

Generation of ultra-compressed solitons with a high tunable wavelength shift in Raman-inactive hollow-core photonic crystal fibers

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Abstract:

The fission of N-solitons is a key mechanism leading to the supercontinuum generation and creation of hyper-compressed pulses and solitons featuring strong wavelength tuning [1, 2]. Recent advances in manufacturing PCFs filled with Raman-inactive gases [3] provide a strong motivation to focus the studies on the fission driven by the third-order dispersion (TOD) [4], and potential applications of this setting to photonics. In the absence of the Raman-induced self-frequency shift and the Raman-associated noise, fission-produced strongly compressed solitons, once generated, may propagate keeping constant internal frequencies. A higher-order N-soliton, $u = N\sqrt{P_0} \text{sech}(T/T_0)$ is launched into the fiber, where T_0 and P_0 are width and peak power of the corresponding fundamental soliton. If the TOD dispersion β_3 is very small, it can be considered as a perturbation added to the second-order dispersion, β_2 . In this case, the peak power of each fundamental soliton emerging after the splitting of the N-soliton is given by the classical result, $P_j = P_0(2N - 2j + 1)^2$ [5]. If δ_3 is larger, it leads to a significant increase in the largest fundamental-soliton's peak power and compression degree, along with the increase of the wavelength shift. However, further increase of β_3 leads to a loss of the peak-power enhancement. For optimal pulse compression and wavelength conversion, universal optimal value of TOD strength parameter $\delta_3 \equiv \beta_3/(6\beta_2 t_0)$ was found. This optimal value is valid for any pulse duration, second- and third-order dispersion coefficients, depending solely on order N of the injected soliton. The optimal pulse-compression degree significantly exceeds the well-known analytical prediction [5].

On the contrary to the insensitivity of the solitons' peak powers to the signs of δ_3 , the wavelength upshift and downshift significantly differ for the opposite signs of δ_3 . The downshift of strongest solitons wavelength is opposed by the proximity to the zero dispersion point of PCF, beyond which solitons cannot propagate. The highest order in Fig.1 is $N = 15$. For still larger N, one can achieve even higher peak-power ratio and wavelength shift, but, due to interaction between multiple generated solitons and trapped dispersive waves [6, 7], the control over parameters of the emerging solitons deteriorates. The physical mechanism beyond the phenomenon, which is valid also in the presence of the self-steepening effect, is the power and momentum absorption by the strongest newly born soliton in the course of inelastic interactions with weaker pulses after the decomposition of the initial N-soliton [4,8].

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