1\|^2_{\mathscr{X}_1'} + \epsilon^{-2}\|f_2\|^2_{\mathscr{X}_2'}\right)^{\frac{1}{2}}.$$

Apparently, $\mathscr{X}_\epsilon := \mathscr{X}_1 \bigcap \epsilon\mathscr{X}_2$ and $\mathscr{X}_\epsilon' := \mathscr{X}_1 + \epsilon^{-1}\mathscr{X}_2$ are the same as $\mathscr{X}_1$ as sets. As ϵ tends to zero, the norms for $\mathscr{X}_\epsilon$ and $\mathscr{X}_\epsilon'$ approaches the norms $\|\cdot\|_{\mathscr{X}_1}$ and $\|\cdot\|_{\mathscr{X}_1'}$, respectively. In particular, for the reaction-diffusion problem, we have

$$\mathscr{X}_\epsilon = L^2(\Omega) \bigcap \epsilon H^1_0(\Omega) \qquad \text{and} \qquad \mathscr{X}_\epsilon' = L^2(\Omega) + \epsilon^{-1}H^{-1}(\Omega).$$

As ϵ goes to zero, both norms approaches the L^2-norm. Furthermore, $\|f\|^2_{\mathscr{X}_\epsilon'}$ is equivalent to $\langle f, (\mathcal{I} - \epsilon^2\Delta)^{-1}f\rangle = \langle (\mathcal{I} - \epsilon^2\Delta)u, u\rangle$.

Now we are in position to develop preconditioners for the time-dependent Stokes problem (7.23):

① In view of §7.3, we know that $\tilde{\mathcal{A}}_0$ is bounded from $H_0(\mathrm{div}, \Omega) \times L^2_0(\Omega)$ into its dual space. Hence we consider the operator $\tilde{\mathcal{A}}_\epsilon$ on

$$\mathscr{X}_\epsilon := \left(H_0(\mathrm{div}, \Omega) \bigcap \epsilon[H^1_0(\Omega)]^d\right) \times L^2_0(\Omega) \quad \text{and} \quad \mathscr{X}_\epsilon' := \left(H_0(\mathrm{div}, \Omega)' + \epsilon^{-1}[H^{-1}(\Omega)]^d\right) \times L^2_0(\Omega).$$

In this case, the two Brezzi conditions holds and $\tilde{\mathcal{A}}_\epsilon$ is an isomorphism. In turn, the canonical preconditioner is of the form

$$\tilde{\mathcal{D}}^{(1)}_\epsilon = \begin{pmatrix} (\mathcal{I} - \mathrm{grad}\,\mathrm{div} - \epsilon^2\Delta)^{-1} & 0 \\ 0 & \mathcal{I} \end{pmatrix}.$$

② We have seen that $\tilde{\mathcal{A}}_0$ is also bounded on $[L^2(\Omega)]^d \times \left(H^1(\Omega) \bigcap L_0^2(\Omega)\right)$ into its dual space. Furthermore, in order to guarantee the inf-sup condition, the proper norm for the pressure unknown is [43, 44]:

$$\sup_{\mathbf{v}\in[H_0^1(\Omega)]^d} \frac{(q, \nabla\cdot\mathbf{v})}{\|\mathbf{v}\|_{L^2\bigcap \epsilon H^1}} = \|\nabla q\|_{L^2+\epsilon^{-1}H^{-1}} \sim \|q\|_{H^1+\epsilon^{-1}L^2}.$$

Motivated by these observations, we can consider

$$\mathscr{X}_\epsilon := \left[L^2(\Omega)\bigcap \epsilon H_0^1(\Omega)\right]^d \times \left(H^1(\Omega)\bigcap L_0^2(\Omega) + \epsilon^{-1}L_0^2(\Omega)\right)$$

and

$$\mathscr{X}_\epsilon' := \left[L^2(\Omega) + \epsilon^{-1}H^{-1}(\Omega)\right]^d \times \left((H^1(\Omega)\bigcap L_0^2(\Omega))'\bigcap \epsilon L_0^2(\Omega)\right).$$

This choice of spaces gives a preconditioner of the form

$$\tilde{\mathcal{D}}_\epsilon^{(2)} = \begin{pmatrix} (\mathcal{I}-\epsilon^2\Delta)^{-1} & 0 \\ 0 & (-\Delta)^{-1}+\epsilon^2\mathcal{I} \end{pmatrix}.$$

Along this line, we can construct discrete block diagonal preconditioners for the time-dependent Stokes problem [29, 12].

7.5 Block preconditioners

In the previous section, we discussed how to construct canonical (natural) preconditioners based on the mapping property of the continuous Stokes equation. Now we shall consider the discrete Stokes problem arising in the mixed finite element method (such as the Taylor–Hood finite element method) in algebraic setting, i.e.,

$$\tilde{A}\begin{pmatrix} u \\ p \end{pmatrix} = \begin{pmatrix} f \\ g \end{pmatrix} \quad \text{and} \quad \tilde{A} := \begin{pmatrix} A & B^T \\ B & 0 \end{pmatrix}. \tag{7.43}$$ eqn:dis-saddl

Suppose $A \in \mathbb{R}^{n\times n}$, $B \in \mathbb{R}^{m\times n}$, $u \in \mathbb{R}^n$, and $p \in \mathbb{R}^n$. Let $N = n + m$. Assume that A is SPD and B has full rank. It is well-known that the coupled system $\tilde{A}$ is symmetric, indefinite, and non-singular.

Block diagonal and lower triangular method

If we consider the block diagonal preconditioner given in the previous section, the preconditioner can be written as

$$\tilde{D} := \begin{pmatrix} A^{-1} & 0 \\ 0 & M_p^{-1} \end{pmatrix}, \tag{7.44}$$ eqn:diag-prec

where M_p is the mass matrix corresponding to the pressure approximation space and, hence, it is well-conditioned; see Remark 3.7. It is easy to check that (7.44) is exactly the algebraic form of (7.42). Because both A and M_p are symmetric positive definite matrices, the preconditioner is well-defined.

factorization **Remark 7.5** (Block factorizations)**.** We can apply the following block factorizations to the matrix $\tilde{A}$ such that

$$\begin{pmatrix} A & B^T \\ B & 0 \end{pmatrix} = \begin{pmatrix} I_u & 0 \\ BA^{-1} & I_p \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & S \end{pmatrix} \begin{pmatrix} I_u & A^{-1}B^T \\ 0 & -I_p \end{pmatrix}$$
$$= \begin{pmatrix} A & 0 \\ B & S \end{pmatrix} \begin{pmatrix} I_u & A^{-1}B^T \\ 0 & -I_p \end{pmatrix} = \begin{pmatrix} I_u & 0 \\ BA^{-1} & -I_p \end{pmatrix} \begin{pmatrix} A & B^T \\ 0 & S \end{pmatrix},$$

where the matrix $S := BA^{-1}B^T$ is the Schur complement. In fact, $\tilde{D}$ in (7.44) can be viewed as an approximation of $\mathrm{diag}(A^{-1}, S^{-1})$.

rem:Schur **Remark 7.6** (Schur complement)**.** Since the A is SPD, the Schur complement $S = BA^{-1}B^T$ is symmetric and positive semi-definite. Moreover, if B has full rank, S is also SPD and we can apply the CG method to solve the Schur complement equation. However, generally speaking, $S^{-1}p$ cannot be computed efficiently with acceptable computational cost. Hence the Schur complement S should be approximated by some approximation $\hat{S}$. There are many different ways based on approximation of the Schur complement; see the survey paper [6].

We can also use the block lower triangular matrix to construct a preconditioner

$$\tilde{T} := \begin{pmatrix} A & 0 \\ B & \hat{S} \end{pmatrix}^{-1}. \tag{7.45}$$ eqn:lowertrig

In particular, if we replace A by its diagonal part D in the LU decomposition of Remark 7.5, then we get the so-called SIMPLE preconditioner

$$\tilde{T}_{\mathrm{SIMPLE}} := \begin{pmatrix} I_u & D^{-1}B^T \\ 0 & -I_p \end{pmatrix}^{-1} \begin{pmatrix} A & 0 \\ B & BD^{-1}B^T \end{pmatrix}^{-1}. \tag{7.46}$$ eqn:simple-pr

The name comes from the widely-used SIMPLE method for fluid problems.

Augmented Lagrangian method

One of the most well-known iterative method for solving (7.43) is probably the Uzawa method. As the last decomposition in Remark 7.5, we can factorize the coefficient matrix as

$$\begin{pmatrix} A & B^T \\ B & 0 \end{pmatrix} = \begin{pmatrix} I_u & 0 \\ BA^{-1} & -I_p \end{pmatrix} \begin{pmatrix} A & B^T \\ 0 & S \end{pmatrix}.$$

This means the original linear system can be rewritten as

$$\begin{pmatrix} A & B^T \\ 0 & S \end{pmatrix} \begin{pmatrix} u \\ p \end{pmatrix} = \begin{pmatrix} f \\ BA^{-1}f - g \end{pmatrix}.$$

As discussed in Remark 7.6, the pressure Schur complement equation might be too expensive to be solved exactly. We can apply an iterative method to solve it. For example, we can apply the Richardson's iteration for the second equation in the above system, i.e.,

$$p^{\text{new}} = p^{\text{old}} + \omega\Big(BA^{-1}f - g - Sp^{\text{old}}\Big) = p^{\text{old}} - \omega\Big(g - BA^{-1}f + BA^{-1}B^Tp^{\text{old}}\Big).$$

Hence we can write the above iteration as an alternative direction method

$$Au^{\text{new}} = f - B^Tp^{\text{old}}, \qquad p^{\text{new}} = p^{\text{old}} - \omega(g - Bu^{\text{new}}). \tag{7.47}$$ eqn:uzawa

The method (7.47) is called the Uzawa iteration and it is just the Richardson iteration for the Schur complement equation. As we have discussed in §2.1, the method converges with an appropriate scaling factor ω but the convergence rate is usually very slow. One way to speedup the convergence is to apply the Augmented Lagrangian method (cf., for example, [32]):

$$(A + \epsilon^{-1}B^TB)u^{\text{new}} = f + \epsilon^{-1}B^Tg - B^Tp^{\text{old}}, \qquad p^{\text{new}} = p^{\text{old}} - \epsilon^{-1}(g - Bu^{\text{new}}). \tag{7.48}$$ eqn:AL

rem:UzawaAL **Remark 7.7** (Uzawa method and Augmented Lagrangian method)**.** It is easy to see that the Augmented Lagrangian (AL) method is just the Uzawa method for the modified equation

$$\tilde{A}_\epsilon \begin{pmatrix} u \\ p \end{pmatrix} = \begin{pmatrix} f + \epsilon^{-1}B^Tg \\ g \end{pmatrix}, \qquad \text{where } \tilde{A}_\epsilon := \begin{pmatrix} A + \epsilon^{-1}B^TB & B^T \\ B & 0 \end{pmatrix}. \tag{7.49}$$ eqn:dis-saddl

Furthermore, the damping factor ω is chosen to be ϵ^{-1}.

thm:AL **Theorem 7.4** (Convergence rate of Augmented Lagrangian method)**.** *Let $(u^{(0)}, p^{(0)})$ be a given initial guess and $(u^{(m)}, p^{(m)})$ be the iterates obtained via the Augmented Lagrangian method (7.48). Then we have*

$$\begin{aligned} \|p - p^{(m)}\|_0 &\leqslant \Big(\frac{\epsilon}{\epsilon + \lambda_1}\Big)^m \|p - p^{(0)}\|_0, \\ \|u - u^{(m)}\|_A &\leqslant \sqrt{\epsilon}\|p - p^{(m-1)}\|_0 \leqslant \sqrt{\epsilon}\Big(\frac{\epsilon}{\epsilon + \lambda_1}\Big)^{m-1} \|p - p^{(0)}\|_0, \end{aligned}$$

where λ_1 is the minimal eigenvalue of $S = BA^{-1}B^T$.

Sketch of proof. From (7.48) and (7.49), we have

$$(A + \epsilon^{-1}B^TB)\big(u - u^{(m)}\big) = -B^T\big(p - p^{(m-1)}\big)$$

and

$$p - p^{(m)} = \Big(I - B(\epsilon A + B^T B)^{-1} B^T\Big)\big(p - p^{(m-1)}\big).$$

By the Shermann–Morrison–Woodburry formula, we have

$$Z := B(\epsilon A + B^T B)^{-1} B^T = S_\epsilon - S_\epsilon (I + S_\epsilon)^{-1} S_\epsilon, \qquad S_\epsilon := \epsilon^{-1} B A^{-1} B^T.$$

It is easy to verify that

$$I - B(\epsilon A + B^T B)^{-1} B^T = I - S_\epsilon + S_\epsilon (I + S_\epsilon)^{-1} S_\epsilon = (I + S_\epsilon)^{-1}.$$

The above equality shows $\rho(Z) \leqslant 1$ and $p - p^{(m)} = (I + S_\epsilon)^{-1}\big(p - p^{(m-1)}\big)$. So the first estimate follows immediately. The second estimate is obtained by observing

$$\big\|u - u^{(m)}\big\|_A^2 = \Big((A + \epsilon^{-1} B^T B - \epsilon^{-1} B^T B)\big(u - u^{(m)}\big), u - u^{(m)}\Big) \leqslant \epsilon\big(Z(p - p^{(m-1)}), p - p^{(m-1)}\big)$$

and then applying the first estimate. □

According to Theorem 7.4, we can make the convergence as fast as we want by adjusting the parameter ϵ. However, the price to pay is that, in each iteration, we have to solve a nearly-singular system with coefficient matrix $A + \epsilon^{-1} B^T B$, which was discussed in [41]. We can also apply the Augmented Lagrangian method as a preconditioner

$$\tilde{T}_{\mathrm{AL}} := \begin{pmatrix} A + \epsilon^{-1} B^T B & 0 \\ B & \epsilon I \end{pmatrix}^{-1}, \tag{7.50}$$

eqn:AL-precon

which is often referred to as the AL preconditioner [7].

The method is closely related to the grad-div stabilization [22] of the Stokes (or Navier–Stokes) problem:

$$\begin{cases} (I - \mu\Delta)\mathbf{u} - \epsilon^{-1}\nabla\nabla\cdot\mathbf{u} + \nabla p &= \mathbf{f}, \quad \Omega; \\ \nabla\cdot\mathbf{u} &= 0, \quad \Omega; \\ \mathbf{u} &= 0, \quad \partial\Omega. \end{cases} \tag{7.51}$$

eqn:Stokes2

In this modified problem, the coercivity condition automatically holds on the discrete level for the $H_0(\mathrm{div})$-norm defined by (7.39). After discrezation by some mixed finite element method, we obtain discrete systems in the form of (7.48). We can apply the block preconditioners discussed in the previous subsection to solve the resulting discrete problems; see the survey and numerical experiments by He and Vuik [36].

7.6 Multigrid methods for Stokes equation

We can construct coupled multigrid methods for the saddle-point problem (7.43) as well. For the transfer operators, by applying the similar ideas as in multigrid methods for scalar equations, we can construct prolongations and restrictions for velocity and pressure variables separately. Coarse level solvers can also apply the same nested iterations as in §6.2. So we only discuss smoothers for the Stokes system. Analysis and numerical experiments using different smoothers have been reviewed in the survey by Larin and Reusken [40]. Apparently, the block preconditioners discussed in the previous section can also be applied as smoothers for coupled multigrid methods. In this section, we discuss two other widely-used smoothers in practice.

Braess–Sarazin smoother

The Braess–Sarazin smoother is introduced in [11] and can be written as

$$\begin{pmatrix} u^{(m+1)} \\ p^{(m+1)} \end{pmatrix} = \begin{pmatrix} u^{(m)} \\ p^{(m)} \end{pmatrix} + \begin{pmatrix} \omega D & B^T \\ B & 0 \end{pmatrix}^{-1} \left[\begin{pmatrix} f \\ 0 \end{pmatrix} - \begin{pmatrix} A & B^T \\ B & 0 \end{pmatrix} \begin{pmatrix} u^{(m)} \\ p^{(m)} \end{pmatrix} \right], \tag{7.52}$$

eqn:Braess

where ω is a positive parameter. This method mimics the damped Jacobi smoother for the Poisson's equation.

We need to solve, in each smoothing step, the following the linear system

$$\begin{pmatrix} \omega D & B^T \\ B & 0 \end{pmatrix} \begin{pmatrix} \delta u^{(m)} \\ \delta p^{(m)} \end{pmatrix} = \begin{pmatrix} f - Au^{(m)} - B^T p^{(m)} \\ -Bu^{(m)} \end{pmatrix}.$$

The second equation ensures the discrete divergence free condition, i.e.,

$$Bu^{(m+1)} = B\big(u^{(m)} + \delta u^{(m)}\big) = 0, \qquad m = 1, 2, \ldots$$

Apparently, the Braess–Sarazin smoother can be reduced to an auxiliary pressure equation

$$\big(BD^{-1}B^T\big)\,\delta p^{(m)} = \omega Bu^{(m)} + BD^{-1}\big(f - Au^{(m)} - B^T p^{(m)}\big).$$

The coefficient matrix $\hat{S} := BD^{-1}B^T$ is similar to a scaled discrete Laplace operator on the pressure space. In practice, we can solve it approximately using an iterative method.

Vanka smoother

Next we introduce a smoother originally proposed by Vanka [54]. In the context of finite element methods, the Vanka-type smoothers are just block Gauss–Seidel (or Jacobi) methods. Each block contains degrees of freedom in an element or a set of elements. One of the popular variant of

Vanka-type smoothers is the so-called *pressure-oriented Vanka smoother* for continuous pressure approximations. We only discuss this special case of Vanka smoother here.

For each pressure variable indexed by i ($1 \leqslant i \leqslant m$), let the set of velocity indices that are "connected" to i as

$$\mathsf{S}_i := \{1 \leqslant j \leqslant n \,:\, b_{i,j} \neq 0\},$$

where $b_{i,j}$ is the (i,j)-entry of the matrix B. So we can define an injection to the set of variables $\{u_j \,(j \in \mathsf{S}_i),\, p_i\}$, i.e.,

$$I_i = \begin{pmatrix} I_{u,i} & 0 \\ 0 & I_{p,i} \end{pmatrix} \in \mathbb{R}^{(|\mathsf{S}_i|+1)\times(n+m)},$$

where $I_{p,i}p = p_i$ and $I_{u,i}u = (u_j)_{j\in\mathsf{S}_i}$ are the corresponding injection matrices for velocity and pressure, respectively.

We can then apply a multiplicative Schwarz method (or the so-called Full Vanka smoother):

$$I - \tilde{T}_{\text{FVanka}}\tilde{A} = \prod_{i=1}^{m}\left(I - I_i^T \tilde{A}_i^{-1} I_i \tilde{A}\right), \tag{7.53}$$

eqn:FVanka

where

$$\tilde{A}_i = I_i \tilde{A} I_i^T = \begin{pmatrix} A_i & B_i^T \\ B_i & 0 \end{pmatrix} \in \mathbb{R}^{(|\mathsf{S}_i|+1)\times(|\mathsf{S}_i|+1)}.$$

We can also use a simplified version (i.e., the Diagonal Vanka smoother):

$$I - \tilde{T}_{\text{DVanka}}\tilde{A} = \prod_{i=1}^{m}\left(I - I_i^T \tilde{D}_i^{-1} I_i \tilde{A}\right), \tag{7.54}$$

eqn:DVanka

where

$$\tilde{D}_i = \begin{pmatrix} D_i & B_i^T \\ B_i & 0 \end{pmatrix} \in \mathbb{R}^{(|\mathsf{S}_i|+1)\times(|\mathsf{S}_i|+1)}.$$

In this case, due to the special nonzero pattern of $\tilde{D}_i$, it can be solved very efficiently.

7.7 Homework problems

hw:div-eps **HW 7.1.** Show the equation (7.16). Hint: In $\mathbb{R}^2$, taking divergence of the symmetric gradient, we get

$$\begin{aligned} \nabla\cdot\boldsymbol{\varepsilon}(\mathbf{u}) &= \begin{pmatrix} \partial_1^2 u_1 + \frac{1}{2}\partial_2(\partial_2 u_1 + \partial_1 u_2) \\ \partial_2^2 u_2 + \frac{1}{2}\partial_1(\partial_1 u_2 + \partial_2 u_1) \end{pmatrix} \\ &= \begin{pmatrix} \frac{1}{2}(\partial_1^2 u_1 + \partial_2^2 u_1) + \frac{1}{2}\partial_1(\partial_1 u_1 + \partial_2 u_2) \\ \frac{1}{2}(\partial_1^2 u_2 + \partial_2^2 u_2) + \frac{1}{2}\partial_2(\partial_1 u_1 + \partial_2 u_2) \end{pmatrix} = \frac{1}{2}\Delta\mathbf{u} + \frac{1}{2}\nabla\nabla\cdot\mathbf{u}. \end{aligned}$$

hw:Stokes_weak **HW 7.2.** Derive the weak form (7.26) of the Stokes equations (7.25).

hw:AL **HW 7.3.** Give the complete proof of Theorem 7.4.

Chapter 8

Optimization Problems

Mathematical optimization (mathematical programming or optimization) is the selection of a "best" element (with regard to certain criterion) from some set of available alternatives. Many optimization problems can be written as variational inequalities (VIs); for example, many problems in economics, operations research, and transportation equilibrium problems. In this chapter, we discuss multilevel iterative methods for solving finite-dimensional variational inequalities.

8.1 Model problems

sec:vi

VIs arise from a wide range of application areas, like mechanics, control theory, engineering, and finance. After several decades of development, this subject has become very rich on both theory and numerics. For a general discussion on the existence and regularity, we refer the interested readers to [38]. For a comprehensive discussion on numerical methods for VIs, we refer to Glowinski [34].

A model variational inequality

Let $a[\cdot,\cdot]$ and $f(\cdot)$ be a symmetric bilinear form and a linear form, respectively, and $\chi \in H_0^1(\Omega)$ be an admissible obstacle (for simplicity, we assume the zero boundary condition). Consider the following elliptic variational inequality (or the obstacle problem): Find $u \in \mathcal{K}_\chi := \{v \in H_0^1(\Omega) : v \geqslant \chi\}$, such that

$$a[u, v-u] \geqslant f(v-u), \quad \forall\ v \in \mathcal{K}_\chi. \tag{8.1}$$

eqn:evi1

After transformation $w := u - \chi$, we arrive at a new problem with a simple inequality constraint: Find $w \in \mathcal{K}_0 := \{v \in H_0^1(\Omega) : v \geqslant 0\}$, such that

$$a[w, v-w] \geqslant f_0(v-w) := f(v-w) - a[\chi, v-w], \quad \forall\ v \in \mathcal{K}_0. \tag{8.2}$$

eqn:evi2

For problem (8.1), the *Lagrange multiplier* can be defined as σ_1 such that

$$\langle \sigma_1(u), \varphi \rangle := f(\varphi) - a[u, \varphi], \quad \forall\ \varphi \in H_0^1(\Omega). \tag{8.3}$$ eqn:sigma

On the other hand, for (8.2), notice, for any $\varphi \in H_0^1(\Omega)$, that

$$\langle \sigma_2(w), \varphi \rangle = f_0(\varphi) - a[w, \varphi] = f(\varphi) - a[u, \varphi] = \langle \sigma_1(u), \varphi \rangle .$$

It is easy to see that

$$\langle \sigma_1(u), v - u \rangle \leqslant 0, \quad \forall\ v \in \mathcal{K}_\chi, \tag{8.4}$$ eqn:nonpos

or

$$\langle \sigma_2(w), v - w \rangle \leqslant 0, \quad \forall\ v \in \mathcal{K}_0.$$

On the other hand, if σ is the Lagrange multiplier of (8.1), we have

$$\langle \sigma(v) - \sigma(u), \varphi \rangle = -a[v - u, \varphi], \quad \forall\ \varphi \in H_0^1(\Omega).$$

Hence,

$$\langle \sigma(v) - \sigma(u), v - u \rangle = -a[v - u, v - u] = -\left\| v - u \right\|^2, \quad \forall\ v, u \in H_0^1(\Omega). \tag{8.5}$$ eqn:mono

Hence, we have $\langle \sigma(v) - \sigma(u), v - u \rangle \leqslant 0$, for any $v, u \in H_0^1(\Omega)$, i.e., σ is a *monotone* operator.

Remark 8.1 (Uniqueness of solution)**.** Notice that if both u_1 and u_2 are solutions of the variational inequality (8.1), by the monotonicity of σ, $\left\| u_1 - u_2 \right\| = 0$ and then we obtain the uniqueness.

As before, we assume that $\mathcal{A} : H_0^1(\Omega) \mapsto H^{-1}(\Omega)$ be the operator corresponding to $a[\cdot, \cdot]$. An frequently equivalent formulation of (8.1) is the so-called linear complementarity problem (LCP): Find a solution $u \in H_0^1(\Omega)$ such that

$$\begin{cases} \mathcal{A}u - f \geqslant 0 \\ u - \chi \geqslant 0 \\ \langle \mathcal{A}u - f, u - \chi \rangle = 0. \end{cases} \tag{8.6}$$ eqn:lcp

The last equation is the so-called *complementarity condition.*

Proof. If u is a solution of LCP (8.6), then for any $v \in H_0^1(\Omega)$ and $v \geqslant \chi$ we have

$$\langle \mathcal{A}u - f, u - v \rangle = \langle \mathcal{A}u - f, \chi - v \rangle \leqslant 0,$$

in view of the complementarity condition and the sign condition of $\mathcal{A}u - f$. On the other hand, if u is solution of (8.1), it is trivial to see that u satisfies the first two conditions of LCP. The complementarity condition is obtained by taking $v = u + (u - \chi)$ and $v = \chi$. □

Finite element discretization for VIs

As discussed in §3.1, the domain Ω is partitioned into a quasi-uniform simplexes of size h; this mesh is denoted by $\mathcal{M}_h$. Let $V_h \subset W_0^{1,\infty}(\Omega)$ be the continuous piecewise linear finite element space associated with $\mathcal{M}_h$. The obstacle problem (8.2) can be approximated by a finite element function $u_h \in \mathcal{K}_0 \bigcap V_h$ satisfying:

$$a[u_h, v_h - u_h] \geqslant f_0(v_h - u_h), \quad \forall\ v_h \in \mathcal{K}_0 \bigcap V_h. \tag{8.7}$$ eqn:fem2

As before, we denote all the interior nodes of the partition $\mathcal{M}_h$ by $\mathring{G}(\mathcal{M}_h)$. Let $\{\phi_z\}_{z\in\mathring{G}(\mathcal{M}_h)}$ be the canonical linear finite element basis of the mesh $\mathcal{M}_h$. Let $u = u_h := \sum_{z\in\mathring{G}(\mathcal{M}_h)} u_z\phi_z$ and $\underline{u} = (u_z)_{z\in\mathring{G}(\mathcal{M}_h)}$, the discrete solution and its nodal value vector (primal vector form), respectively. Hence we have the following linear system

$$(\underline{v} - \underline{u})^T(A\underline{u} - \vec{f}_0) \geqslant 0, \quad \forall\ \underline{v} \geqslant 0, \tag{8.8}$$ eqn:dis2

where A is the corresponding stiffness matrix of the bilinear form and $\vec{f}_0$ is the dual vector form of f_0.

Remark 8.2. One can prove (see for example [18]) that the l^2-error between the exact solution $\underline{u}$ of (8.8) and any approximation solution $\underline{v}$ satisfies that

$$\|\underline{v} - \underline{u}\|_0 \lesssim \|(\vec{f}_0 - A\underline{v})_+\|_0,$$

where the vector $(\vec{f}_0 - A\underline{v})_+$ is defined element-wise by

$$(\vec{f}_0 - A\underline{v})_{+,i} = \begin{cases} (\vec{f}_0 - A\underline{v})_i & \text{if } \underline{v}_i > 0 \\ \min\{(\vec{f}_0 - A\underline{v})_i, 0\} & \text{if } \underline{v}_i = 0. \end{cases}$$

Error and residual

ssc:residual

As usual, we define the energy functional as following

$$\mathcal{F}(v) := \frac{1}{2}a[v, v] - f(v).$$

Then it follows that

$$\mathcal{F}(v) - \mathcal{F}(u) = \frac{1}{2}\left|\!\left|\!\left| v - u \right|\!\right|\!\right|^2 - \langle \sigma, v - u\rangle, \quad \forall\ v \in \mathcal{K}_\chi. \tag{8.9}$$ eqn:Idiff

Consider finite element solutions, u_h and w_h for problems (8.1) and (8.2), respectively. The differences, in terms of energy, between the finite element solutions and the exact solutions can be written as

$$\begin{aligned} \mathcal{F}(u_h) - \mathcal{F}(u) &= \frac{1}{2}\left|\!\left|\!\left| u_h - u \right|\!\right|\!\right|^2 - \langle \sigma, u_h - u\rangle \\ \mathcal{F}(w_h) - \mathcal{F}(w) &= \frac{1}{2}\left|\!\left|\!\left| w_h - w \right|\!\right|\!\right|^2 - \langle \sigma, w_h - w\rangle. \end{aligned} \tag{8.10}$$ eqn:diff

It is easy to see that the variational inequality (8.2) can be written as the following quadratic minimization problem:

$$\min_{w\in\mathcal{K}_0} \frac{1}{2}a[w,w] - f_0(w). \tag{8.11}$$ eqn:min2

For finite element approximation, we compute the finite dimensional minimization problem

$$\min_{w_h\in V_h\bigcap\mathcal{K}_0} \frac{1}{2}a[w_h,w_h] - f_0(w_h). \tag{8.12}$$ eqn:femmin2

Suppose $\hat{w}_h$ is an approximate solution of the above minimization problem. Then the defect $e_h := w_h - \hat{w}_h$ satisfies

$$\min_{\hat{w}_h+e_h\in V_h\bigcap\mathcal{K}_0} \frac{1}{2}a[\hat{w}_h+e_h,\hat{w}_h+e_h] - f_0(\hat{w}_h+e_h) = \frac{1}{2}a[e_h,e_h] - f_0(e_h) + a[\hat{w}_h,e_h] + C,$$

i.e.,

$$\min_{\hat{w}_h+e_h\in V_h\bigcap\mathcal{K}_0} \frac{1}{2}a[e_h,e_h] - \langle\sigma(\hat{w}_h),e_h\rangle\,. \tag{8.13}$$ eqn:ErrVI

Notice that it is in the same form as (8.12) but replacing f_0 by $\sigma(\hat{w}_h)$. Hence the above problem can be viewed as the error problem; compare this with the error equation in the linear case (1.35). Whence we have e_h, we can update $w_h = \hat{w}_h + e_h$ as in the linear case.

8.2 Nonlinear equation and unconstrained minimization

We first consider the unconstrained optimization problem

$$u = \operatorname*{argmin}_{v\in\mathcal{V}} \mathcal{F}(v). \tag{8.14}$$ eqn:min

If $\mathcal{F} : \mathcal{V} \mapsto \mathbb{R}$ is a convex function, then the problem is called a convex optimization (or convex programming). If $\mathcal{F}$ is differentiable, a minimizer satisfies the well-known *first-order optimization condition*

$$\mathcal{G}(u) := \mathcal{F}'(u) = 0, \tag{8.15}$$ eqn:1st-cond

where $\mathcal{G} : \mathcal{V} \mapsto \mathbb{R}$ is the Frechet derivative of $\mathcal{F}$. If $\mathcal{F}$ is convex, then (8.14) is equivalent for solving the nonlinear equation (8.15). In particular, if $\mathcal{F}$ is quadratic, then the problem is called a quadratic optimization. Apparently, if $\mathcal{F}$ is a convex quadratic functional, then the problem (8.14) is equivalent to our model problem (2.1), $\mathcal{A}u = f$, with an SPD operator $\mathcal{A} = \mathcal{G}'$.

Nonlinear solvers

In general, the problem (8.14) is much more difficult to solve than (2.1) due to its nonlinearity. We can employ a nonlinear iterative solver to linearize (8.15) to obtain a linear (differential) equation, i.e., linearization then discretization. For example, we may use the standard approaches,

like the Picard method or the Newton–Raphson method. Another strategy is to discretize the continuous problem (8.14) or (8.15) in order to obtain a nonlinear algebraic problem

$$u = \operatorname*{argmin}_{v\in\mathbb{R}^N} \mathcal{F}(v) \tag{8.16}$$ eqn:dmin

or

$$\mathcal{G}(u) = 0. \tag{8.17}$$ eqn:d1st-cond

The idea of coarse-grid correction used in Algorithm 3.1 does not apply any more here because the classical residual equation is linear. There are basically two approaches to apply the multilevel idea on this problem—The first approach is to linearize the problem and then apply multigrid methods to linear problems; The second one is to apply multigrid directly to the nonlinear problem using the so-called *Full Approximation Scheme* (FAS).

Newton–Raphson method

There are different ways to linearize a nonlinear problem like (8.15). For simplicity, we now only consider discrete version of the nonlinear equation, i.e., $\mathscr{V} = \mathbb{R}^N$. The most popular approach is the so-called Newton–Raphson (or Newton) linearization. We apply second-order Taylor expansion to approximate the objective function near the current iteration $u^{(k)} \in \mathbb{R}^N$, i.e.,

$$\mathcal{F}(u^{(k)} + e) \approx \mathcal{F}(u^{(k)}) + (\nabla\mathcal{F}(u^{(k)}),\, e) + \frac{1}{2}(\nabla^2\mathcal{F}(u^{(k)})e,\, e).$$

In order to find a good incremental correction step, we can consider

$$e^{(k)} = \operatorname*{argmin}_{e\in\mathbb{R}^N} \frac{1}{2}(\nabla^2\mathcal{F}(u^{(k)})e,\, e) + (\nabla\mathcal{F}(u^{(k)}),\, e) = -\big[\nabla^2\mathcal{F}(u^{(k)})\big]^{-1}\nabla\mathcal{F}(u^{(k)}).$$

This is the Newton–Raphson iteration

$$u^{(k+1)} = u^{(k)} - \big[\nabla^2\mathcal{F}(u^{(k)})\big]^{-1}\nabla\mathcal{F}(u^{(k)}). \tag{8.18}$$ eqn:Newton

In the above iteration step, we need to solve a linear system, the Jacobian equation:

$$\mathcal{A}e^{(k)} := \big[\nabla^2\mathcal{F}(u^{(k)})\big]e^{(k)} = -\nabla\mathcal{F}(u^{(k)}) =: r^{(k)}. \tag{8.19}$$ eqn:Jeqn

We can employ the methods discussed in the previous chapters to solve such equations.

Listing 8.1: Newton–Raphson method

```
1 Given an initial guess u ∈ 𝒱 and set r ← −∇ℱ(u);
2 while ‖r‖ > ε
3     solve the Jacobian equation ∇²ℱ(u)e = r;
4     find a good stepsize α > 0;
5     u ← u + α e;  r ← −∇ℱ(u);
6 end
```

The Newton-Raphson method converges very fast (second-order convergence) if the initial guess is close enough to the exact solution. So if a good initial guess is available, the main computation cost of the above algorithm is assembling the Jacobian systems and solving it to acceptable accuracy. If we apply a multigrid algorithm to solve the Jacobian systems, then this method is usually called Newton-Multigrid method. Similarly, another wide-used approach to apply a domain decomposition preconditioned Krylov method to solve the Jacobian systems, then this method is called Newton-Schwarz-Krylov method. Note that we might not need to assemble the Jacobian system explicitly; instead, we can use a Jacobian-free scheme.

Full approximation scheme

For the nonlinear equation (8.15), the residual corresponding to an approximate solution v can be defined as

$$r := -\mathcal{G}(v) = \mathcal{G}(u) - \mathcal{G}(v) \tag{8.20}$$ eqn:nonlinear

However, because $\mathcal{G}$ is not linear, $r \neq \mathcal{G}(u - v)$. In FAS, instead of considering the residual equation as in the linear case, the full equation is solved on the coarse grids.

We now use the following two-grid method to demonstrate the basic idea of FAS. Let $u^{(1)}$ be an approximate solution on the fine grid after several steps of relaxation. On the coarse grid, according to (8.20), we need to solve the following nonlinear equation

$$\mathcal{G}_c\big(u_c^{(1)}\big) - \mathcal{G}_c\big(\mathcal{I}_c^T u^{(1)}\big) = r_c = \mathcal{I}_c^T r = -\mathcal{I}_c^T \mathcal{G}\big(u^{(1)}\big). \tag{8.21}$$ eqn:nonlinear

This means, on the coarse level, a problem similar to the original problem (with different right-hand side) should be solved

$$\mathcal{G}_c\big(u_c^{(1)}\big) = \mathcal{G}_c\big(\mathcal{I}_c^T u^{(1)}\big) - \mathcal{I}_c^T \mathcal{G}\big(u^{(1)}\big). \tag{8.22}$$ eqn:nonlinear

Usually the right-hand side of the above equation is denoted as $\tau_c(u^{(1)})$ and is called the *tau correction*. Note that the coarse-level equation $\mathcal{G}_c$ can be obtained from the discretization on the coarse grid. We can also use the Galerkin method

$$\mathcal{G}_c(u_c) := \mathcal{I}_c^T \mathcal{G}(\mathcal{I}_c u_c).$$

Once the problem (8.22) is solved, we correct the approximation as

$$u^{(2)} = u^{(1)} + \mathcal{I}_c\big(u_c^{(1)} - \mathcal{I}_c^T u^{(1)}\big). \tag{8.23}$$ eqn:nonlinear

Apparently the above idea can be applied recursively as we discussed in §6.2. Because the coarse-grid problem is solved for the full approximation, rather than the error, the method is named as the Full Approximation Scheme. In this algorithm, evaluating the nonlinear function is usually the most expensive part computationally. We summarize the two-grid FAS algorithm as follows:

Listing 8.2: Full Approximation Scheme

```
1 Given an initial guess u ∈ V;
2 Solve the nonlinear equation G_c(u_c) = G_c(I_c^T u) − I_c^T G(u);
3 u ← u + I_c(u_c − I_c^T u);
```

Subspace correction methods for convex minimization

Apparently, the idea of subspace correction methods can be easily extended to unconstrained convex minimization problems here. The convergence analysis of SSC and PSC methods has been given by Tai and Xu [52].

8.3 Constrained minimization

In this section, we consider multilevel solvers for constrained minimization problems

$$u = \underset{v\in\mathcal{K}_0}{\operatorname{argmin}}\, \mathcal{F}(v) := \frac{1}{2}a[v,v] - f(v), \tag{8.24}$$ eqn:EVI2

which is equivalent to the variational inequality (8.2).

Projected full approximation method

Since the the above problem is nonlinear, we can apply the Full Approximation Scheme introduced in the previous section to solve this problem. And this is the so-called Projected Fully Approximation Scheme (PFAS) by Brandt and Cryer [18].

As we have discussed in the previous chapters, we first need to find a relatively simple iterative procedure which is able to dump the high-frequency part of the error quickly. In order to obtain a smoother for (8.24), we can employ the simple iterative methods discussed in §2.1 and then apply a projection step to ensure the new iteration stays in the feasible set. For example, if u^{old} is the previous iteration and u^{GS} is the iteration after one or several Gauss-Seidel sweeps, then $u^{\text{new}} := \max\{0, u^{\text{GS}}\} \in \mathcal{K}_0$ is the new iteration. This method is naturally called the Projected Gauss-Seidel (PGS) method.

At some point PGS will not reduce error efficiently any more, we then apply FAS to approximate the error on a coarser level and continue this procedure until the coarsest level where the nonlinear problem can be solved quickly and accurately. To ease the notation, we explain the idea using a two-grid algorithm for now. We first solve the general LCP problem on a fine level with a given right-hand side f_l

$$\begin{cases} \mathcal{A}u \geqslant f \\ u \geqslant 0 \\ \langle \mathcal{A}u - f, u\rangle = 0. \end{cases}$$

using the PGS method or some other smoother to obtain an approximate solution $u^{(1)}$. Then we solve the above LCP on a coarse level with the right-hand side

$$f_c := \mathcal{I}_c^T \left(f - \mathcal{A}u^{(1)}\right) + \mathcal{A}_c \mathcal{I}_c^T u^{(1)}$$

to obtain an approximation $u_c^{(1)}$. In turn, an improved approximation is given by

$$u^{(2)} = u^{(1)} + \mathcal{I}_c\big(u_c^{(1)} - \mathcal{I}_c^T u^{(1)}\big).$$

Interior point method

For simplicity, we now consider the constrained minimization problem (8.2) on the finite dimensional space $\mathbb{R}^N$, that is to say

$$u = \underset{v\geqslant 0,\, v\in\mathbb{R}^N}{\operatorname{argmin}} \mathcal{F}(v) := \frac{1}{2} v^T A v - f^T v. \tag{8.25}$$ eqn:disEVI2

In this case, the Lagrange multiplier $\sigma \in \mathbb{R}^N$ satisfies that $\sigma = -\mathcal{G}(u)$. Then we have the first-order optimality condition

$$\begin{aligned} &\sigma + \mathcal{G}(u) = 0, \qquad && \sigma \leqslant 0, \\ &U\sigma = 0, && u \geqslant 0. \end{aligned}$$

Here we use a convention often employed in the literature $U := \operatorname{diag}\{u_1, \ldots, u_N\}$; similarly, we will denote $\Sigma := \operatorname{diag}\{\sigma_1, \ldots, \sigma_N\}$.

The condition $U\sigma = 0$ (or equivalently, $u_i\sigma_i = 0$ for any $i = 1, \ldots, N$) is usually called the complementarity condition. We now try to relax this condition such that $U\sigma = \mu\mathbf{1}$, where μ is a positive penalty parameter and $\mathbf{1}$ is an all-one vector. At the same time, we try to maintain the iterative solution (u, σ) strictly in the primal-dual feasible set, i.e., $u > 0$ and $\sigma < 0$. Hence we need to solve a system of nonlinear equations:

$$\begin{cases} \sigma + \mathcal{G}(u) &= 0, \\ U\sigma - \mu\mathbf{1} &= 0. \end{cases}$$

We apply the Newton's method for this system and obtain an iterative method

$$\begin{cases} A\delta u + \delta\sigma &= -\sigma - \mathcal{G}(u) \\ \Sigma\delta u + U\delta\sigma &= \mu\mathbf{1} - U\sigma \end{cases} \quad \text{or} \quad \begin{pmatrix} A & I \\ \Sigma & U \end{pmatrix} \begin{pmatrix} \delta u \\ \delta\sigma \end{pmatrix} = \begin{pmatrix} f - Au - \sigma \\ \mu\mathbf{1} - U\sigma \end{pmatrix}.$$

Upon solving this linear system, we can obtain a new iteration. Furthermore, in the above system, I, Σ, and U are all known diagonal matrices, we only need to solve the Schur complement problem

$$(A - U^{-1}\Sigma)\delta u = \mu U^{-1}\mathbf{1} + f - Au. \tag{8.26}$$ eqn:IPmethod

Moreover, since $\sigma < 0$ and $u > 0$, the above equation is well-defined and the coefficient matrix is SPD. We can then apply a multilevel iterative method discussed in the previous chapters to solve it efficiently; see [5] for details.

Monotone multigrid method

Now suppose we hierarchical meshes, $\{\mathcal{M}_h^0, \ldots, \mathcal{M}_h^j\}$ and let A_l, b_l, $l = 0, \ldots, j$ are the stiffness matrices and right-hand-side vectors corresponding to the partition $\mathcal{M}_h^l$, respectively. As usual, $\mathcal{M}_h^j$ is the finest mesh. We denote the linear finite element space by V_h^l associated with mesh $\mathcal{M}_h^l$.

We need two kinds of orthogonal projections onto the finite element space V_h^l. The L^2-projections $Q_l : V_h^j \to V_h^l$ are defined by

$$(Q_l v_h, \phi_l) = (v_h, \phi_l), \quad \phi_l \in V_h^l, \tag{8.27}$$ eqn:L2proj

and the energy projections $\Pi_l : V_h^j \to V_h^l$ by

$$a[\Pi_l v_h, \phi_l] = a[v_h, \phi_l], \quad \phi_l \in V_h^l. \tag{8.28}$$ eqn:proj

We first define multigrid methods recursively. For a given initial guess $w_j^{(0)} \in V_h^j \bigcap \mathcal{K}_0$. A coarse grid correction is performed: computing the approximate defect $e_{j-1}^{(0)} = \Pi_{j-1}(w_h - w_j^{(0)}) \in V_h^{j-1}$ as the solution of the quadratic programming problem

$$\min_{e_{j-1}^{(0)} \in V_h^{j-1},\ w_j^{(0)} + e_{j-1}^{(0)} \bigcap \mathcal{K}_0} \frac{1}{2} a[e_{j-1}^{(0)}, e_{j-1}^{(0)}] - \langle \sigma(w_j^{(0)}), e_{j-1}^{(0)} \rangle. \tag{8.29}$$ eqn:corrction

Then let $w_j^{(1)} = w_j^{(0)} + e_{j-1}^{(0)}$. Then we apply m steps of post-smoothing scheme, like *projected SOR* to obtain $w_j^{(m+1)}$. For the coarse correction step, instead of solving the problem on the coarser level $j-1$ exactly, we can solve it by the same multigrid procedure described here. In this way, we obtain a recursive multigrid V-cycle. If we perform coarse grid correction twice at each level, then we get a W-cycle.

One problem with this procedure is that e_{j-1} and w_j are in different levels. To avoid this difficulty, we propose the following coarse grid correction scheme instead of (8.29):

$$\min_{d_{j-1}^{(0)} \in V_h^{j-1} \bigcap \mathcal{K}_0} \frac{1}{2} a[d_{j-1}^{(0)}, d_{j-1}^{(0)}] - \langle \sigma(w_j^{(0)}), d_{j-1}^{(0)} \rangle. \tag{8.30}$$ eqn:newcorrct

And then $w_j^{(1)} = w_j^{(0)} + d_{j-1}^{(0)}$ which is always in $\mathcal{K}_0$ because both $w_j^{(0)}$ and $d_{j-1}^{(0)}$ are in $\mathcal{K}_0$ by definition. It is easy to check that the local obstacles in this method are monotone in the sense of Kornhuber [39]. Then we get the similar V-cycle or W-cycle multigrid method as for linear problems expect we need to add a projection step to project the iterates to $\mathcal{K}_0$.

Remark 8.3. This method is shown to be not very good by Tai's test example. The reason is that the coarse grid correction only works when the current approximation is less than the exact solution in the method.

8.4 Constraint decomposition method

It is known the general V-cycle can be written as a successive subspace correction method. For a sequence of search directions $\{\phi_i\}_{i=1}^N$ such that $V_h^j := \text{span}\{\phi_i\}_{i=1}^N$. We can construct a numerical method for find the minimizer of (8.12) as a sequential quadratic programming method. Starting from an initial guess $w_j^{(0)} \in V_h^j \bigcap \mathcal{K}_0$, at each iteration, we solve

$$\min_{w_j^{(0)}+\alpha\phi_1 \in V_h^j \bigcap \mathcal{K}_0} \frac{1}{2} a[w_j^{(0)} + \alpha\phi_1, w_j^{(0)} + \alpha\phi_1] - f_0(w_j^{(0)} + \alpha\phi_1). \tag{8.31}$$ eqn:ssc

Similar to the discussion in the previous section, we need to solve a discrete problem

$$\min_{w_j^{(0)}+\alpha\phi_1 \in V_h^j \bigcap \mathcal{K}_0} \frac{1}{2} a[\phi_1, \phi_1]\alpha^2 - \langle \sigma(w_j^{(0)}), \phi_1 \rangle \alpha. \tag{8.32}$$ eqn:ssc2

Then the new iterate is obtained by $w_j^{(1)} = w_j^{(0)} + \alpha\phi_1$. Similarly, we start from $w_j^{(1)}$ and search in the direction ϕ_2 to obtain $w_j^{(2)}$, and so on.

If we choose $\text{span}\{\phi_i\}_{i=1}^N$ as the canonical nodal basis of V_h^j, then it is just usual nonlinear or projected Gauss-Seidel method. To take advantage of multilevel basis, it is natural to choose $\text{span}\{\phi_i\}_{i=1}^N = \{\phi_1^j, \ldots, \phi_{N_j}^j, \phi_1^{j-1}, \ldots, \phi_{N_{j-1}}^{j-1}, \ldots, \phi_1^1, \ldots, \phi_{N_1}^1\}$. It falls into the category of *extended relaxation methods*. The problem with this procedure is that ϕ_i might not be in the finest level j, which costs extra computation effort to enforce the constraints $w_j^{(i-1)} + \alpha\phi_i \in V_h^j \bigcap \mathcal{K}_0$.

See Tai [51] for details.

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